

1 **Barrier islands face a gradual path toward drowning**  
2 **under most sea level rise scenarios**

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6 **Key Points:**

- 7 • Rapid sea level rise makes inlets expand until barrier islands disappear  
8 • Barrier drowning lags increases in rates of sea level rise by hundreds of years  
9 • Higher rates of sea level rise cause earlier, intensified drowning while higher waves  
10 cause more drowning or more frequent inlet closure

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## Abstract

The expected increase in rates of sea level rise during the 21st century and beyond may cause tidal inlets to expand and barrier islands to drown. However, many aspects remain unclear, e.g., the timescales involved in the drowning process have received little attention. To gain insight into the morphodynamics of barrier systems subject to sea level rise, we here present results obtained with a novel barrier island model, BRIE-D. This new model allows for changes in the alongshore extent of the barrier lying below sea level. These concern reductions in barrier width, barrier height, as well as lateral expansion of tidal inlets. Model results show that the evolution of barrier islands is susceptible to the wave height and the rate of sea level rise that they experience. It takes hundreds of years for barrier islands to drown in response to high rates of sea level rise (more than 15 mm/yr). Furthermore, increasing rates of sea level rise cause an earlier and more severe barrier drowning in environments with low waves. Barrier systems that face higher waves can undergo more frequent inlet closures (due to a larger amount of sediment imported into the inlets), but also the degree of barrier drowning might increase (due to a deepening of the toe of the shoreface). The latter process dominates over the former when rates of sea level rise are higher than 5 mm/yr.

## Plain Language Summary

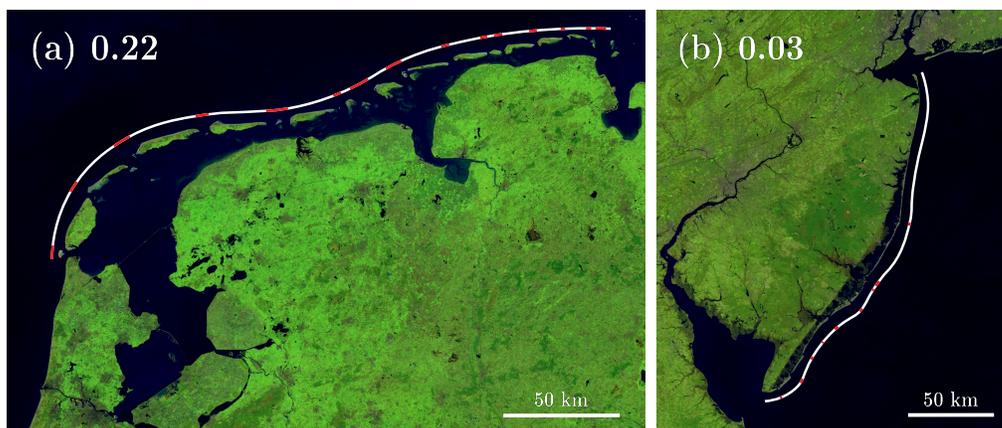
Barrier islands respond to an increase in sea level by migrating landward. Furthermore, inlet width and/or number may increase as portions of the barrier drown. In extreme sea level rise scenarios (like those predicted during the 21st century and beyond) barrier islands may not be able to migrate landward fast enough to keep up with sea level rise. In these situations, barrier islands will (partially) drown. Nowadays, it is difficult to predict when and under which conditions this drowning may occur. To better understand the dynamics of a drowning barrier island, we adapted a pre-existing numerical model. The new aspects of the model are that inlet widths may change gradually, depending on the barrier response to sea level rise and sediment availability. It was found that rates of sea level rise and the height of incoming waves are the key drivers that determine the long-term fate of barrier systems. Higher rates of sea level rise result in a larger portion of the barrier that is drowned. When wave heights are increased, the inlets tend to close more easily (when rates of sea level rise are small to moderate, i.e., less than 5 mm/yr), or there is more drowning (for higher rates of sea level rise).

## 1 Introduction

Barrier islands are low-lying coastal land forms that constitute 10 – 15% of the world’s coasts. They lie parallel to the mainland coast, thereby they protect it from coastal hazards such as storm surges (FitzGerald et al., 2006). As most coastal lowlands are densely populated, they are thus of great socio-economic importance.

Most barrier islands were created during the Holocene, when rates of sea level rise (SLR) decreased from 7–15 mm/yr to  $\sim 2$  mm/yr (Leatherman, 1983; Beets & van der Spek, 2000; Mariotti, 2021). With this change, sediment deposition could gradually catch up with SLR and compensate for its creation of accommodation. If all inlets filled up, then uninterrupted barrier coasts were formed. However, in some areas sediment supply was insufficient to fill up the inlets completely, then barrier islands were formed. The latter was the case of the Wadden Islands along the Dutch, German and Danish coast, which formed around 7000 yrs BP (Beets & van der Spek, 2000), as well as that of the barrier islands along the US east coast (Figure 1).

Future projected SLR is an important threat to most coastal systems in the world. Worst case scenarios predict a global mean sea level (MSL) increase of roughly 2.7 m by the year 2300 with respect to that of the year 2000 (Palmer et al., 2020). Furthermore,



**Figure 1.** Examples of barrier islands and their respective fraction below mean sea level: (a) 0.22 for the Wadden Islands along the coasts of the Netherlands and Germany, and (b) 0.03 for the barrier islands along the US east coast of New Jersey. White curves represent the islands extent, while red curves represent that of inlets. The given fraction below MSL is computed as the ratio between inlet extent and the total barrier chain extent (inlets and islands). Extracted from Google Earth (images provided by TerraMetrics).

60 the effects of e.g., vertical land motion should also be considered when studying the re-  
 61 sponse of coastal systems to changes in sea level. Given that many barrier islands are  
 62 located near deltas, where land is sinking, they may experience high relative rates of SLR.  
 63 Another possible consequence of climate change are variations in storm return periods,  
 64 which can also affect barrier coasts through changes in barrier breaching and sediment  
 65 transport during overwash events (Reef et al., 2020).

