

# Surface Quasi Geostrophic Reconstruction of Vertical Velocities and Vertical Heat Fluxes in the Southern Ocean: Perspectives for SWOT

Elisa Carli<sup>1</sup>, Lia Siegelman<sup>2</sup>, Rosemary Morrow<sup>1</sup>, Oscar Vergara<sup>1,3</sup>

<sup>1</sup>Laboratoire d'Etudes en Geophysique et Oceanographie Spatiale (LEGOS) [CNES, UT3, CNRS, IRD]

<sup>2</sup>Scripps Institution of Oceanography, UC San Diego

<sup>3</sup>Collecte Localisation Satellites

## Key Points:

- Vertical velocities can be coherently reconstructed in the Southern Ocean at scales down to 30-40 km
- Vertical heat fluxes can be retrieved using effective surface quasi-geostrophic reconstruction and are dominated by vertical velocities
- SWOT fine-scale observations allow for a reconstruction of vertical velocities with 45% more variance compared to traditional altimetry

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Corresponding author: Elisa Carli, [elisa.carli@univ-tlse3.fr](mailto:elisa.carli@univ-tlse3.fr)

## Abstract

Mesoscale currents account for 80% of the ocean’s kinetic energy, whereas submesoscale currents capture 50% of the vertical velocity variance. SWOT’s first sea surface height (SSH) observations have a spatial resolution an order of magnitude greater than traditional nadir-looking altimeters and capture mesoscale and submesoscale features. This enables the derivation of submesoscale vertical velocities, crucial for the vertical transport of heat, carbon and nutrients between the ocean interior and the surface. This work focuses on a mesoscale energetic region south of Tasmania using a coupled ocean-atmosphere simulation at km-scale resolution and preliminary SWOT SSH observations. Vertical velocities ( $w$ ), temperature anomalies and vertical heat fluxes (VHF) from the surface down to 1000 m are reconstructed using effective surface Quasi-Geostrophic (sQG) theory. An independent method for reconstructing temperature anomalies, mimicking an operational gridded product, is also developed. Results show that sQG reconstructs 90% of the total  $w$  and VHF rms at scales down to 30 km just below the mixed layer and 50-70% of the rms for scales larger than 70 km at greater depth, with a spatial correlation of  $\sim 0.6$ . The reconstruction is spectrally coherent ( $> 0.65$ ) for scales larger than 30-40 km at the surface, slightly degrading ( $\sim 0.55$ ) at depth. The two temperature anomalies datasets yield similar results, indicating the dominance of  $w$  on VHF. RMS of sQG  $w$  and VHF derived from SWOT are twice as large as those derived from conventional altimetry, highlighting the potential of SWOT for reconstructing energetic meso and submesoscale dynamics in the ocean interior.

## Plain Language Summary

This work focuses on vertical ocean dynamics induced by small-scale ocean surface currents, especially those at the mesoscale and submesoscale (20-100 km in diameter). These ocean vertical velocities pump heat and carbon from the surface to depth and carry nutrients important for biomass development. Our motivation is the first data from the new satellite SWOT (Surface Water and Ocean Topography), which provides observations of surface geostrophic currents at much higher resolution compared to traditional satellites. We study a Southern Ocean area near Tasmania using a realistic coupled ocean-atmosphere simulation at a km-scale resolution and initial SWOT data and reconstruct vertical velocities, temperature, and heat fluxes using a vertical projection theory. Results show that the vertical reconstruction coherently represents ocean vertical velocities and heat fluxes below the surface mixed layer larger than 30-40 km and that smaller scales are not well reconstructed. In the computation of vertical heat fluxes, we compare two different ways of measuring temperature anomalies, finding that, overall, the difference is not significant because vertical velocities dominate the heat fluxes. Additionally, this study highlights the potential of SWOT for studying small-scale ocean dynamics.

## 1 Introduction

Studying oceanic vertical exchanges is paramount because the global ocean is responsible for storing up to 93% of the energy imbalance due to global warming (Rhein et al., 2013), and has absorbed about 30% of total anthropogenic carbon dioxide emissions since the 1980’s (Intergovernmental Panel On Climate Change (Ipcc), 2022). Observing the dynamical processes that carry this heat and carbon into the ocean interior is still a major challenge in oceanography. The ocean is also highly turbulent. Cyclonic and anticyclonic eddies, ubiquitous and very energetic mesoscale structures, contribute to the energetics of the oceanic circulation, particularly in the western boundary currents and the Antarctic Circumpolar Current (ACC). Mesoscale eddies (50-300 km) impact the lateral stirring and structure of tracers, setting the ocean distribution of heat, salt, carbon, nutrients and biomass (Franks et al., 1986; Strass, 1992; Mahadevan & Archer, 2000; A. P. Martin et al., 2002). Their primary role is now recognised also in driving global

upward vertical heat transport (Rintoul, 2018; Llorc et al., 2018). Griffies et al. (2015) found that the representation of the mesoscale in climate models can largely influence lateral and vertical heat transport. Therefore, it is fundamental to include a rich representation of the ocean mesoscale in model simulations or employ a correct parametrization. However, numerical models do not fully agree on the magnitude of the eddies' impact on the vertical tracer distribution (McGillicuddy et al., 2007) and observations of these 3D dynamical structures are limited. *In-situ* observations of vertical velocities are rare, difficult to measure and have limited coverage in space and time. Satellite infrared and colour images show submesoscale horizontal filamentary structures between the eddies, which greatly contribute to the nutrient supply, closing the nutrient budget (Lévy, 2024), but their observations are limited to the ocean surface and in clear sky conditions. Traditional altimetry surface geostrophic velocities have a resolution of around 150 km (Ballarotta et al., 2019) and are thus not able to resolve the small and very energetic scales predominant in the Southern Ocean, where the Rossby radius is small.

The Surface Water and Ocean Topography (SWOT) wide-swath altimeter mission launched in December 2022 (Morrow et al., 2019; Fu et al., 2024) brings a new 2D observation of these scales, being able to observe small ocean structures of 5-10 km with a time coverage in the order of a few days or better at high latitudes. The first SWOT data indicate that the spectral level noise is even lower than the requirement by one order of magnitude at small scales, showing ocean circulation processes at scales down to  $\sim 5$  km (Fu et al., 2024). After careful editing and processing (Dibarboue et al., 2023), this allows a more precise estimate of ocean sea level elevation, and its derived geostrophic currents, vorticity and strain. These are key properties for reconstructing the 3D ocean circulation from 2D surface satellite observations, which has long been a goal in oceanography.

Different methodologies have been developed to extend surface observations into the ocean interior and estimate vertical velocities as the vector for the vertical transport of heat, carbon, and nutrients. The classical Omega equation (Hoskins et al., 1978) has been widely used to study the balanced motions associated with fronts (Tintoré et al., 1991; Rudnick, 1996; A. Martin & Richards, 2001; Allen et al., 2001), but the availability of high-resolution data within mesoscale eddies remains a challenge for its application on a larger scale. To explain energy redistribution through different scales, the interior-quasi-geostrophic (QG) and surface-quasi-geostrophic (sQG) theories are often applied. QG is based on a small Rossby number approximation (Charney, 1971; Rhines, 1979). Motions are assumed not to be influenced by vertical boundary conditions and density anomalies are assumed to be zero (Klein & Lapeyre, 2009). McWilliams (1989) found that increasing the resolution in three dimensions does not significantly change the QG turbulence. sQG, on the other hand, considers nonzero surface density perturbations, which cascade to small scales leading to small-scale surface frontogenesis (Lapeyre & Klein, 2006). In this case, the potential vorticity (PV) is assumed to be uniform in the interior of the domain. The main limitation of sQG is that when the resolution is high, small-scale structures may have high vorticity values of the order of  $f$  (the Coriolis parameter). These structures depart from sQG and are subject to ageostrophic frontogenesis (Klein et al., 2008). Lapeyre and Klein (2006) define an "effective" sQG model based on the same assumptions, which allows the diagnosis of the 3D dynamics in the upper ocean from either the sea surface height (SSH) or surface buoyancy. This present work explores this "effective" sQG solution. The gain in spatial resolution from the new SWOT observations motivates renewed interest in using sQG with altimetry data to infer the 3D dynamical structures in the energetic Southern Ocean region (Wang et al., 2013; Su et al., 2018; Qiu et al., 2020).

The primary objective of this work is to assess the potential of the effective sQG methodology to reconstruct the vertical velocities and vertical heat fluxes within the first 1000 m of the water column in the energetic Polar Front region of the ACC, downstream

of the SouthEast Indian Ridge. This region was chosen since an intensive field campaign was undertaken in November and December 2023, which will provide in-situ validation of the vertical reconstruction technique in a second study. The region is dynamically inhomogeneous, with strong water mass variations across the energetic Polar front. Whilst it has relatively flat bottom bathymetry, the energetic dynamics have strong topographic control from ocean ridges upstream and downstream (Figure 1), and we will explore the sQG reconstruction capabilities in this complex dynamical regime. We analyze two five-day periods in March and October taken from a realistic coupled ocean-atmosphere simulation having a similar resolution to SWOT, to assess the reconstructed dynamics in two contrasting seasons. This analysis is not intended to be a full quantitative representation of all the seasonal dynamics and variability in the region, which is beyond the scope of this work.

Section 2 provides some background on the region of interest, motivating its choice for this study. The data and methods are presented in Section 3. Section 4 includes the modelling study of the 3D dynamics of the region and the reconstructions of vertical velocity, temperature and vertical heat flux based on a realistic model. Finally, the conclusion and discussion are provided in Section 5, together with a preliminary first test with SWOT data in preparation for a full validation with in-situ data that will be part of a follow-up study.

## 2 Meso- and submesoscale dynamics in the Southern Ocean

The Southern Ocean is a critical region for its strong vertical atmospheric heat and carbon uptake. Even though the Southern Ocean represents only 30% of the ocean surface, it accounts for 60% of the global heat content trend since 1970. Sallée (2018) suggests that the poor observations in the region might bias this result and that in recent times, with the increasing measurements from the autonomous Argo profilers, the Southern Ocean contribution to heat storage is between 67% and 98% of the global mean. The mechanisms that influence the capacity of the ocean to move heat and carbon are various, and cover spatial scales from centimetres to kilometres and temporal scales from minutes to decades (Ferrari & Wunsch, 2008).