66 One of the possible consequences of this increase in sea level, is that barrier islands  
 67 will not be able to migrate landward fast enough, thus resulting in “whole scale” bar-  
 68 rier island drowning (Gilbert, 1885; Storms et al., 2002; Lorenzo-Trueba & Ashton, 2014;  
 69 Mellett & Plater, 2018; Nienhuis & Lorenzo-Trueba, 2019a). Observations of drowning  
 70 of barrier systems as a whole are scarce. There is an example in the English Channel,  
 71 where a barrier was formed around 9500–8800 yrs BP, when MSL was at  $-22$  m rela-  
 72 tive to that of present day, and it drowned around 8300 yrs BP when MSL reached  $-17$  m  
 73 (Sanders & Kumar, 1975).

74 Partial barrier drowning is more common. With high rates of SLR, the fraction of  
 75 a barrier island chain lying below MSL (inlets, see Figure 1) is expected to increase. More  
 76 inlets are formed or existing inlets become wider (FitzGerald et al., 2018). Inlets that  
 77 might have been in equilibrium, due to a balance between sediment export by tidal cur-  
 78 rents and sediment import by littoral drift (Escoffier, 1940), will deviate from that equi-  
 79 librium when the MSL changes. An example of such a situation is the evolution of the  
 80 Isles Dernières, Louisiana since the mid-1800s. At that time, a drowning process started  
 81 to take place, resulting in narrowing of the barrier and creation of new inlets, as well as  
 82 widening of existing inlets (FitzGerald et al., 2008).

83 In part because of the scarcity of observations, we remain unable to predict when  
 84 and where barriers will drown. Modern barrier islands might be at equilibrium with the  
 85 tides and waves, but they might also already show signs of drowning. The timescales in-

86 volved in barrier island drowning from present-day SLR are mostly unknown, and could  
87 be long (Mariotti & Hein, 2022).

88 Models have been developed to better understand barrier dynamics, and barrier  
89 drowning. Cowell et al. (1992) developed a cross-shore model that describes the evolu-  
90 tion of the active shoreface profile subject to SLR and to losses and gains of sediment  
91 beyond the active profile. Masetti et al. (2008) developed a more detailed cross-shore model,  
92 where additional sediment dynamics are included. Their model is able to represent over-  
93 wash, eolian processes, and sediment input from rivers in the back-barrier lagoon. This  
94 model was used to study the dependency of barrier island drowning on overwash fluxes.  
95 Lorenzo-Trueba and Ashton (2014) also designed a cross-shore model to study barrier  
96 island drowning and retreat due to SLR. Their model uses a more idealized domain than  
97 that of Masetti et al. (2008), but is able to represent different dynamic equilibrium con-  
98 figurations of the barrier.

99 These cross-sectional models can represent the process of barrier drowning, but to  
100 study chains of barrier islands, two horizontal dimensions are required. Such models have  
101 been recently developed (Ashton & Lorenzo-Trueba, 2018; Nienhuis & Lorenzo-Trueba,  
102 2019b; Mariotti, 2021). The model by Ashton and Lorenzo-Trueba (2018) follows the same  
103 parameterized cross-shore dynamics as that of Lorenzo-Trueba and Ashton (2014), and  
104 couples them in the alongshore direction with a shoreline diffusivity. The BRIE model  
105 of Nienhuis and Lorenzo-Trueba (2019b) presents a more realistic picture by also account-  
106 ing for inlet dynamics. Similarly, the model of Mariotti (2021) also accounts for inlet dy-  
107 namics, but uses a more process-based approach. The advantage of the parameterized  
108 BRIE model is that it is fast, so it is a suitable tool for performing extensive sensitiv-  
109 ity studies. Furthermore, since it explicitly accounts for inlet dynamics, such as open-  
110 ing, closing, or migration, it can offer insights into partial drowning, where only a frac-  
111 tion of the chain is below sea level. Tidal inlets in BRIE, however, cannot expand be-  
112 yond their equilibrium state, and therefore the model struggles to appropriately quan-  
113 tify the effect of SLR on the size and/or number of inlets.

114 Motivated by these knowledge gaps, we modify and expand the BRIE model into  
115 the BRIE-D model to allow for SLR-driven transformations of tidal inlets on barrier is-  
116 land chains. The new aspects implemented in the BRIE-D model concern a better rep-  
117 resentation of the process of inlet expansion, therefore allowing a gradual increase in the  
118 alongshore extent of the barrier lying below MSL. The evolution of the inlet width de-  
119 pends on the distribution of alongshore sediment transport within the inlet, the exchange  
120 of sediment with the flood-tidal delta, merging with other inlets, and drowning of por-  
121 tions of the barrier.

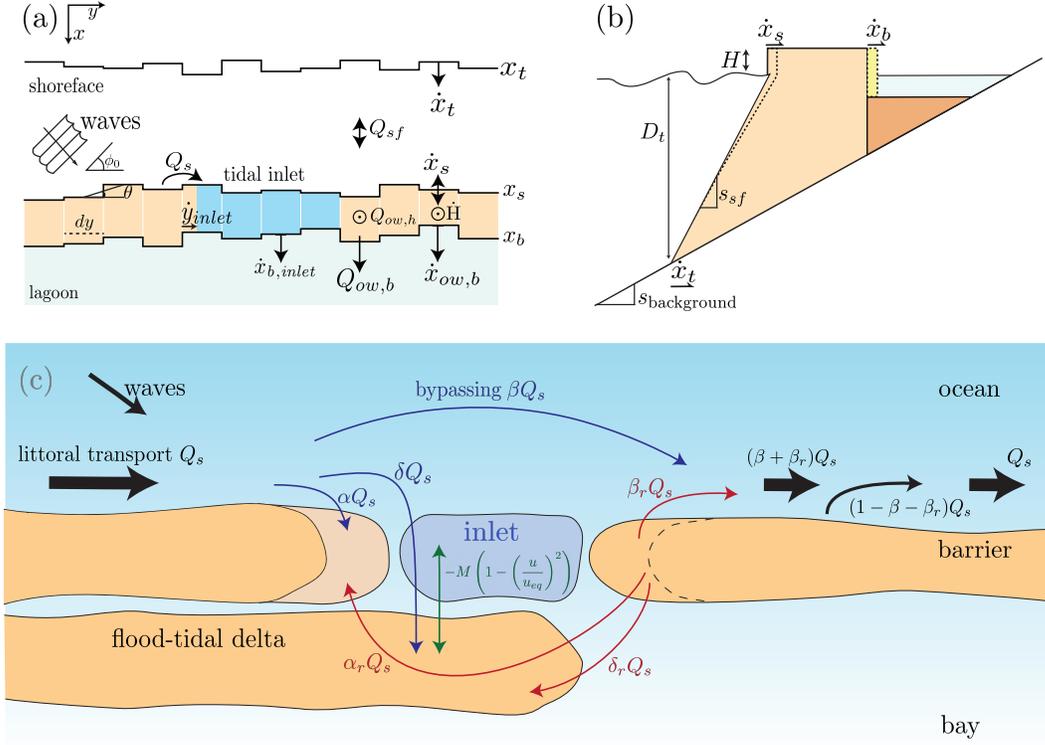
122 Our study objectives are to (1) understand how the fraction of the barrier lying  
123 below MSL is affected by SLR-driven inlet widening by comparing the outcomes of the  
124 BRIE model with those of the BRIE-D model, (2) examine the temporal evolution of  
125 the fraction of the barrier that is drowned, as well as to quantify drowning timescales,  
126 and (3) explore the dependence of barrier island drowning on model parameters.