Within the Southern Ocean, the role of mesoscale to submesoscale eddy processes is critical. The meandering ACC promotes exchanges between the three ocean basins contributing to the global overturning circulation, which transports heat, moisture and carbon dioxide around the globe. The ACC also isolates cooler, fresher Antarctica waters from the warmer, saltier subtropical waters in the north. Watts et al. (2016) found that south of the ACC the ocean loses 0.4 PW of net heat to the atmosphere, and a large horizontal poleward oceanic heat flux is required to balance this loss. However, due to the zonal flow of the ACC, no strong meridional (N-S) mean currents carry this heat poleward in the upper 2000 m above the main bathymetric ridges. Instead, strong poleward heat fluxes are carried by mesoscale eddies across the ACC. These have been estimated from models (Trenberth & Caron, 2001; Boccaletti et al., 2005; Wolfe et al., 2008), from in-situ moorings and altimetry (Phillips & Rintoul, 2000; Marshall et al., 2006; Watts et al., 2016; Foppert et al., 2017) and from Argo float observations (Sun et al., 2019). Mesoscale eddies also play a role in modulating the ACC dynamics (Sallée, 2018), and the Southern Ocean's eddy kinetic energy (EKE) evolves interannually (Martínez Moreno et al., 2020) and contributes to the dynamical adjustment to climate variations (Hogg et al., 2008). Today, most of these Southern Ocean observational studies on the role of the horizontal transfers of heat, salt, carbon and nutrients come from large-scale Argo observations or altimetric currents that only resolve the large mesoscales. The role of horizontal stirring of fluxes at small mesoscales or submesoscales has been estimated only at localised sites (A. P. Martin et al., 2001; Naveira Garabato et al., 2011; Nencioli et al., 2013; Jakes et al., 2024) or from models (Wolfe et al., 2008; Le Sommer et al., 2011).

Recent studies show that the submesoscale also plays an important role. Su et al. (2020) found that the upper-ocean submesoscale produces a mean winter upward heat transport of more than  $100 \text{ Wm}^{-1}$  at latitudes of  $60^\circ$ , that is five times higher than the mean mesoscale heat transport, based on a high-resolution  $1/48^\circ$  model. Compared to a lower resolution model, once the submesoscale is resolved, it contributes to warming the sea surface by up to  $0.3^\circ\text{C}$  and an annual air-sea heat flux anomaly of  $10 \text{ Wm}^{-2}$ . Submesoscale structures were initially thought to be confined within the mixed layer (Callies et al., 2015; Fox-Kemper et al., 2008) because of their assumed quasigeostrophic (QG) balance which would prevent the formation of strong density gradients at depth. However, recent modelling studies by Siegelman (2020) and Siegelman et al. (2020) have shown an enhanced vertical heat flux within deep submesoscale ocean fronts in the Kerguelen region in spring and summer, and Yu et al. (2019) found the largest submesoscale vertical velocities extending down to at least 200 m below the mixed layer base in winter and spring. Frontogenesis also penetrates from the ocean surface to below the ML base in the presence of deep fronts (Thomas & Ferrari, 2008). In addition to their dynamical role, Brannigan (2016) found that submesoscale resolving models present an enhanced biological activity in anticyclonic eddies that could lead to the upwelling of nutrient-rich waters and other biological activity within the eddies. Similar results arise with in-situ observations. Using a combination of satellite altimetry and CTD data, Siegelman et al. (2020) reported vertical heat fluxes with a local amplitude reaching  $2000 \text{ Wm}^{-2}$  at the location of submesoscale fronts in the Kerguelen region, with vertical exchanges occurring both within and below the mixed layer (ML). The mechanisms we will focus on in this work are those driven by the vertical velocity from deep-reaching eddy stirring below the ML.

The area we analyse is located southeast of Tasmania (Figure 1), within a large meandering of the Polar Front. This is a region of energetic mesoscale and submesoscale dynamics (Foppert et al., 2017), downstream of the Southeast Indian Ridge (Cyriac et al., 2023). In this region, characterised by a topographic obstacle, standing meanders or lee waves are generated (Hughes et al., 1998) and many of the key physical processes central to the Southern Ocean heat and momentum balance are concentrated (Rintoul, 2018). Particularly, a poleward transport of heat (Meijer et al., 2022) and biogeochemical tracers (Patel et al., 2020) are generated by the strong mesoscale and submesoscale activity in the region, which makes the ACC “leak” (Naveira Garabato et al., 2011).

Numerous in-situ data were collected here in late 2023, simultaneously with SWOT observations, in a collective effort for the calibration and validation of the satellite (d’Ovidio et al., 2019). In late 2023 the FOCUS campaign sampled the area at high resolution with a towed Triaxus, CTDs (conductivity, temperature, and depth), gliders, and deployed moorings (<https://www.swot-adac.org/campaigns/acc-smst/>). In the same period, as part of the Survostral campaign (<https://www.legos.omp.eu/survostral>), more than 400 expendable bathythermographs (XBTs) have been launched between Tasmania and Antarctica to sample the upper-ocean temperature structure at very high resolution (5-7 km) in the region of strong frontal activity. These data will be used in a more extensive validation of SWOT data, including for the vertical velocity and heat flux reconstruction in future work.

### 3 Data and methods

#### 3.1 Coupled Ocean Atmosphere Simulation (COAS)

Here, we use a new global coupled ocean-atmosphere simulation at high resolution. COAS has been used by Torres et al. (2022), and we refer the reader to this publication for more details on the simulation and its validation. Briefly, however, the COAS simulation is generated by coupling the Goddard Earth Observing System (GEOS) atmospheric model to the Massachusetts Institute of Technology general circulation ocean model

(MITgcm). For the atmospheric component, GEOS is configured to use the C1440 cubed-spherical grid, and a detailed description of the model configuration can be found in Molod et al. (2015). The atmosphere has a nominal horizontal grid spacing of 6.9 km and 72 vertical levels. The oceanic component uses the same configuration as the Latitude Longitude-polar Cap 2160 (LLC2160) simulation, detailed in Arbic et al. (2018), permitting mesoscale and submesoscales as well as more realistic coupled upper ocean exchanges than a forced simulation. We solely analyse the ocean component, configured with a nominal horizontal grid spacing of  $1/24^\circ$  and 90 vertical levels with a non-uniform grid spacing of 1 m at the surface, increasing to around 300 m at the bottom and tidal forcing. The simulation period covers 14 months, from 20 January 2020 to 25 March 2021, with the first two months corresponding to the spin-up period. In this study, to test the vertical sQG reconstruction, we extract two five-day periods in contrasting seasons, 25-29 March 2020 and 25-29 October 2020. In the study region, the resolution of  $1/24^\circ$  corresponds to an average grid point of 2.3 km, comparable to the SWOT L3 ocean product's 2 km grid. Although other simulations exist with higher resolution (e.g., LLC4320 at  $1/48^\circ$  (Rocha et al., 2016)), COAS' horizontal resolution at our latitudes is sufficient to capture the physical features of interest for our study and its observation with SWOT. In addition, COAS's weaker (and more realistic) tidal forcing compared to LLC4320 allows for better comparison between model outputs of  $w$  and reconstructed  $w$ .

### 3.2 Effective Surface Quasi Geostrophic (sQG) theory

Surface quasi-geostrophic (sQG) theory used in this work is the "effective" version defined by Lapeyre and Klein (2006). It is based on the reduction of Eady's model for baroclinic instabilities from one boundary (here the ocean surface), used to describe the oceanic mesoscale eddy fields in the surface layers (Klein & Lapeyre, 2009). This approach is valid for the surface intensified eddies. The reconstruction of other eddies, namely subsurface intensified eddies, requires the combination of the effective sQG with the interior sQG (iSQG) as in Wang et al. (2013), where both have a signature in SSH. The major assumption for effective sQG is that potential vorticity (PV) is uniform and constant in the ocean interior. Using PV inversion, one can retrieve a geostrophic stream function ( $\psi$ ) related to PV using an elliptic operator (Klein et al., 2009). Following Klein et al. (2009) and assuming a doubly periodic domain, the geostrophic stream function is retrieved at all depths ( $z$ ) from SSH ( $\eta$ ) using an exponential decay as:

$$\widehat{\psi}(\mathbf{k}, z) = \frac{g}{f} \widehat{\eta}(\mathbf{k}) \exp\left(\frac{N_0}{f} kz\right) \quad (1)$$

where  $\widehat{(\cdot)}$  corresponds to the Fourier transform,  $\mathbf{k} = (k_x, k_y)$  is the wavenumber vector and  $k = |\mathbf{k}|$  is its norm,  $g$  is the gravity constant,  $f$  is the Coriolis frequency,  $N_0$  is a Brunt-Väissälä frequency that accounts for the contribution of the interior PV. Vertical motions are retrieved from the buoyancy equations as in Klein et al. (2009)

$$\widehat{w}(\mathbf{k}, z) = -\frac{c^2}{N_0^2} \left[ -J(\widehat{\psi}_s, b_s) \exp\left(\frac{N_0}{f} kz\right) + J(\widehat{\psi}, b) \right] \quad (2)$$

where the subscript  $s$  refers to the surface values,  $b$  is the buoyancy,  $c$  is a unitless constant that helps optimise the reconstructed fields' amplitude and  $J(A, B) \equiv (\partial_x A \partial_y B - \partial_y A \partial_x B)$ . This methodology requires the optimization of the constants  $N_0$  and  $c$ . We refer the reader to Lapeyre (2017) for a full derivation of sQG theory.

Following Klein et al. (2009), the SSH field is filtered in space before applying sQG. We apply a two-dimensional Lanczos low-pass filter in the spectral domain to eliminate all spatial scales smaller than 20 km (Duchon, 1979). The reconstructions obtained with the filtered SSH are compared to the COAS fields also filtered with the same method. This choice derives from the study of Klein et al. (2009), who assesses that sQG is non-realistic below these scales. This is confirmed by comparing the COAS and sQG verti-



cal velocity wavenumber spectra before the 20 km filter is applied (not shown), revealing a clear drop in energy in the sQG w reconstruction at scales smaller than 20 km at the surface and at larger scales at depth. To reduce the residual noise from the SWOT data due to surface waves and other perturbations (Tréboutte et al., 2023), the same filter is later applied to the SSH fields of SWOT before application of sQG, ensuring consistency with the model results. We estimate that 20 km is a reasonable scale to match the effective capability of the gridded SWOT product.

In this study, since we have access to the full model vorticity and velocity field, we optimize  $N_0$  so that the shape of the model and sQG vorticity rms has a minimal mean squared error from the bottom of the mixed layer to 1000 m. The mixed layer average is around 100 m in March and 250 m in October over the full domain. The optimal  $N_0$  for this domain is  $3.72\text{e-}6 \text{ s}^{-1}$  in March and  $3.98\text{e-}6 \text{ s}^{-1}$  in October. We keep separate values to reproduce the seasonality of the vertical stratification. With these values, the correlation between the COAS and sQG vorticity fields is around 0.9 for the full water column (not shown).

To optimize  $c$  we work directly with the density anomalies and  $w$  fields, having an SSH filtered at 20 km before applying sQG and obtaining one optimal parameter  $c$  for the density anomaly reconstruction and one for the vertical velocity for each season. For the optimization, following Klein et al. (2009), we aim to align the maximum amplitudes so that, on average, the difference between the 2D COAS and the reconstructed physical fields between the bottom of the mixed layer and 1000 m is minimized. First, the model density anomaly and  $w$  corresponding to the maximum amplitude values of magnitude occurring outside  $1\sigma$  are extracted. This selects 30% of density anomaly points in March and 35% in October for each vertical layer on the water column and 23% (March and October) for the vertical velocity. We find an average optimal  $c$  of 1.2 for the vertical velocity and 1.51 for the density, which applies to both seasons.