127 The next section includes the model description, design of simulations and anal-  
128 ysis of model output. Section 3 contains the results, followed by a discussion in Section 4.  
129 The final section contains the conclusions.

## 130 2 Methods

### 131 2.1 Model Description

132 Here we give an overview of the physics represented in the BRIE-D model, with  
133 a detailed model description and relevant equations given in the Supplementary Infor-



**Figure 2.** Schematized model domain: (a) plan view highlighting the three moving boundaries (toe of the shoreface  $x_t$ , shoreline  $x_s$  and back-barrier shoreline  $x_b$ ) and barrier height  $H$ , as well as the sediment transports determining their evolution. Vectors indicate the direction of potential changes, with the dot symbolizing barrier heightening. We use  $\dot{x}$  to represent the local time derivative. (b) Cross-shore view of the barrier showing landward migration in terms of  $\dot{x}_s$  and  $\dot{x}_b$ . (c) Mass balance of the boxes used in modeling inlet dynamics, including sediment exchange with updrift and downdrift tips of the barrier, as well as with the flood-tidal delta. The parameters  $\alpha$ ,  $\beta$ ,  $\delta$ ,  $\alpha_r$ ,  $\beta_r$  and  $\delta_r$  denote fractions of the littoral transport  $Q_s$ . Note that the flood-tidal delta extends through the updrift barrier because it has been building up as the inlet was migrating. Modified from Nienhuis and Lorenzo-Trueba (2019b) and Nienhuis and Ashton (2016). A detailed description on the moving boundaries and the sediment exchange within the inlet is given in SI1.

134 mation (SI1). Differences between the BRIE-D and the BRIE model involve the inlet  
 135 evolution (Section 2.1.4), all other routines are equivalent in both models.

136 **2.1.1 Domain**

137 The BRIE-D model uses an idealized model domain with the  $x$ - and  $y$ - axis point-  
 138 ing perpendicular and parallel to the barrier, respectively. The  $z$ - axis points upward,  
 139 with  $z = 0$  representing MSL at  $t = 0$ . The domain comprises the shoreface, the sub-  
 140 aerial part of the barrier, as well as the back-barrier lagoon (see Figure 2).

141 The model is driven by tides (with a prescribed range), rate of SLR and waves that  
 142 propagate from deep water toward the barrier, among other boundary conditions. The  
 143 waves are characterized by a significant wave height (assumed to be constant) and by

144 an angle of incidence with respect to the  $y$ -axis. This angle is a stochastic variable that  
 145 is determined by a probability density distribution that depends on wave asymmetry (frac-  
 146 tion of waves approaching from the left, looking offshore) and wave highness (fraction  
 147 of waves approaching at a high angle, i.e.,  $|\phi_0| > 45^\circ$ , see Figure 2a).

### 148 **2.1.2 Cross-shore Dynamics**

149 The BRIE-D model describes the barrier system in the cross-shore direction as the  
 150 time evolution of barrier height and the location of three boundaries: the toe of the shoreface,  
 151 and the shorelines of the barrier on the sea side,  $x_s$ , and back-barrier side,  $x_b$ . We de-  
 152 fine the toe of the shoreface as the location  $x = x_t$  where there is negligible cross-shore  
 153 sediment exchange between the shoreface and the shelf (Ortiz & Ashton, 2016). This cor-  
 154 responds to a depth  $z = -D_t$  at  $t = 0$ . The depth of closure is set to  $D_t = 8.9H_s$ , so  
 155 it depends on the significant wave height  $H_s$  at deep water (Houston, 1995), which is held  
 156 constant in time in our simulations.

157 The cross-shore depth profile of the shoreface tends toward a prescribed equilib-  
 158 rium defined by the balance between onshore sediment transport by waves and offshore  
 159 directed transport due to gravity. The shoreface sediment transport  $Q_{sf}$  is directed off-  
 160 shore when the bottom slope of the shoreface  $s_{sf}$  is larger than the equilibrium slope  $s_{sf,eq}$ ,  
 161 otherwise it is directed onshore. As revealed by Figure 2,  $s_{sf} = D_t/(x_s - x_t)$ , where  
 162  $x = x_t$  and  $x = x_s$  are the cross-shore positions of the toe of the shoreface and of the  
 163 seaward shoreline of the barrier, respectively. Furthermore,  $s_{sf,eq}$  is determined by the  
 164 long-term wave conditions. The transport  $Q_{sf}$  is one of the drivers that determine time  
 165 changes in both  $x_t$  and  $x_s$  (see S11 for details). SLR is a second driver of changes in  $x_t$ ,  
 166 and it causes a decrease in barrier height  $H$ .

167 Sediment transport during overwash connects the shoreface with the back-barrier  
 168 lagoon. Part of the sediment overwash accumulates on top of the barrier ( $Q_{ow,h}$ ) and  
 169 part of it is deposited in the back-barrier lagoon ( $Q_{ow,b}$ ). The first contribution results  
 170 in an increase in the height  $H$  of the barrier, while the latter results in a landward mi-  
 171 gration of the back-barrier shoreline, denoted by  $x = x_b$ . Both sediment overwash trans-  
 172 port fluxes scale with their associated deficit volumes, which represent the difference be-  
 173 tween a current barrier configuration and one that is both high and wide enough such  
 174 that overwash is presumed not to occur. The latter situation is characterized by the crit-  
 175 ical barrier width and critical barrier height (Jiménez & Sánchez-Arcilla, 2004; Lorenzo-  
 176 Trueba & Ashton, 2014).

### 177 **2.1.3 Inlet Opening**

178 Inlets may open due to barrier breaching caused by a storm or due to barrier drown-  
 179 ing. Breaching is imposed every 10 yrs where the barrier volume is at a minimum, and  
 180 at a location at least 5 km from existing inlets. Alternatively, inlets appear when a por-  
 181 tion of the barrier drowns (either because the width or the height of the barrier becomes  
 182 negative), which is not restricted to its proximity to other inlets. We set the initial width  
 183 of a breached inlet to 1 km, while for a drowned inlet we set it to the width of the por-  
 184 tion of the barrier that drowned.

### 185 **2.1.4 Inlet Evolution**

186 We depart from the formulation used in the BRIE model by allowing for variations  
 187 in the cross-sectional area of the inlet. In the BRIE model, the inlets were assumed to  
 188 instantly attain their equilibrium cross-sectional area (following Escoffier, 1940). In the  
 189 BRIE-D model, we impose a gradual temporal evolution of the inlet size. We distinguish

between four different sources of variations in the cross-sectional area of the inlet  $A_{inlet}$ ,

$$\frac{dA_{inlet}}{dt} = G_{sd} + G_{Esc} + G_m + G_d , \quad (1)$$

where  $G_{sd}$  represents the variations due to the distribution of the alongshore transport within the inlet,  $G_{Esc}$  represents the changes in cross-sectional area due to sediment exchange between the inlet and the flood-tidal delta (which is partly parameterized following the theory by Escoffier, 1940), and  $G_m$  and  $G_d$  account for the increase in the cross-sectional area of the inlet due to merging with another inlet and due to drowning of the barrier, respectively.

Each inlet receives an amount of sediment per time unit that is a fraction of the alongshore sediment transport  $Q_s$ , where the latter depends on wave climate and shoreline angle. A fraction  $\beta$  of  $Q_s$  bypasses sediment to the coast of the downdrift island, another fraction  $\delta$  toward the flood tidal delta, where the sediment gets deposited. The remaining fraction  $\alpha$  transports sediment to the tip of the updrift island, where sediment gets deposited (see Figure 2c). Sediment eroded from the tip of the downdrift island is distributed similarly. This sediment distribution, which depends on tidal prism and wave characteristics, causes inlet migration and variations in the volume of the downdrift and updrift tips of the barrier,  $V_{down}$  and  $V_{up}$ , respectively. These variations in volume of the barrier cause changes in the cross-sectional area of the inlet,

$$G_{sd} = D_{inlet} \left( \frac{1}{A_{b,down}} \frac{dV_{down}}{dt} - \frac{1}{A_{b,up}} \frac{dV_{up}}{dt} \right) . \quad (2)$$

Here,  $A_{b,down}$  and  $A_{b,up}$  correspond to the cross-sectional area of the barrier downdrift and updrift of the inlet, and  $D_{inlet}$  represents the inlet depth. In the BRIE model this sediment distribution was such that the cross-sectional area of the inlet was maintained constant. Differently, in the BRIE-D model we allow for both tips of the barrier to be disconnected, and grow or shrink the inlet. A detailed description of the variations in updrift and downdrift volumes is given in SI1.

We also allow for variations in the cross-sectional area of the inlet depending on the balance in sediment exchange with the flood-tidal delta. This balance depends on a prescribed transport from the flood-tidal delta to the inlet and the export of sediment from the inlet to the flood-tidal delta due to tidal currents. We use a simple model for an inlet-bay system, as was derived by Brown (1928) and used by Escoffier (1940) to explain the stability of tidal inlets. The changes on the cross-sectional area of the inlet governed by these dynamics are described by

$$G_{Esc} = -\frac{M}{W_b} \left( 1 - \left( \frac{U}{U_e} \right)^2 \right) . \quad (3)$$

In this equation,  $W_b$  is the width of the barrier,  $U$  is the amplitude of the tidal current in the inlet (which depends on the imposed tidal range at sea, the cross-sectional area of the inlet, the barrier width, and the wetted surface of the back-barrier lagoon),  $U_e$  is the amplitude of the tidal current at equilibrium (set at 1 m/s for all simulations), and  $M$  is the volume of sediment per time unit that the inlet receives from the flood tidal delta. With this representation of tidal dynamics we allow for the inlet to evolve toward an equilibrium configuration, using a parametrisation of the net sediment transport due to tides that was earlier used by van de Kreeke (1998, 2004). This evolution differs from the behavior imposed in the BRIE model, where the inlet instantly attained its equilibrium cross-sectional area.

The increase in the cross-sectional area of the inlet due to merging with other inlets is such that the total cross-sectional area is conserved,

$$G_m = \sum_{i=1}^{N_m} \frac{dA_i}{dt} . \quad (4)$$

236 Here,  $N_m$  is the number of inlets with which the considered inlet is merging, and  $A_i$  are  
 237 their respective cross-sectional areas.

238 Lastly, the increase in the cross-sectional area of the inlet due to barrier drown-  
 239 ing depends on the length of the portion of the barrier that drowned  $W_d$  (either due to  
 240 negative barrier width or negative barrier height),

$$241 \quad G_d = \frac{dW_d}{dt} \gamma_{aspect}^2 W_d, \quad (5)$$

242 where the corresponding depth is computed using the inlet aspect ratio  $\gamma_{aspect}$  (with  $\gamma_{aspect}^2 =$   
 243  $D_{inlet}/W_{inlet}$ ).