We use various diagnostics to assess the quality of the sQG reconstruction. These include the 5-day average of the spatial correlation of the physical parameters between the COAS and the sQG fields and the vertical profile of average rms values for each depth. Wavenumber spectra at different depths are used to quantify which scales are well represented and where we lose the spectral energy, identifying those scales where the sQG reconstruction is ineffective in representing the dynamics.

### 3.3 Reconstruction of temperature anomalies

In order to calculate vertical heat fluxes, we need to know the vertical profile of temperature anomaly. In this work, we investigate the use of temperature anomalies retrieved from two different sources: from a pseudo-gridded product from observational data (section 3.3.1) and from a sQG reconstruction of density anomalies (section 3.3.2).

#### 3.3.1 Temperature anomalies from an observational gridded product

We test the computation of the VHF with temperature anomalies based on existing gridded products like ARMOR 3D (Guinehut et al., 2012). This product is built by merging in-situ observations of temperature and salinity (XBTs, CTDs, moorings and Argo profiles) with remote sensing observations of sea level anomaly and sea surface temperature (Barceló-Llull et al., 2018). The operational product provides weekly 3D temperature, salinity, geopotential height and geostrophic currents globally on a  $1/4^\circ$  grid. It has been validated by Mulet et al. (2012) against model reanalysis and in-situ observations (ANDRO current velocities from Argo float displacements and velocity measurements from the RAPID-MOCHA current meter array). The effective resolution at our latitudes is around 150 km, comparable to the DUACS altimetry products (Ballarotta et al., 2019). Here, we build a similar product with the same resolution in our region by

filtering the full model temperature at 150 km and using this field for the reconstruction of the VHF. The anomaly is computed spatially, with respect to the original field linearly detrended on longitudes and latitudes. In the following, this temperature product will be referred to as Argo-like temperature anomaly.

### 3.3.2 Temperature anomalies from sQG density anomalies

In some applications, we could be interested in computing the VHF only using satellite altimetry. In this case, temperature anomaly can be reconstructed from the sQG density anomaly by constructing a linear regression between the model temperature and density for each depth level. For this study, one regression is built for each vertical layer and season, and it is used to reconstruct the temperature anomaly profiles from the 3D sQG density anomalies. The regression is stable in time over the 5 days for each season. The temperature anomalies computed with this method are then used for the computation of the reconstructed VHF. In this frontal region, we test the dominance of temperature or salinity in determining the density field structure and tune the regression for values below the mixed layer above 1000 m. In March and October, the temperature gradients dominate the density gradients above 600 m and 800 m, respectively. Below these levels, the temperature derivation is less accurate. The strengths and drawbacks of the method are discussed in Section 5 along with alternative methods to retrieve temperature profiles from high-resolution SSH observations. The details of this methodology are given Appendix A.

### 3.4 Vertical Heat Fluxes (VHF)

VHF are computed following Su et al. (2018); Siegelman et al. (2020) as  $VHF = \rho_0 C_p w' T'$ , with  $\rho_0 = 1023 \text{ kg m}^{-3}$  the reference seawater density,  $C_p = 3985 \text{ J kg}^{-1} \text{ K}^{-1}$  the specific heat capacity of seawater,  $w$  the vertical velocity,  $T$  the temperature. The prime refers to the spatial anomaly of the linearly detrended 2D domain averaged values for each vertical layer. Positive (negative) values indicate an upward (downward) flux. VHF from COAS are computed with the filtered (with the low-pass filter of 20 km introduced in Section 3.2) outputs of vertical velocities and temperature, consistent with the vertical velocity field mentioned above. VHF reconstructions are computed with sQG vertical velocities and with temperature anomalies either from the Argo-like product or from sQG reconstruction, as detailed above. As for the vertical velocity reconstruction, diagnostics of the reconstructed VHF in the two cases are compared with VHF derived from COAS, using 2D spatial correlations and spatial rms for each layer of the water column to assess the quality of the VHF reconstructions.

## 4 Results

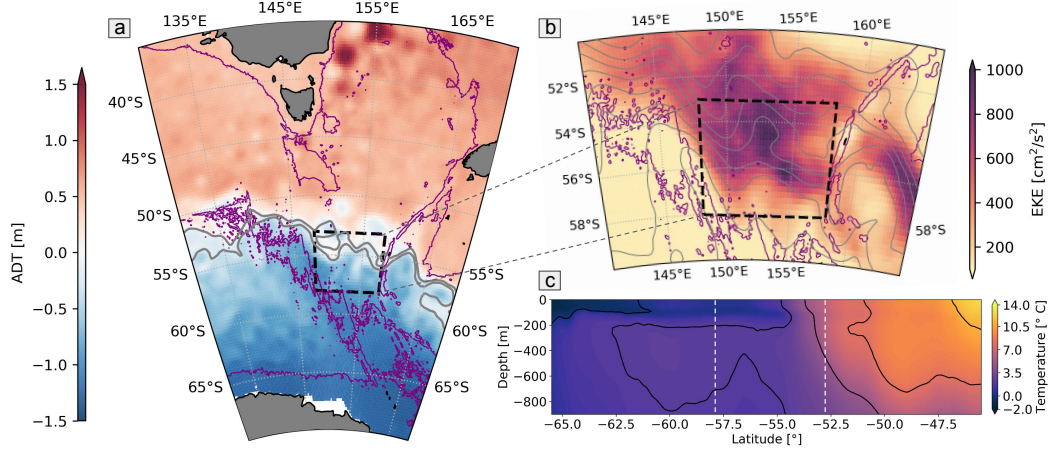
Our aim is to understand how SWOT data may be used to reconstruct vertical heat fluxes from the SSH observations alone in the region of the energetic Polar Front south of Tasmania. First, we investigate the main dynamical structures in the region. Then, we focus on the sQG reconstruction of vertical velocities and vertical temperature anomalies from the high-resolution COAS coupled simulation. Finally, the potential of reconstructing vertical heat fluxes from SSH fields is discussed, and a test case is run with early SWOT data.

### 4.1 Dynamical properties of the study region south of Tasmania

The study region is the energetic EKE hot-spot of the Antarctic Circumpolar Current (ACC) southeast of Tasmania (Figure 1a), extending North and South of the local Polar Front (PF) (Dong et al., 2006). This region has strong EKE, reaching over  $1000 \text{ cm}^2 \text{ m}^{-2}$  on average over the last 30 years (Figure 1b). Large mesoscale eddies have been



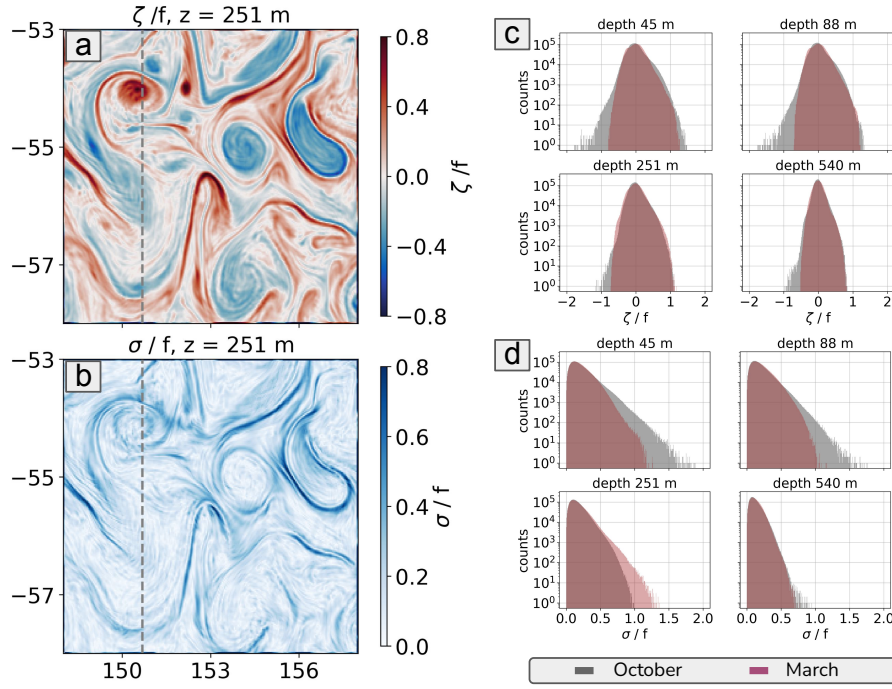
tracked crossing the front, cold-core eddies transport cold water from south of the front, and vice versa for warm-core eddies: both contributing a strong poleward heat transport (Morrow et al., 2004; Patel et al., 2019). Recent studies have also highlighted the important vertical velocities generated in such regions, leading to enhanced upwelling of Circumpolar Deep Water around the major topographic Southern Ocean features such as Southeast Indian Ridge or the Macquarie Ridge around 160°E (Tamsitt et al., 2017).



**Figure 1.** (a) CMEMS Absolute dynamic topography (ADT) average over seven years (2015 to 2022) in the energetic Southern Ocean region between Tasmania and Antarctica. The purple lines represent the 3000 m isobath, and the grey contours represent the Subantarctic Front (SAF) and the Polar Front (PF) taken from the CTOH/AVISO products. The black dashed line delimits the region of interest for the present study. (b) Average of DUACS dt2021 Eddy Kinetic Energy (EKE) from 1993 to 2022. The grey contours represent the SSH average from -1 m to 1 m spaced by 0.2 m. (c) CMEMS vertical temperature section between 0-900m, averaged over longitudes 148-158°E and between 2020 and 2021 between Tasmania 43°S and Antarctica 67°S. The black contours represent the 0°C, 2°C, 5°C, 8°C, 11°C isotherms. The white lines represent the latitude limit of the studied region