244 The difference between the BRIE and the BRIE-D models is the temporal evolu-  
 245 tion of the width of the inlets. In the BRIE model, inlets instantly achieved a width de-  
 246 fined by the equilibrium between sediment import by the littoral drift and sediment ex-  
 247 port by tidal currents. In contrast, the width of inlets in the BRIE-D model gradually  
 248 adapts depending on the four terms present in Equation 1, and defined in Equations 2 –  
 249 5. Given that these differences affect the dynamics of the whole barrier chain, not only  
 250 inlet widths are different in both models, but also inlet position and number may dif-  
 251 fer.

### 252 **2.1.5 Evolution of the Shoreline**

253 The shoreline evolves as a result of divergence of alongshore sediment transport,  
 254 which is parameterized using the CERC formula and the presence of cross-shore sedi-  
 255 ment motion. Overwash and inlet dynamics result in a sink of sediment for the shore-  
 256 line, while shoreface transport may result in a sink or source of sediment, depending on  
 257 its direction. Following Ashton and Murray (2006), this results in a forced diffusion equa-  
 258 tion (see Equation 70 of the SI), in which the diffusion coefficient depends on wave char-  
 259 acteristics and the orientation of the shoreline.

### 260 **2.1.6 Initial and Boundary Conditions**

261 Simulations are initialized with a barrier without inlets. The position of the sea-  
 262 ward shoreline  $x = x_s$  is computed imposing the equilibrium shoreface slope between  
 263  $x = x_t$  and  $x = x_s$  and adding a random perturbation following a uniform distribu-  
 264 tion between 0 and 1 m. The back-barrier shoreline is set such that the barrier width  
 265 equals its critical value. The barrier height is also set equal to its critical value. These  
 266 are representative values for barrier width and height (Leatherman, 1979). We apply pe-  
 267 riodic boundary conditions in the alongshore direction.

### 268 **2.1.7 Numerical Aspects**

269 The alongshore extent of the domain covers 50 km with a grid size of 100 m. We  
 270 solve the equation for the cross-sectional area of the inlet (Equation 1) using an Euler  
 271 forward scheme with a time step  $\Delta t = 0.05$  yr  $\sim$  18 days. The diffusion equation defin-  
 272 ing the shoreline evolution (Equation 70 of the SI) is solved using a Crank-Nicolson scheme  
 273 (Crank & Nicolson, 1947).

## 274 **2.2 Design of Simulations**

275 We run all simulations for 2500 yrs, taking a rate of SLR of  $\dot{\xi} = 2$  mm/yr dur-  
 276 ing the first 2000 yrs, which serve as model spin-up. After model spin-up, when the sys-  
 277 tem reaches a statistically stationary state in terms of inlet number and inlet migration  
 278 rates, we change  $\dot{\xi}$  in order to study the system response for another 500 yrs. All other  
 279 parameters are kept constant during the entire 2500 yrs (see Appendix A for a full overview

**Table 1.** Overview of simulations performed, imposing different values for the rate of SLR ( $\dot{\xi}$ ) and significant wave height ( $H_s$ ).

Aim	Model used	Parameter range <sup>a</sup>	Figure
Effects of inlet widening	BRIE, BRIE-D	$\dot{\xi} = 4, 17$ mm/yr	3, 4
Temporal evolution of barrier drowning	BRIE-D	$\dot{\xi} = 4, 17$ mm/yr	5
Dependence on model parameters	BRIE-D	$\dot{\xi}$ varying between 2 and 20 mm/yr, $H_s$ varying between 0.75 and 3 m <sup>b</sup>	6, 7, 8

<sup>a</sup> If not specified parameters take their default values (see Appendix A).

<sup>b</sup> Mulhern et al. (2017)

280 of the model parameters and their default values). The new  $\dot{\xi}$  is not changed during the  
281 last 500 yrs of model evolution.

282 We simulate barrier response to rates of SLR  $\dot{\xi}$  between 2 and 20 mm/yr. The fol-  
283 lowing equivalences can be considered at global scale over the next centuries: RCP2.6  
284 and  $\sim 5$  mm/yr, RCP4.5 and  $\sim 6$  mm/yr, RCP8.5 and  $\sim 10$  mm/yr (IPCC, 2021; Palmer  
285 et al., 2020). Simulations span a broad range of significant wave heights  $H_s$  (between  
286 0.75 and 3 m). Since this is a stochastic system, where randomness originates from the  
287 wave angle and from the initial value of  $x_t$ , we perform five model realizations for each  
288 parameter setting. The default parameter set includes a tidal amplitude of 0.8 m and  
289 a wave height of 1.5 m, which are representative values for a typical barrier island sys-  
290 tem, such as that in the Dutch Wadden Sea. Table 1 presents an overview of the sim-  
291 ulations performed.

292 We investigate the effects of barrier drowning on the widening of inlets by compar-  
293 ing the outcome of the BRIE-D model with that of the BRIE model for low and high  
294  $\dot{\xi}$  (4 and 17 mm/yr). We also study the temporal evolution of barrier island drowning  
295 for these same two rates of SLR. Moreover, we present the dependence of barrier drown-  
296 ing on a broad range of significant wave height and rates of SLR.

297 We present the model results as the mean of the five realizations for each param-  
298 eters setting. Errors are quantified using the standard error of the mean. Experiments  
299 performed with an ensemble size of 100 showed no significant differences in model out-  
300 come when compared to results computed with only five simulations.

301 We explored the sensitivity of model output to halving the grid size  $\Delta y$  and halv-  
302 ing the time step  $\Delta t$  and found that differences were smaller than 3% for the situation  
303 with default parameter values.

### 304 2.3 Analysis of Model Output

305 We quantify barrier drowning by the fraction of the barrier lying below MSL due  
306 to tide-wave imbalance in the inlet,

$$307 \quad \Delta F(t) = F(t) - F_{eq}(t) . \quad (6)$$

308 Here,  $F$  and  $F_{eq}$  are both fractions of the barrier lying below MSL, computed as sums  
309 of the widths of all inlets (a total of  $N(t)$ ) divided by the alongshore extent of the bar-  
310 rier,

$$311 \quad F(t) = \frac{\sum_{i=1}^{N(t)} W_{inlet,i}(t)}{L_b} , \quad F_{eq}(t) = \frac{\sum_{i=1}^{N(t)} W_{inlet,eq,i}(t)}{L_b} . \quad (7)$$

312 Their difference concerns the fact that  $F_{eq}$  is calculated by setting the tidal current am-  
 313 plitude in the inlet equal to its equilibrium value ( $U_e = 1$  m/s), from which the cor-  
 314 responding equilibrium cross-sectional area is computed (further details are given in SI1).  
 315 In Equation 7,  $L_b$  is the alongshore extent of the barrier,  $N(t)$  is the number of inlets  
 316 at time  $t$ ,  $W_{inlet}$  is the inlet width, and  $W_{inlet,eq}$  is the equilibrium inlet width defined  
 317 by the balance in sediment exchange in the inlet. Thus,  $\Delta F$  mainly represents the in-  
 318 crease of the fraction of the barrier lying below MSL due to the effects of flooding caused  
 319 by SLR.

320 Note that  $F_{eq}$  is not identical to the  $F$  that would be obtained when using the BRIE  
 321 model. This is because the time evolution of both models is governed by different dyn-  
 322 amics, so the number and distribution of inlets will be different in both models.

323 Increasing SLR does not only affect  $\Delta F(t)$ , but it may also affect  $F_{eq}(t)$ . The in-  
 324 let equilibrium width depends on the tidal prism, among other system characteristics,  
 325 which can also be affected by SLR. Indeed, an increase in sea level may result in an in-  
 326 crease in tidal prism due to a widening of the back-barrier lagoon. Thus, an increase in  
 327 SLR may induce an increase in the equilibrium inlet width, even if the system is not drown-  
 328 ing.

329 With these definitions for  $F$ ,  $F_{eq}$  and  $\Delta F$  we are able to quantify the different ef-  
 330 fects of SLR on barrier systems, from changes in equilibrium width, up to drowning of  
 331 the barrier due to a decrease in sediment availability.

332 We define drowning timescales as the time needed for  $\Delta F$  to take a value of 0.1  
 333 or 0.3. We also investigate the evolution of the number of inlets as well as that of bar-  
 334 rier width, summarized by its alongshore mean through time. The barrier width is com-  
 335 puted as the distance between the seaward shoreline and the back-barrier shoreline  $W_b =$   
 336  $x_b - x_s$ . We compute the barrier width only along the portions corresponding to sub-  
 337 aerial barrier, i.e., where  $W_b > 0$ .