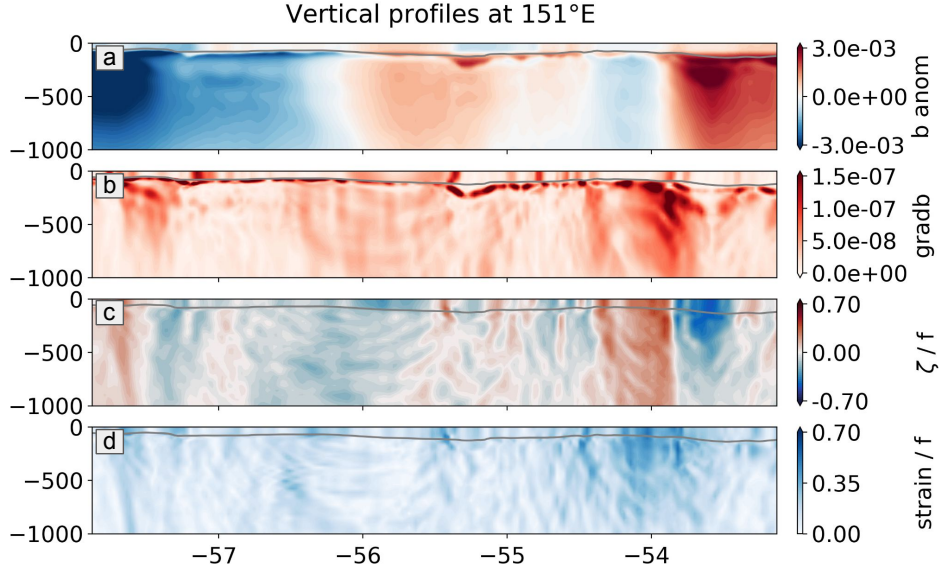
The interaction of multiple mesoscale and submesoscale eddies contributes to the formation of strong horizontal velocity gradients. Figure 2 shows COAS modelled fields at a selected daily snapshot on March 29th 2020 (the October fields show very similar results and are therefore not shown) for (a) relative vorticity ( $\zeta = v_x - u_y$  with the subscript  $x$  of  $y$  referring to the derivative of the variable along the  $x,y$  direction), (b) the strain rate ( $\sigma = \sqrt{\sigma_n^2 + \sigma_s^2}$  with  $\sigma_n = u_x - v_y$  the normal strain and  $\sigma_s = v_x + u_y$  the shear strain rate), at 250 m depth and their distribution at different depths. These diagnostics give insight into the dominant dynamics in the region, and how they extend in the water column. The sign of the relative vorticity in Figure 2c is used to separate cyclones (negative) and anticyclones (positive) (Frenger et al., 2015). Filaments of vorticity have a larger positive skewness than vortices due to ageostrophic frontogenesis (Siegelman et al., 2020). The COAS fields show positive skewness indicating strong vorticity filaments down to 1000m for both seasons (skewness is quite similar at all depths in March, 0.39, 0.41, 0.47, 0.37 from the shallowest to the deepest, slightly decreasing in October, 0.66, 0.66, 0.72, 0.57). Regions of strong strain in Figure 2d are of interest as they lead to the formation of strong horizontal gradients of buoyancy, or fronts, which cause the

development of vertical velocities (Hua & Klein, 1998). Both the vorticity and the strain rate show strong filaments and fronts with a magnitude close to the Coriolis frequency  $f$ , especially in the first 100 m, highlighting the energetic variability of this region, consistent with the strong EKE shown in Figure 1. In October the distribution shows high values of vorticity and strain in the upper ocean, especially within the October average mixed layer of 250 m. Regions of strong strain at the edge of the eddies are also regions where strong horizontal buoyancy gradients develop (Figure 3) along the full water column.



**Figure 2.** (a) Normalised relative vorticity, (b) strain rate and (c) and (d) are respectively the relative vorticity and strain distributions at 45 m, 88 m, 251 m, and 540 m for hourly snapshots over five days in March (25th to 29th) and October (25th to 29th). The grey dashed line represents the longitude at which the vertical profiles are taken throughout the paper

Vertical sections from the model reveal complex dynamical features in the ocean interior. Figure 3 shows the vertical extension of mesoscale eddies below the mixed layer (ML). Here, the mixed layer depth (MLD) is chosen where the density increases of  $0.03 \text{ kg m}^{-3}$  with respect to the density at 7 m depth (de Boyer Montégut et al., 2004). The average MLD in the region for the selected dates is 103 m in March (late summer/early autumn) and 250 m in October (late winter/early spring). We show a transect at a longitude of  $151^\circ$  E down to a depth of 1000 m because it crosses specific eddy structures of this region (see dashed line in Figure 2a, b), with different characteristics. A large eddy develops from the surface down to 1000 m south of  $54^\circ$  S, and has a strong positive signature in vorticity and a strong horizontal buoyancy gradient. Smaller eddies north of  $54^\circ$  S or south of  $55^\circ$  S are trapped in the subsurface and develop down to 500 m and 300 m. They have a strong signature of buoyancy anomaly in their core and a strong horizontal buoyancy gradient at their edge. This vertical section also shows internal wave structures at depth, which will not be reconstructed via the sQG methodology.

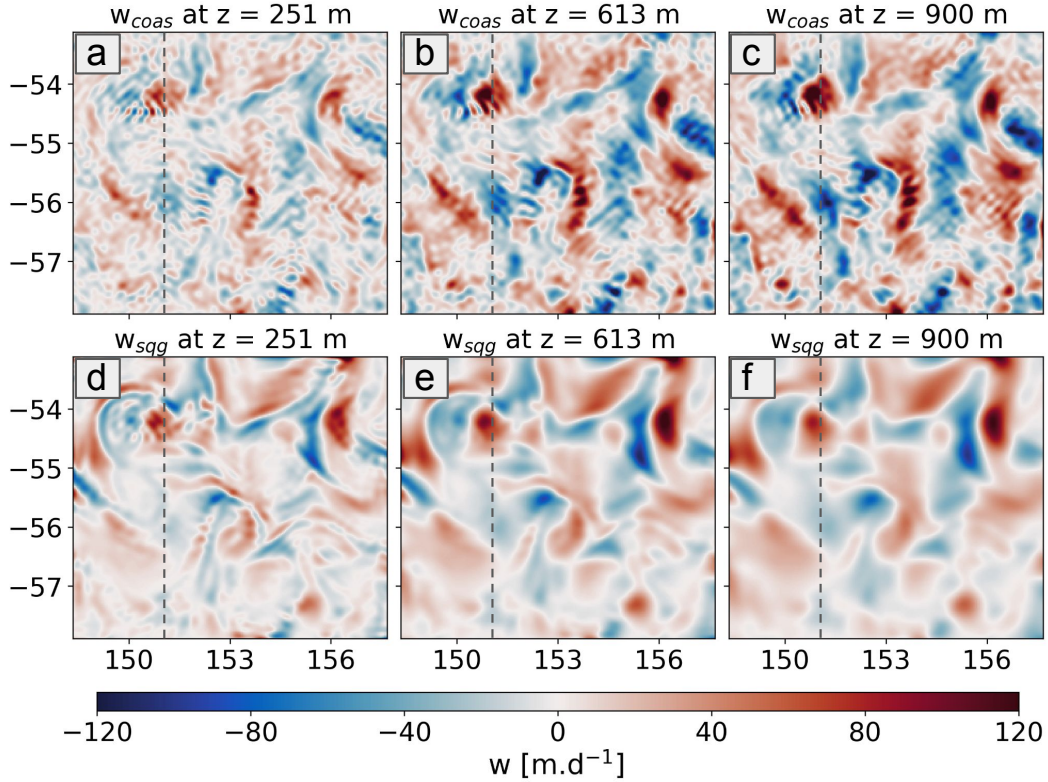


**Figure 3.** Vertical profiles on the first 1000 m at 151°E of (a) buoyancy anomaly, (b) horizontal buoyancy gradient, (c) normalized vorticity, (d) normalized strain rate. Results refer to one hourly snapshot at 00UTC on the 29th of March. The grey contour is the average mixed layer in the region over 5 days

## 4.2 Vertical velocity reconstruction

Once the stratification and amplitude parameters are optimised, the sQG methodology is applied to reconstruct vertical velocities in the domain for both seasons, using daily averaged COAS SSH. The physical sQG fields in Figure 4 show that the small eddy structures are present below the mixed layer at 250 m, and become progressively smoothed at depth. The 2D geographical fields of the COAS and sQG vertical velocity at three different depths are represented in Figure 4 a) to f). sQG can reconstruct some of the small-scale eddy structures closer to the surface, see Figure 4d at 251m between 54°S and 55°S for instance, but these features are gradually smoothed out with depth in the sQG reconstruction, in Figures 4 e and f. sQG does not reconstruct the small wave-like instabilities in the COAS simulation at each depth. The correlation remains consistently around 0.6 at all depths below the mixed layer, representing the mesoscale and large submesoscale filaments. The reduced energy in the reconstructed vertical velocities at depth is also quantified by the vertical profiles of rms in Figure 6d.

Vertical sections of the COAS versus sQG  $w$  fields are shown in Figure 5 b and c, and highlight one of the key aspects of this region – the vertical velocities in COAS are stronger at depth, in this regions of strong bathymetric control downstream of a major ridge, as noted by Tamsitt et al. (2017). Above 500 m, the sQG shows some of the small-scale variability that is gradually filtered at depth. The COAS features at depth are large-scale structures with a large SSH signature. In fact, the vertical profile between 53°S and 54°S shows that while the COAS simulation has an inversion of velocity around 500 m, the sQG projects the surface dynamics to the ocean interior. South of 54°S there is a large ascending  $w$  hotspot colocalized with a large horizontal buoyancy gradient and vorticity (Figure 3 b and c). This feature is captured by the SSH and projected by the sQG exponentially at depth, since the sQG, by construction, is not able to represent the increasing amplitude at depth. This is highlighted by the difference between the COAS and sQG fields (Figure 5d). In fact, in this specific transect, the SQG explains up to 90%

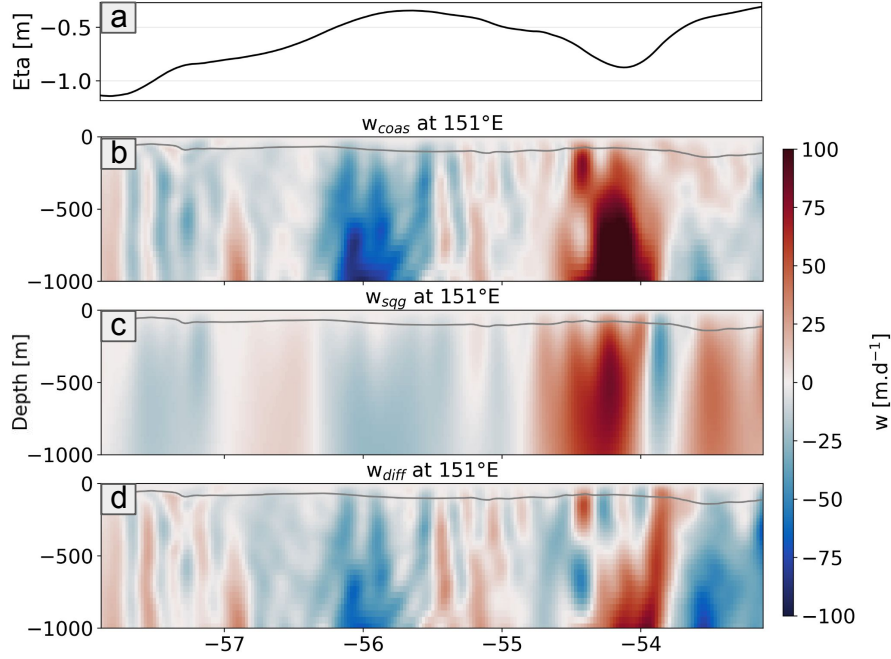


**Figure 4.** Vertical velocity from the COAS simulation for depths of (a) 250m, (b) 613 m and (c) 900 m. Vertical velocity sQG reconstruction at (d) 250m, (e) 613 m and (f) 900 m. The grey dashed line in (a) represents the longitude at which the vertical profiles are taken in Figure 3. Figures a to f refer to the daily averaged outputs on the 29th of March 2020