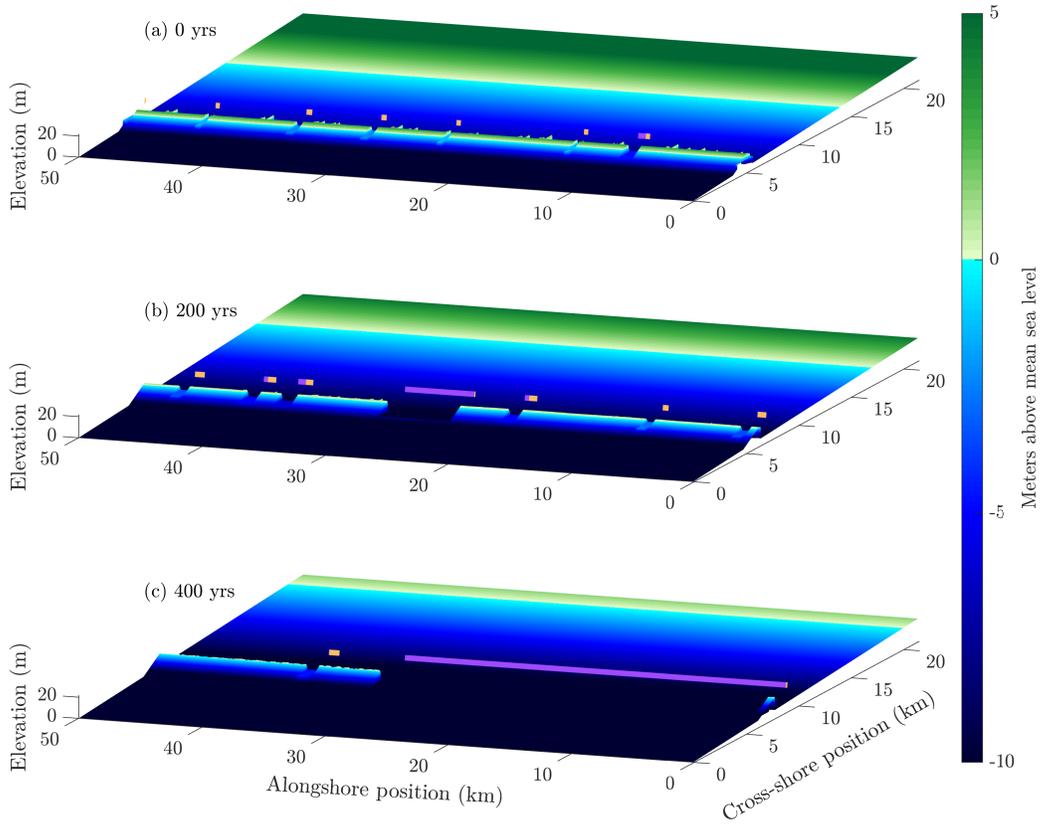
### 338 3 Results

#### 339 3.1 Manifestation of Drowning in BRIE-D Compared to BRIE

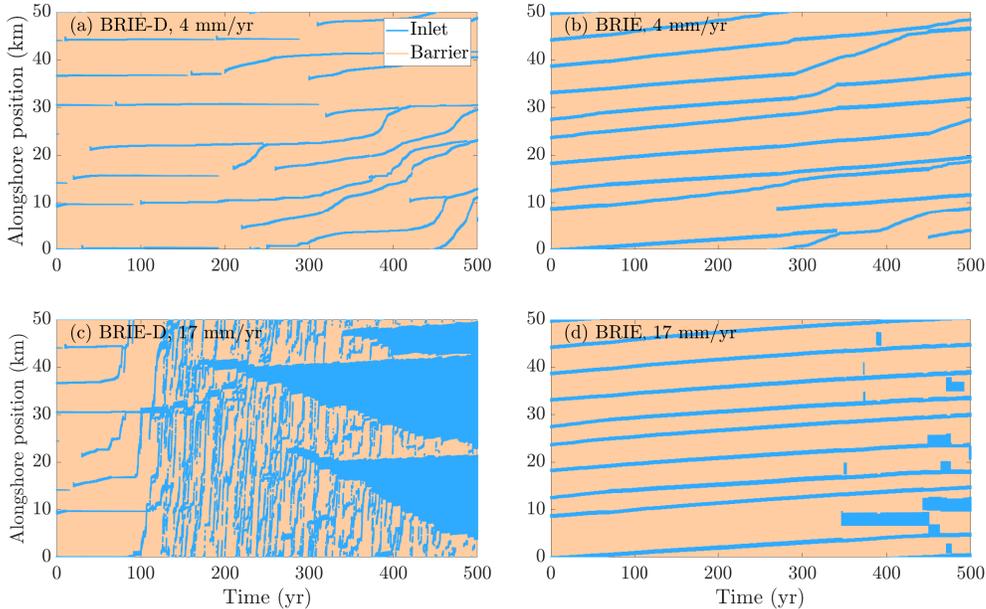
340 An example model simulation under a rate of SLR  $\dot{\xi} = 17$  mm/yr shows a grad-  
 341 ual expansion of inlets during 500 yrs of barrier evolution (Figure 3). There is a grad-  
 342 ual evolution towards a drowned barrier: initially (after the model spin-up period) the  
 343 barrier has achieved a statistical equilibrium state, after 200 yrs drowning has started,  
 344 and after 400 yrs more than half of the alongshore extent of the barrier lies below MSL.  
 345 The transition from a state in which inlets are in morphodynamic equilibrium towards  
 346 a state of drowning is evident after 200 yrs, as the number of inlets increases and some  
 347 inlets become much wider than in the equilibrium situation. From there on, some of the  
 348 inlets merge together until reaching a width of the order of tenths of km by the year 400.

349 There are differences between the outcomes of the BRIE and the BRIE-D model  
 350 already at low  $\dot{\xi}$  (e.g., 4 mm/yr, see Figure 4a,b), i.e., in situations where the inlets are  
 351 close to equilibrium and there is no drowning. Inlets tend to close more easily in the BRIE-D  
 352 model than in the BRIE model as a result of the new sediment dynamics imposed in the  
 353 inlet. Nevertheless, under these circumstances inlet width remains approximately con-  
 354 stant in time in the BRIE-D model, as it is the case with the BRIE model. Furthermore,  
 355 inlet migration rates are generally similar in both models ( $\sim 1-2$  m/yr), with the ex-  
 356 ception of short periods in which the BRIE-D model yields migration rates of order 200 m/yr,  
 357 due to local narrowing of the barrier.

358 The increase in inlet width observed in Figure 3 is the main result of our modifi-  
 359 cations in the BRIE model. The BRIE model is by definition not able to model situa-



**Figure 3.** Modeled barrier island evolution at years (a) 0, (b) 200 and (c) 400 after the model spin-up period. Simulation is for  $\dot{\xi} = 17$  mm/yr in order to visualize inlet widening caused by barrier drowning. The number of inlets increases substantially between years 0 and 200, and inlets get wider between years 200 and 400. Orange lines represent the equilibrium inlet width for each inlet ( $W_{inlet,eq}$ ) and purple lines the difference between the actual inlet width and that of equilibrium. Note the differences between actual and equilibrium inlet widths for years 200 and 400. All parameters except  $\dot{\xi}$  have their default values (see Table A1); in particular the offshore significant wave height is  $H_s = 1.5$  m and the tidal amplitude is  $a_0 = 0.8$  m.



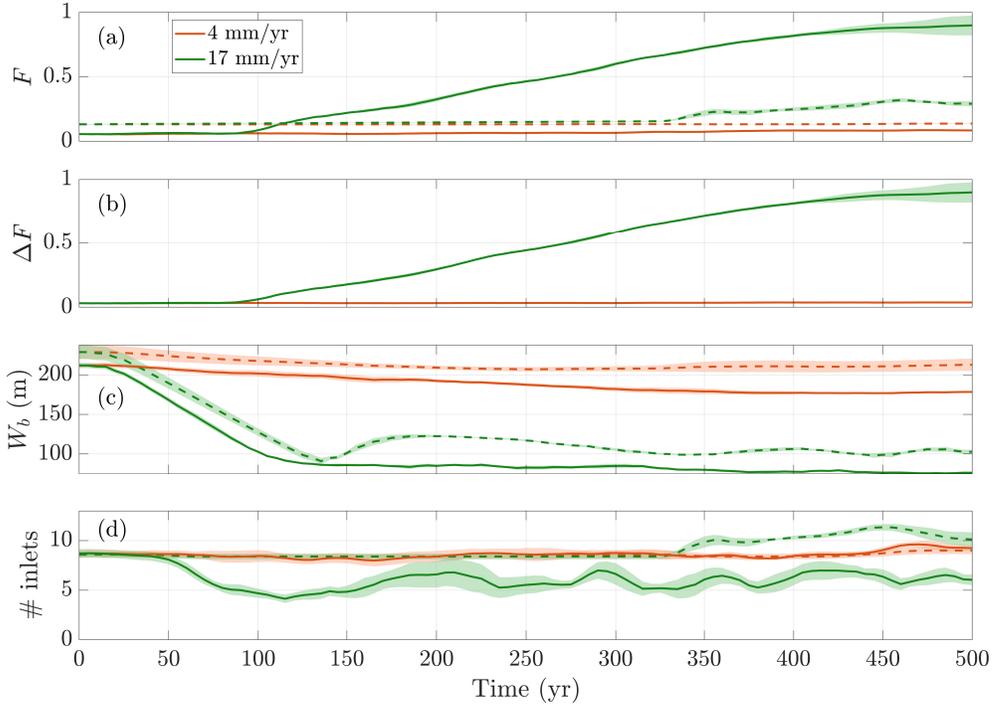
**Figure 4.** Comparison between the BRIE-D model and the BRIE model: Temporal evolution of barrier systems during 500 yrs in a 50 km long domain for a  $\dot{\xi}$  of (a,b) 4 mm/yr (barrier drowning is not occurring) and (c,d) 17 mm/yr (there is barrier drowning causing widening the inlets). Simulations (a,c) were performed using the BRIE-D model, whilst simulations (b,d) were performed using the BRIE model. All parameters except  $\dot{\xi}$  have their default values (see Table A1); in particular the offshore significant wave height is  $H_s = 1.5$  m and the tidal amplitude is  $a_0 = 0.8$  m.