of the signal in the first 500 m gradually decreasing to 50% at 1000 m. The case at  $56^\circ\text{S}$  is different, where we find a negative (downwards)  $w$  hotspot whose sign is captured by sQG, whereas the amplitude is underestimated. This corresponds to a region where the SSH gradients (Figure 5a) are weaker, which explains the magnitude misfit in the reconstructed  $w$ , especially since it develops from the bottom. The difference between the two fields is in the order of half of the COAS  $w$  value, underlining a limitation of the sQG model. South of  $57^\circ\text{S}$  the main difference between the two fields is due to some rapidly changing (20 - 30 km) vertical velocity structures that cannot be effectively represented by the sQG reconstruction because of their small scale. The specific processes involved in the  $w$  intensification at depth observed in the model are not studied in detail, the contribution could come from the enhanced dynamics downstream of the Southeastern Indian ridge, wave effects, or different instabilities from the flow fields. What stands out from Figure 4 and Figure 5 is that if these instabilities in the ocean interior have a signature in SSH (e.g. south of  $54^\circ\text{S}$ ), they are picked up and projected at depth by the sQG, with the correct intensity for the first 500 m and underestimated up to 50% below. The effective sQG cannot reconstruct other dynamical features and a more sophisticated representation of the ocean interior would be necessary for this purpose. This means that in such an energetic region in the presence of these dynamics, we might be missing part of the transport of heat between the ocean interior and the surface.

The spatial scales being resolved are quantified by the wavenumber spectra for the COAS and sQG reconstructed  $w$  fields at two different depths in March, shown in Fig-



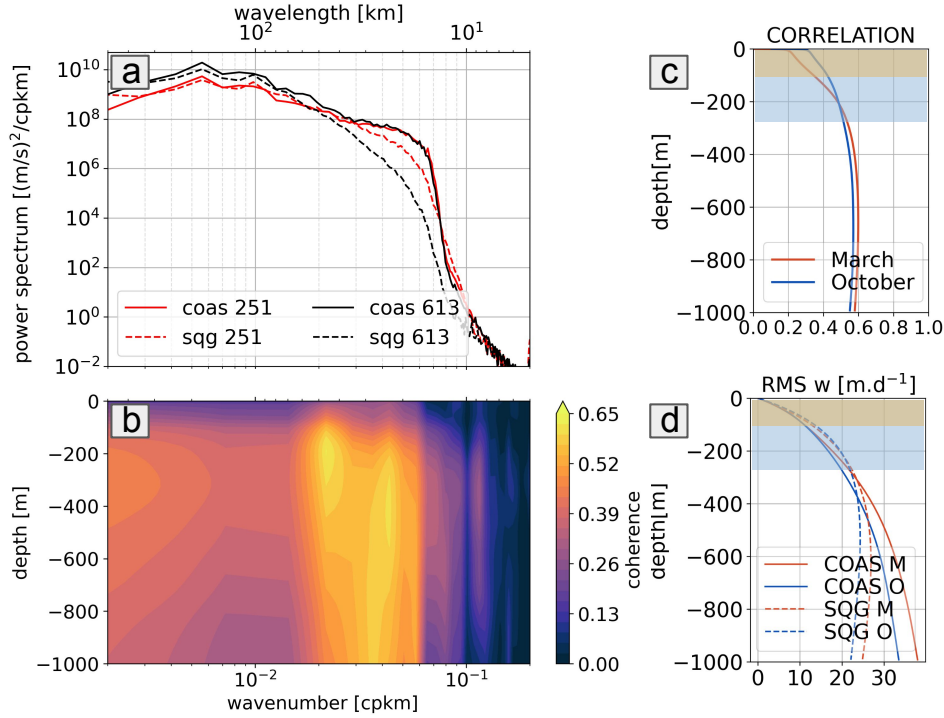


**Figure 5.** (a) SSH profile at 151°E. Vertical profile of (b) COAS, (c) sQG vertical velocity and (d) their difference at 151°E, where the x-axis is the latitude [°S]. The grey contour corresponds to the average mixed layer (ML) in the region for the five days between March 25th and 29th

ure 6a, characterised by a relatively good spectral energy of the two fields at large scales at these two depths. At 251 m depth, the sQG  $w$  clearly shows a reduction in energy at scales  $<30$  km; at 613 m depth, the sQG starts losing energy below 70 km, underestimating smaller scales.

Figure 6b shows spectral coherence for each depth. Although scales larger than 70 km have similar spectral energy (Figure 6a), they are less coherent over the full depth range. As shown in Figure 5, the large scale sQG  $w$  energy is projected from the surface downward, whereas the COAS  $w$  is often depth-intensified. Although the sQG  $w$  has reduced spectral energy at scales smaller than 50-70 km, it has the highest  $w$  spectral coherence. The spectral coherence remains above 0.5 for scales between 30 km and 60 km for the full water column, slightly decreasing with depth. This is in line with the 2D fields of Figure 4 which show that the sQG field at depth is smoothed compared to 250 m, but some of the large mesoscale structures ( $\sim 50$  km) are resolved with a lower magnitude (see the rms at depth Figure 6d).

The sQG reconstruction is evaluated in terms of the spatial correlation between the COAS and sQG  $w$  field, and their associated rms. Figure 6c shows that spatially the structures are well localized, with a spatial correlation reaching 0.6 for the full water column, both in March and October, in line with other sQG reconstructions in the Pacific Ocean (Qiu et al., 2020). At the surface the result is not exploitable down to the bottom of the ML, indicated by the shaded grey sections (darker grey for the shallower ML in March, lighter grey for the average ML in October). The rms (Figure 6d) is comparable between sQG and COAS in the top 300 m (400 m) in March (October), although the sQG has slightly higher values than COAS in the first 200 m (300 m). Below 400 m, the sQG rms slightly decreases but maintains  $w$  rms values of around  $25 \text{ m} \cdot \text{d}^{-1}$  in both seasons, which is lower than COAS by  $15 \text{ m} \cdot \text{d}^{-1}$  and  $10 \text{ m} \cdot \text{d}^{-1}$  respectively in March and



**Figure 6.** (a) Wavenumber spectrum of COAS (filtered at 20 km) and sQG vertical velocity at 251 m and 613 m on March 29th. (b) spectral coherence between the COAS and sQG field for each depth. (c) spatial correlation between the COAS and sQG vertical velocity fields for each depth. (d) rms of the COAS and sQG vertical velocity fields. The shading corresponds to the average mixed layer (ML) in the region for the five days between March 25th and 29th (orange), and between October 25th and 29th (blue).

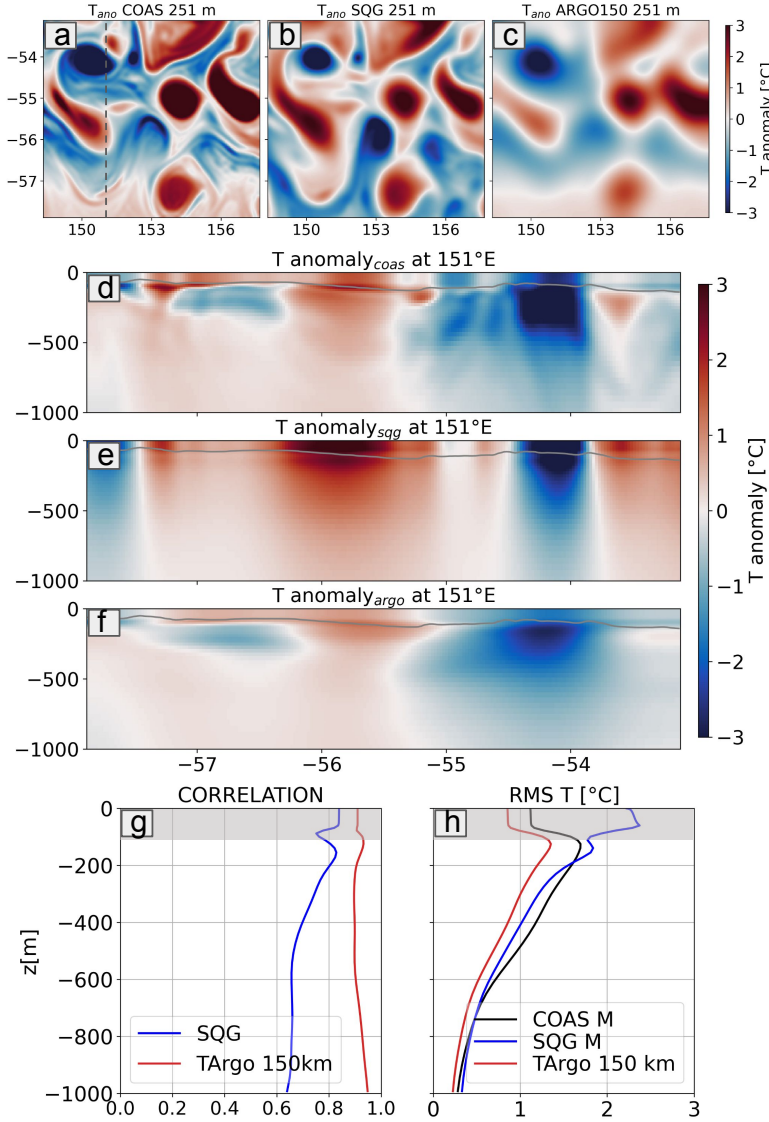
October at 1000 m. This means that on the overall field, the SQG reproduces up to 70% of the COAS rms at depth. Note that, as in Klein et al. (2009), the sQG diagnosed vertical velocity rms does not present any local maximum close to the bottom of the ML. This feature is also absent in the filtered version of the COAS vertical velocity rms, but present in the non-filtered version (not shown). This is likely to be due to the sQG lack of representing ML dynamics and Ekman pumping at the bottom of the ML.

### 4.3 Temperature anomalies reconstruction

We consider two techniques for estimating fine-scale temperature fields in order to calculate vertical heat transport (VHT). The first is an "Argo-like" pseudo-gridded product for the temperature anomaly, and the second is the temperature anomaly derived from the sQG density anomalies by linear regression (see section 3.3). Both temperature anomalies are illustrated in Figure 7 compared with the full-model COAS temperature anomaly. An example of the 2D geographical fields at 250 m depth (Figure 7 a to c) shows that the small-scale structures in the temperature anomaly are smoother in the Argo-like case but with similar polarity, whereas the sQG temperature reconstruction has larger amplitude anomalies, but sometimes of the wrong sign. The diagnostics over the entire region and at depth confirm this. The spatial correlation of the full COAS temperature anomaly with the reconstructed temperature fields (Figure 7g) indicates that these large-scale structures in the Argo-like field are well positioned (correlation  $\sim 0.9$  along the full water column). Whereas the sQG temperature reconstruction has a lower



correlation (0.8 at the surface decreasing to 0.6 at depth) since specific structures and vertical water mass distributions (see Figure 7 e and f ad 57°E compared to d) are not reconstructed by the sQG theory.



**Figure 7.** Temperature anomaly for the 2D field in the (a) full model, (b) sQG reconstruction, (c) Argo-like reconstructed field. d to f vertical profile at 151°E respectively for the same products. (g) spatial correlation between the full COAS simulation and the sQG reconstruction (blue) and the Argo-like field (red). (h) rms of the three fields at each depth. The grey shading is the average mixed layer depth in March

The full model rms (Figure 7h) is closer to the sQG case than the smoother Argo-like reconstruction below the mixed layer. The sQG rms (h) below the ML presents a good agreement with the COAS temperature, with a slight underestimation of 15% in March up to 600 m and 20% in October up to 800 m (not shown) with respect to the reference anomalies. Below, the temperature anomalies are well aligned. The Argo-like temperature differs from the sQG reconstruction by an extra 15% in rms between 200

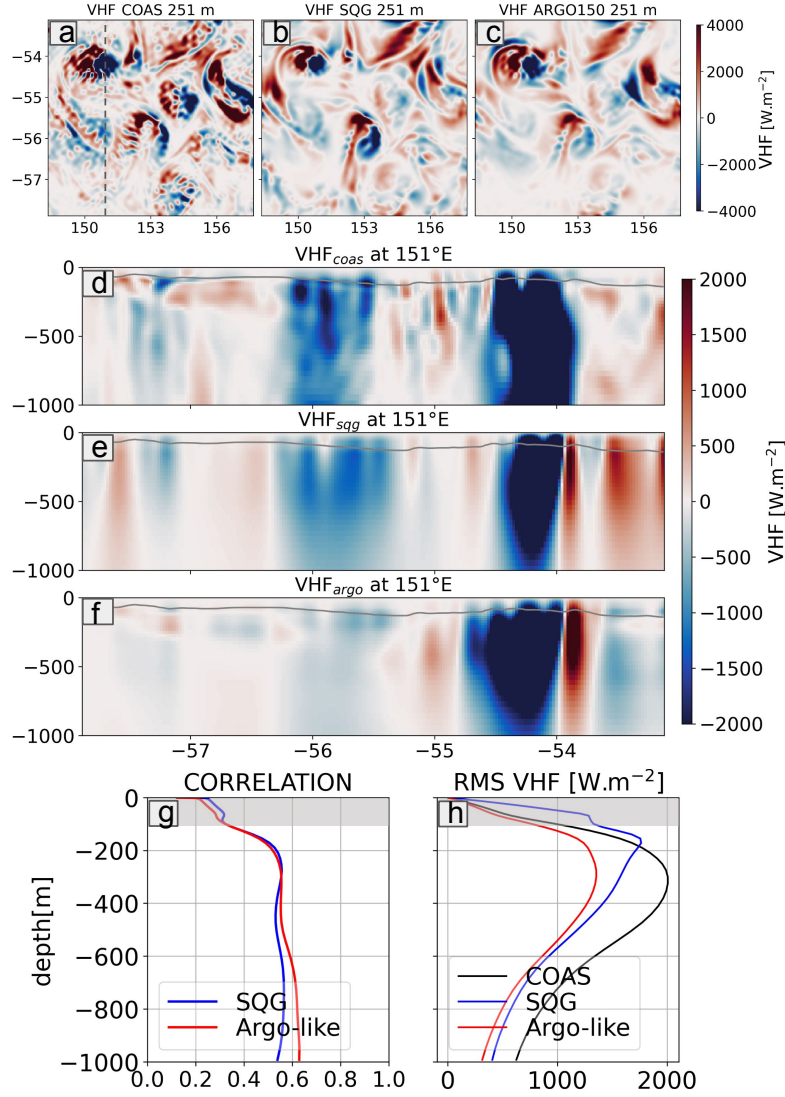
and 600 m and up to 30% at 1000m. Thus, even if the large-scale structures are spatially better represented in the Argo-like field, the small scales ( $<150$  km) missing are bringing much of the variability in the temperature anomaly. Different studies using ARMOR 3D have pointed out the underestimation of inferred QG vertical velocities (Pascual et al., 2015; Barceló-Llull et al., 2016; Mason et al., 2017), due to the smoothing of the field. Within the ML, the sQG temperature anomaly rms presents a large departure from the modelled one ( $\sim 100\%$  at the surface), which would prevent any interpretation of the results in this layer. The Argo-like temperature anomalies are closer to the correct ML structure, with a lower magnitude, overall between 25 and 30% lower than COAS along the full water column.

The comparison of Figure 7 e and f shows one of the limits of the sQG methodology for reconstructing temperature anomalies: while the large-scale structures are well represented and aligned with the modelled ones, sQG is not able to represent all the water mass structure in the mixed layer and at the bottom of the mixed layer. Between  $56^\circ\text{S}$  and  $57^\circ\text{S}$ , for instance, the negative temperature anomaly is not reproduced, and the full sQG water column has a positive temperature anomaly. In some cases, the sQG anomaly has an opposite sign next to the boundaries (e.g. south of  $57^\circ\text{S}$  because it projects the surface negative anomaly at depth). The Argo-like product (f) keeps the large-scale temperature variability along the water column and the distinct water masses that exist north and south of the polar front.

#### 4.4 Vertical Heat Fluxes reconstruction

After reconstructing the vertical velocities and temperature anomalies in the first 1000 m of the water column, we can estimate the amount of vertical heat flux that this methodology can capture compared to the COAS simulation taken as a reference for the study. Figure 8 shows the full COAS model VHF compared to the two reconstructed VHF, both computed with sQG vertical velocity but either using the Argo-like temperature or the sQG reconstructed temperature. The 2D fields (Figure 8 a to c) show that the two reconstruction methodologies for computing the VHF give similar results in terms of reconstructed structures, highlighting that the sQG vertical velocities are the key factor for the VHF. As for the vertical velocities, the small-scale COAS wave patterns are not reconstructed with sQG.

The vertical sections along  $151^\circ\text{E}$  (Figure 8 d to f) show a large eddy strain driving downwelling around  $54^\circ\text{S}$  in COAS, that is well captured by the two reconstruction methodologies. In contrast, north of  $54^\circ\text{S}$ , COAS shows a small sub-surface-trapped eddy inducing upwelling in the COAS VHF that is not captured by either of the correlations. The COAS temperature (Figure 7) shows this eddy below the mixed layer extends down to 400 m, but Argo-like resolution is too coarse to retain it, and the sQG temperature reconstruction projects it from the surface to depth. These two examples highlight that the reconstructed VHF only presents the major coherent structures given by the  $w$  sQG on the full vertical column. Although, in regions south of the polar front like between  $56$  and  $57^\circ\text{S}$  the Argo-like temperature reconstruction allows for a representation of some vertical structures in the VHF. The correlation (Figure 8g) shows a similarity between the sQG and the Argo-like reconstructed VHF down to a depth of 700 m. Between 700 m and 1000 m, the sQG  $w$  field represents large spatial scales, but the sQG temperature reconstruction has reduced correlation (Figure 7g), resulting in a reduced sQG reconstructed VHF correlation at depth ( $< 0.1$  less than the Argo-like reconstruction) at 1000 m. However, due to the smoother reconstructed temperature, the Argo-like case has a lower rms (Figure 8h, around 40% at 200m and around 10% at depth). North and south of the meandering polar front region with its energetic, deep-reaching eddies, the Argo-like temperature anomaly does a better job representing the vertical water masses layers in the upper water column, and the dynamics within and below the ML are better represented than with the sQG reconstruction. However, Figure 8h indicates that over-



**Figure 8.** Vertical heat flux for the 2D field in the (a) full model, (b) sQG reconstruction, (c) Argo-like reconstructed field. (d) spatial correlation between the full COAS simulation and the sQG reconstruction (blue) and the Argo-like field (red). (e) rms of the three fields at each depth. The grey shading is the average mixed layer depth in March

all on the vertical heat fluxes, the rms is dominated by the contribution of the vertical velocity; recall that the sQG reconstructed vertical velocity underestimates the full COAS model below 600 m by 30-40% (Figure 6d). Thus, improving the  $w$  reconstruction is paramount for a better representation of the VHF.

Thus, the reconstruction using the sQG  $w$  and temperature and the one using the sQG  $w$  and the Argo-like product give comparable statistical results in terms of VHF structures, albeit with improved correlation of the vertical structures with the "Argo-like" temperature reconstruction over all depths, but with an overall better rms below the mixed layer provided by the full ( $w$  and temperature) sQG reconstruction. These results suggest that the existing operational gridded temperature products like ARMOR 3D could be used to retrieve VHF associated with mesoscale structures down to 1000

m and offer promising results, especially within the ML where the sQG performs worse. This methodology is a good alternative to the sQG reconstruction, especially in regions where a tight linear regression between density and temperature necessary for the sQG reconstruction cannot be correctly estimated.

## 5 Discussion

This work aims to investigate the use of effective sQG for reconstructing vertical velocities and vertical heat fluxes in an energetic region south of Tasmania. Strong mesoscale strain and relative vorticity in this region make it a hotspot for the generation of sub-mesoscale fronts, leading to intense vertical velocities and enhanced vertical fluxes of heat and other tracers. Given the availability of five days of data in two contrasting times of the year (March and October), the aim is to demonstrate what physical structures the effective sQG is able to reconstruct at these specific times without focusing on the full seasonal cycle.

For the effective sQG reconstruction, we need to tune two parameters. We use a constant value for the vertical stratification parameter and optimize it iteratively against the relative vorticity reconstruction below the mixed layer, finding a small difference between March and October. A different choice could be to take climatological values for the vertical stratification or to compute an average value from existing ARGO floats. Indeed, for a longer time series over multiple years/seasons, this would need to be tested to set a seasonally varying vertical stratification parameter. For the parameter  $c$ , the tuning is performed on the 2D physical fields. In this particular region, non-QG dynamics in the model can impact the stability of this parametrization. More analysis is needed in less energetic regions to test the robustness of this methodology.

After parametrizing the sQG coefficients, this study shows that  $w$  in this region can be reconstructed in the first 1000 m of the water column with a spatial correlation of 0.6 against the full COAS model  $w$  field. The main result of this study is that the reconstructed field captures most of the variance along the water column (fully at the surface, gradually decreasing to 70% at 1000 m depth over the full domain). The small-scale reconstructed structures have the highest spectral coherence at scales between 20 km and 70 km, with coherent scales increasing from below the mixed layer to larger scales at depth. Here, we interpret our effective sQG results only below the mixed layer since we find, as in previous studies (Klein et al., 2009; Qiu et al., 2020), that the sQG methodology cannot fully represent the mixed layer dynamics. This is, in fact, a first estimation of the feasibility of using the effective sQG in this energetic region. It could be expanded to improve the reconstruction both at depth and within the mixed layer, using improved theories.

Within the mixed layer, other parametrisations accounting for improved dynamics representation have been proposed in the literature. Ponte et al. (2013) demonstrate the impact of wind-driven mixed layer dynamics on vertical velocities in the case of small wind-driven velocities compared to those associated with eddies and an Ekman number not small. Callies et al. (2016) highlight the importance of mixed-layer instabilities compared to mesoscale-driven surface frontogenesis. Chavanne and Klein (2016) include diabatic vertical velocities in the surface layer based on thermal wind balance restoration after perturbations due to turbulent vertical mixing of momentum and buoyancy. These could also be explored in this Southern Ocean region in the future.

Below the mixed layer, the major components contributing to the 30% missing rms in the sQG reconstruction are the depth-intensified vertical velocities, the smaller 20-30 km barotropic structures present on the full water column, and the small-scale wave structures at depth highlighted in Figure 3. The analysis over five days in two seasons does not allow a quantitative analysis of the specific dynamics that these features rep-

resent. However, the strong deep vertical velocities could be due to the interaction of the flow with the bathymetry at the level of the Southeast Indian ridge. A barotropic component of the flow downstream could cause the intensification of the vertical velocities, which the effective sQG would not reconstruct. In addition to barotropic structures, sub-surface intensified eddies, as explained in Section 3.2, are not directly reconstructed by the effective sQG even if they present a SSH signature. Assassi et al. (2016) have identified two indexes that define the type of the eddy (surface, subsurface) based on SSH and surface density/SST information, that could be used to gain prior knowledge on the type of surface or sub-surface structures that may be reconstructed by sQG in a region, given a field of SSH and SST. An alternative approach to capture more of the depth-intensified  $w$  could be to use the extended interior and surface quasigeostrophic (isQG) theory from Wang et al. (2013) and Liu et al. (2019), which yields good results in other regions (e.g., South Atlantic Ocean). Miracca-Lage et al. (2022) found a strong dependency of the results on the mixed layer depth and improved results with the isQG compared to the effective sQG. Nevertheless, the 30% degradation of the reconstructed results at depth remains an issue.

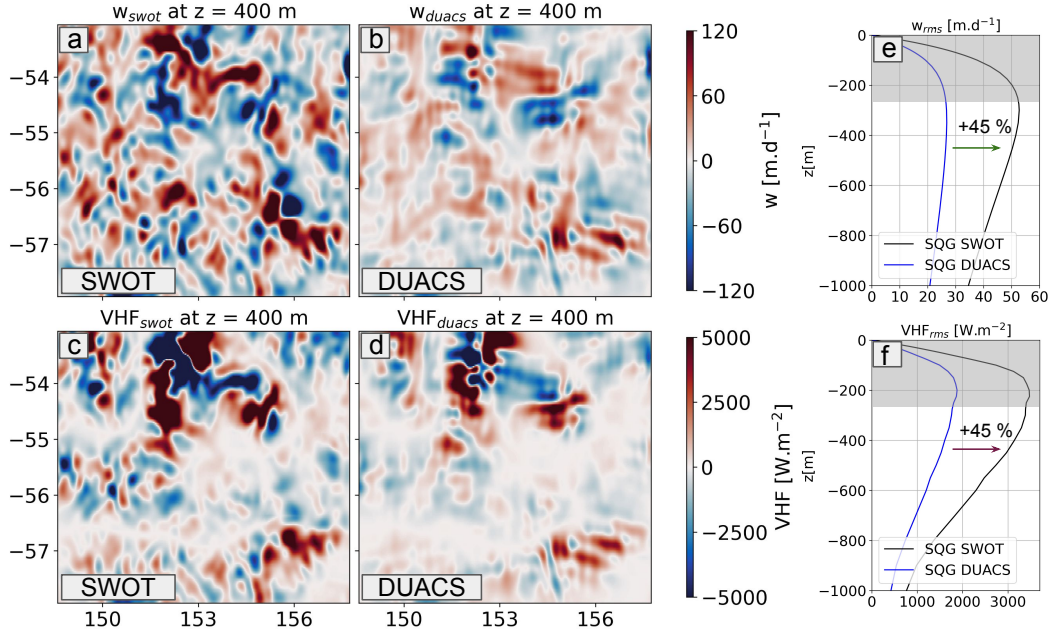
To estimate vertical heat fluxes, having temperature anomalies at the same resolution as the  $w$  field is still challenging. In particular, reconstructing the temperature anomalies with sQG density anomalies presents some limitations in this region of distinct dynamics and vertical water mass layers north and south of the polar front. Specifically, sQG density anomalies do not capture the temperature anomaly inversions along the water column and coherently project the SSH gradients at depth, changing the real distribution of water masses, especially south of the polar front. The sQG reconstruction of temperature anomalies in the domain also requires a specific linear regression between density and temperature tuned to the area for each season. This is a limiting requirement for using this methodology more widely to compute vertical heat fluxes. To overcome this limitation, we have tested using a pseudo-gridded temperature product with the same resolution in our region as the operational ARMOR 3D product (Guinehut et al., 2012), which smooths out the small scales ( $\leq 150$  km) compared to the sQG reconstruction, but is spatially better correlated to the reference temperature from the model and captures the mixed layer variability. The spatial scales represented by ARMOR 3D are limited by the present generation of 2D gridded altimetry maps. In the future, we can expect higher-resolution altimetry maps including SWOT, that can be used to stir and reconstruct the ARMOR 3D temperature and salinity fields, and thus have similar high resolution of 2D surface SSH and 3D T/S fields. Other methodologies exist for reconstructing temperature at depth from SSH, such as the gravest empirical mode (GEM) used by Swart et al. (2010) south of Africa or Meunier et al. (2022) in the Gulf of Mexico, or new machine learning methodologies (Champenois & Sapsis, 2024), which could also be explored in future studies.

VHF are computed for both temperature datasets using the same sQG  $w$  fields. Results indicate that the vertical velocity has a dominant role in the reconstruction. In fact, we obtain comparable spatial correlations and rms with the two methods below the mixed layer, even though the two temperature products have quite different metrics, especially in terms of spatial correlation. Improving the  $w$  correlation to more than 0.6 would yield to better results for the VHF computed with the Argo-like temperature product. The Argo-like product offers improved results within the mixed layer in terms of rms. The downside of using such product is that the 150 km resolution hides small mesoscale and submesoscale dynamics. Zhang et al. (2023) found that the submesoscales ( $< 40$  km) amplify the VHF via the inverse cascade, which explains the underestimation of the VHF in particular with the Argo-like temperature. We have examined two contrasting seasonal examples for this first demonstration of VHF derived from SSH data alone. Although March and October show comparable results, the limited data available (5 days for each month) do not allow for a thorough seasonal comparison.



This study was performed as a preliminary analysis of vertical velocity and VHF reconstruction that may be possible with the fine-scale SSH available from the SWOT mission. Comparing sQG scales with SWOT observability in the area (30-40 km in SSH from the first SWOT KaRIn data) tells us that the main sQG-observable dynamics (highest spectral coherence for velocity between 30 and 70 km) are within the SWOT data observing capabilities, which is encouraging for developing this study even further with real data.

As a first test of the effective sQG reconstruction with SWOT data, we have performed a simple spline filtering to combine 5 consecutive and neighbouring SWOT swaths in our region. These are preliminary beta pre-validated SWOT L3 data available from the 21-day science phase of SWOT in October 2023 (see Appendix B and Figure B1). The first test case of applying effective sQG to these SWOT 2D SSH data is shown in Figure 9, in comparison with the standard 2D SSH fields from DUACS/Copernicus. These preliminary results indicate that the spatial patterns of reconstructed  $w$  are similar to those derived from DUACS/Copernicus maps. The sQG reconstructed vertical velocities have double the rms variability compared to those derived from DUACS  $w$  structures, introduced via the 2D observations and small-scale features in the SWOT SSH. Appendix Appendix B shows how the 2D SSH field was created from the swaths, and more detailed analysis and validation will be shown in a subsequent study. The VHF was computed with the real ARMOR 3D dataset ([https://data.marine.copernicus.eu/product/MULTIOBS\\_GLO\\_PHY\\_TSUV\\_3D\\_MYNRT\\_015\\_012/description](https://data.marine.copernicus.eu/product/MULTIOBS_GLO_PHY_TSUV_3D_MYNRT_015_012/description)) for the same period and region of the SWOT data. We note that these early SWOT data still need to be validated, especially to understand the signal and noise in the new scales < 150 km in wavelength.



**Figure 9.** (a - b) sQG  $w$  and (c - d) VHF computed with the sQG  $w$  and the ready ARMOR 3D product for the week of the 24th of October 2023 from the real SWOT SSH data (a - c) and from the DUACS SSHm (b - d). Rms of (e)  $w$  and (f) VHF for the SWOT and DUACS fields. The grey shading is the average mixed layer depth in October



The potential of this effective sQG reconstruction in the Southern Ocean around 50°S is that every five days during the 3-years of the SWOT science mission, it may be possible to estimate a high-resolution 2D map of the vertical velocities and vertical heat fluxes. This will first require careful validation of the SWOT data and against *in-situ* ground-truth data of 3D velocities and temperature at depth. This methodology could then be extended to larger domains in the energetic Southern Ocean to evaluate the vertical exchanges between the surface and the ocean interior on a regular basis.

Our study has demonstrated that effective sQG can reconstruct up to 70% of the modelled vertical velocities larger than 30-40 km in wavelength in this region and 60-70% of the vertical heat fluxes at these scales. There are limitations in the sQG vertical projection from the surface to depth, which cannot capture the vertical velocities generated at depth downstream of the Southeast East Indian Ridge and cannot capture the small-scale  $w$  ( $< 40\text{km}$ ), nor from internal waves. Another limitation is that fast-evolving structures would still be missed because of the SWOT 21-day coverage (see Appendix Appendix B for SWOT coverage over the region). Indeed, in our simple interpolation, we take the central date as a reference and assume that the observed structures in neighbouring passes offset by 1-day are static over each 5-day period. This assumption is similar to that used with most ship-borne campaigns in reconstructing the 3D field over many days. CNES and Collecte Localisation Satellites (CLS) are working on a higher-resolution 2D multi-altimeter SSH operational product like DUACS/Copernicus that will include SWOT. This product will include more small scales compared to the current Copernicus operational product, but any mapping procedure will still smooth some of the small scales observed within the swaths. Ideally, some regional high-resolution mapping may be developed in the future to reduce this smoothing. These improved future operational 2D gridded SSH data maps that include SWOT data will also help improve derived products like ARMOR 3D due to, for example, better localisation of eddies. This will increase the spatial resolution and potentially avoid the temperature and VHF underestimation of the vertical water column.

## Appendix A sQG reconstruction of temperature anomalies

Temperature anomalies can be reconstructed with only altimetry data, using sQG. Two preliminary analyses have been carried out to ensure that this anomaly can be retrieved from the sQG density anomalies.

### A1 Stability at the base of the mixed layer

The ocean density is a function of the 3D structure of the water temperature, salinity and pressure and depends strongly on the temperature and salinity structure. The Brunt-Väisälä frequency coefficient  $N$ , a function of the density gradient, dictates the stability of the water column. It is defined as

$$N^2 = \frac{g}{\rho} \frac{\partial \rho}{\partial z} \quad (\text{A1})$$

where  $g$  is the gravitational constant,  $\rho$  is the density and  $z$  is the water depth. The necessary condition for water column stability is a positive  $N$ . This stability can be split into the contribution of the thermal and haline parts

$$\frac{1}{\rho} \frac{\partial \rho}{\partial z} = -\alpha \frac{\partial T}{\partial z} + \beta \frac{\partial S}{\partial z} \quad (\text{A2})$$

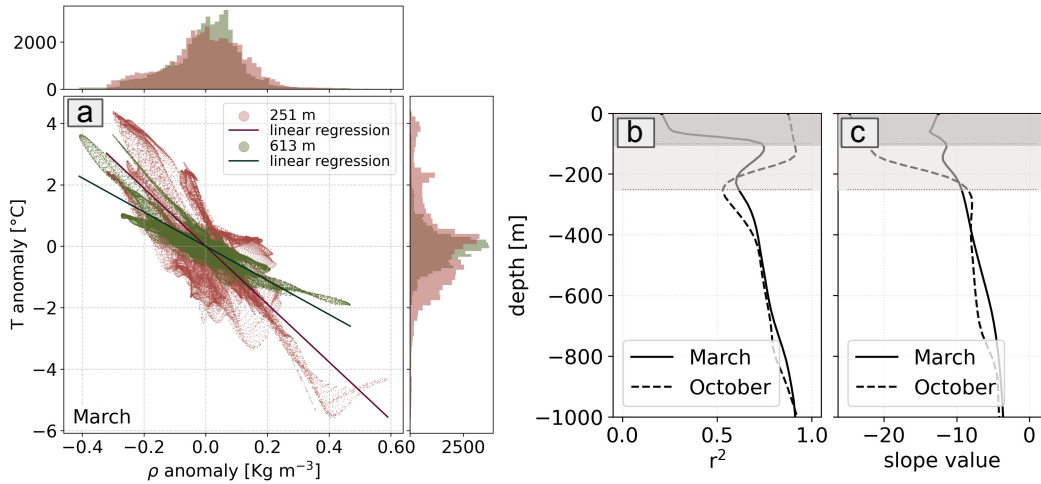
where  $\alpha = \frac{1}{\rho} \frac{\partial \rho}{\partial T}$  and  $\beta = \frac{1}{\rho} \frac{\partial \rho}{\partial S}$ . The two terms representing the role of the thermal and haline stratification on the stability at the base of the mixed layer are the thermal stability  $N_T = -g\alpha \frac{\partial T}{\partial z}$  and the haline stability  $N_S = g\beta \frac{\partial S}{\partial z}$ . Where thermal stability dominates, the density gradient is driven by the temperature, so we can suppose that the temperature measurements in the region are sufficient to retrieve the density. The opposite

is true when the haline stability dominates. We first verified that in this polar front region, the thermal stability dominated over the haline stability at all levels (this is not the case in the sea-ice zone further south).

Then, by building scatter plots between the density and temperature anomalies, we can investigate the relationship between those two variables as a function of depth and season.

## A2 Linear correlation between density and temperature anomalies

The linear regression is built for each vertical layer from the COAS density and temperature spatial anomalies (taken with respect to the mean of the 2D detrended spatial field). Figure A1 shows two examples of the scatter and linear regression below the mixed layer, where the regression improves with depth. Figure A1 c provides the values of the linear regression slope at each depth. Below the mixed layer, these values are similar in March and October and are consistent during the five days. The intercept value is close to zero in all cases. At 251 m, the temperature values still present some extreme values that are overestimated or underestimated by the linear regression. At 613 m, the anomalies are more concentrated around zero (see the distribution of the variables), the scatter itself presents a more linear shape, and there are fewer points with extreme values. The density and temperature have an increasingly linear relation with depth, with fewer points of temperature extremes departing from it. This is shown in Figure A1b, where the significance coefficient  $r^2$  is close to 0.5 at the bottom of the ML for both seasons, gradually increasing up to 0.9 at depth. The p-value of all regressions is below 0.05, a threshold usually taken to test the null hypothesis (not shown). This indicates the statistical significance of the regression. Once the fine-resolution density anomalies are reconstructed with sQG, the temperature anomalies are deduced by implementing the linear regression mentioned above for each depth.

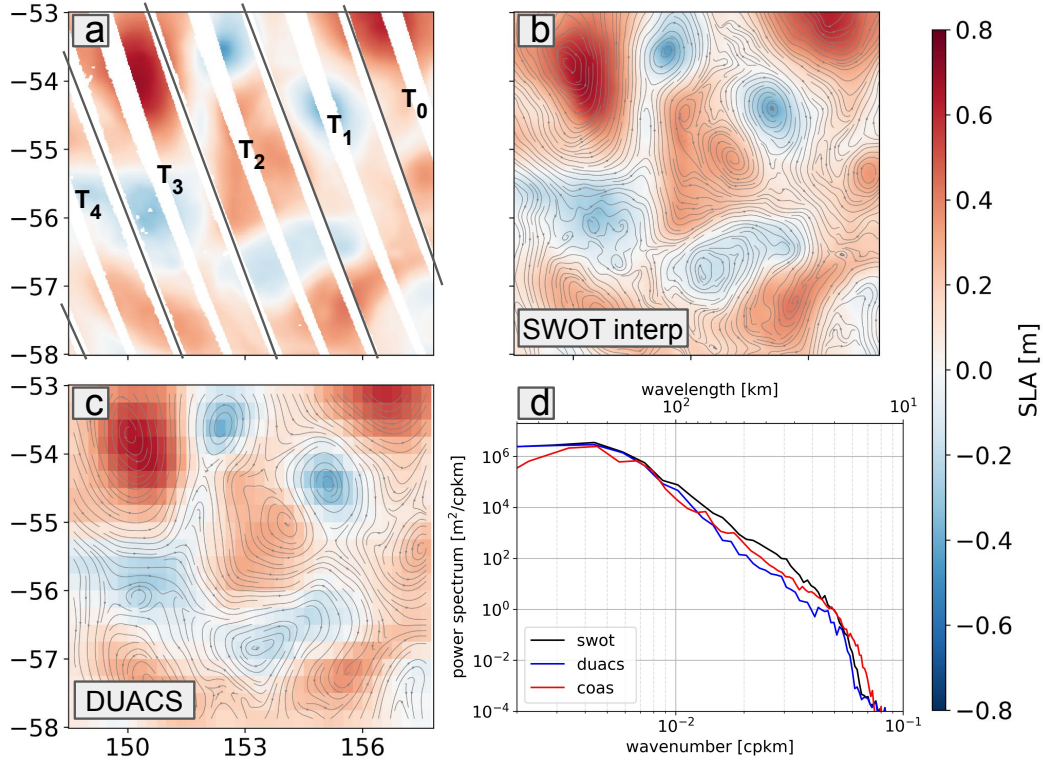


**Figure A1.** (a) Scatter plot between density anomalies and temperature anomalies for COAS at 251 m and 613 m, with the respective linear regression used to reconstruct the temperature anomalies. The distribution of each variable is also shown in the lateral plots. Example for March. (b) significance coefficient ( $r^2$ ) of the linear regression and (c) values of the linear regression at each depth for March and October. The grey shading corresponds to the average mixed layer (ML) in the region for the five days between March 25th and 29th (dark grey), and between October 25th and 29th (light grey)

For this study, we have used the COAS model with full temperature/density field sampling and derived these statistical linear correlations in our region, at each model vertical level, and in contrasting seasons. For a wider application based on observations only, these linear correlations would need to be derived globally from Argo/hydrography profiles, for example, and with monthly variations.

## Appendix B SWOT orbit and coverage in the study area

The Sea Level Anomaly (SLA) swaths used to recreate the 2D field for applying the sQG methodology to the SWOT data are created using the beta pre-validated L3 data (doi:10.24400/527896/A01-2023.017) published by AVISO (AVISO/DUACS, 2023). During the 3-year science phase, SWOT covers the region of interest for this study completely in 21-days. However, due to the orbit coverage, there are four windows in which the satellite observes the region during five consecutive days, with swaths laid down in two sets of alternate ascending and descending passes that fully cover the area (see Figure B1a). These 5-day neighbouring swaths are then used to create a 2D high-resolution SSH map over 600 x 600 km, by interpolating neighbouring swaths using a spline interpolation (see Figure B1b). This field is then filtered at 20 km to imitate the sQG processing described in the main text, and to reduce the SWOT noise, *e.g.* from residual waves, etc. The corresponding DUACS 2D field centred on the same date is illustrated in Figure B1c for comparison.



**Figure B1.** (a) L3 SWOT Sea Level Anomaly over five days in October, (b) spline interpolation of the previous field, (c) DUACS SSH centred in the central day of SWOT observations, (d) SSH spectra of the 2D SWOT field, the DUACS field on the same region and period, and of COAS in the same area and season

The SSH wavelength spectra of the 2D interpolated SWOT field, the DUACS field and the full COAS simulation are represented in Figure B1d. The three spectra match at large scales, down to  $\sim 70$  km. Below, DUACS presents lower energy than COAS and SWOT, and SWOT has slightly higher energy than COAS from 25–50 km. Different processes could cause this. On one hand, SWOT data and COAS have a better representation of small scales than the DUACS product, leading to improved vertical dynamics in the region. On the other hand, this part of the spectrum could also include SWOT errors from new altimetric corrections at 2 km scales and from residual instrumental noise. The increased SWOT energy compared to COAS from at 20–70 km could come from the U-NET filter applied to the data for noise mitigation (Tréboutte et al., 2023). Conclusions on this require further analysis and should be updated when validated data are available later in 2024.

## Open Research

The code developed for the SSH-based diagnostics was created using Python 3.9 and is available upon request. The Jet Propulsion Laboratory has provided the COAS simulation outputs. The SWOT\_L3\_SSH product, derived from the L2 SWOT KaRIn Low rate ocean data products (NASA/JPL and CNES), is produced and made freely available by AVISO and DUACS teams as part of the DESMOS Science Team project”. AVISO/DUACS., 2023. SWOT Level-3 SSH Basic (v0.3) [Data set]. CNES. <https://doi.org/10.24400/527896/A01-2023.017>. The ARMOR 3D data are open source and available at [https://data.marine.copernicus.eu/product/MULTIOBS\\_GLO\\_PHY\\_TSUV\\_3D\\_MYNRT\\_015\\_012/description](https://data.marine.copernicus.eu/product/MULTIOBS_GLO_PHY_TSUV_3D_MYNRT_015_012/description)

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