360 tions in which the inlet width gradually increases due to drowning of the barrier. These  
 361 differences are stressed in Figure 4c,d, for the situation with a high rate of SLR ( $\dot{\xi} =$   
 362 17 mm/yr). In general, the BRIE model does not always present a continuous evolution  
 363 of inlet widths, with abrupt changes taking place for example at years 450 or 465, or with  
 364 inlet closing briefly after opening at years 350 – 400, yielding unrealistic behavior be-  
 365 cause of barrier drowning that was unconnected to (other) inlet dynamics. In contrast,  
 366 with the adaptations implemented in the BRIE-D model in terms of the evolution of the  
 367 cross-sectional area of the inlet, we are able to account for widening of the inlet due to  
 368 barrier drowning, yielding a more smooth barrier behavior.

369 Both models also differ in inlet migration rates, with the BRIE-D model yielding  
 370 higher migration rates ( $\sim 5$  km/yr) than in the BRIE model ( $\sim 10$  m/yr) for narrow  
 371 inlets ( $< 2$  km) in drowning barriers. High migration rates in BRIE-D are due to the  
 372 barrier being very narrow ( $< 100$  m, see Figure 5b). It is unclear whether this result  
 373 is realistic. This discrepancy between the two models is caused by the disconnection of  
 374 the updrift and downdrift barrier tips, imposed to allow for inlet widening beyond its  
 375 equilibrium state. This disconnection causes differences in sediment deposition in the in-  
 376 let, which can alter inlet migration.

### 377 3.2 Temporal Evolution of Barrier Drowning

378 The mean fraction of the barrier lying below MSL,  $F$ , gradually increases from the  
 379 year 100 for a situation with a high rate of SLR ( $\dot{\xi} = 17$  mm/yr, see Figure 5a) up to



**Figure 5.** (a) Time series of the fraction of the barrier lying below MSL  $F$  for  $\dot{\xi} = 4$  mm/yr and  $\dot{\xi} = 17$  mm/yr, for both the BRIE-D model (solid lines) and the BRIE model (dashed lines). (b) As (a), but for the fraction of the barrier lying below MSL due to tide-wave imbalance in the inlet ( $\Delta F$ ) for the BRIE-D model. Note that  $\Delta F = 0$  for the BRIE model. (c) As (a), but for mean barrier width. (d) As (a), but for the number of inlets. Curves represent the mean over five simulations. Shaded areas represent the standard error of the mean, which is very low for some of the variables (e.g.,  $F$  until year 400). Dashed black lines in panel (b) correspond to the situations depicted in Figure 8 ( $\Delta F = 0.1, 0.3$ ).

380  $\sim 0.8$  after 500 yrs. Note that the simulation performed with the BRIE model also shows  
 381 an increase in fraction of the barrier lying below MSL, caused by the increase in tidal  
 382 prism, reaching values up to 0.3 – 0.4. This increase in  $F$  for the BRIE model corre-  
 383 sponds to the sudden inlet creation and inlet widening taking place from the year 350  
 384 onward (see Figure 4d). Regarding the BRIE-D model, the temporal evolution of  $F$  re-  
 385 sults from the gradual inlet widening obtained from the year  $\sim 100$  (see Figure 4c). The  
 386 situation with low rate of SLR ( $\dot{\xi} = 4$  mm/yr) shows a constant fraction of the barrier  
 387 lying below MSL for both models.

388 In the BRIE-D model, inlet widening due to wave-tide imbalance,  $\Delta F$ , is the main  
 389 agent causing the increase in fraction of the barrier lying below MSL (see Figure 5b).  
 390 The fraction  $\Delta F$  starts to deviate from zero after 100 yrs of evolution under  $\dot{\xi} = 17$  mm/yr,  
 391 and achieves a value of 0.85 after 400 yrs more. In contrast, the simulation with  $\dot{\xi} =$   
 392 4 mm/yr never achieves a drowning situation, i.e.,  $\Delta F$  is always close to zero with max-  
 393 imum variations of 0.001. In this situation of low rate of SLR, the barrier is able to adapt  
 394 to the changes in MSL given that both  $F$  and  $\Delta F$  are kept constant. This means that  
 395 landward migration of the barrier is effective enough such that the tidal prism is not changed  
 396 and there is not drowning caused by SLR.

397 After  $\sim 110$  yrs, the simulation with  $\dot{\xi} = 17$  mm/yr attains  $\Delta F = 0.1$ . At this  
 398 same stage, the barrier response starts to differ from its previous behavior (see Figure 4c),  
 399 with inlets becoming notably wider than they were initially.

400 Barrier width rapidly decreases in the simulations with  $\dot{\xi} = 17$  mm/yr during the  
 401 first  $\sim 100$ – $150$  yrs of evolution after spin-up (see Figure 5c). These first  $100$ – $150$  yrs  
 402 represent a transition period in which the barrier is accommodating to the new  $\dot{\xi}$ , which  
 403 can also be seen in Figure 4c. There is also a minor decrease in barrier width for the case  
 404  $\dot{\xi} = 4$  mm/yr regardless of the inlets being always close to equilibrium. The barrier width  
 405 eventually reaches an equilibrium value that depends on the  $\dot{\xi}$  imposed. That value is  
 406 about  $30$  m larger when using the BRIE model instead of the BRIE-D model for the two  
 407 values of  $\dot{\xi}$  shown. This difference is due to the added inlet dynamics in the BRIE-D with  
 408 respect to the BRIE model. In particular, we allow for an exchange of sediment between  
 409 the inlet and the flood-tidal delta in the BRIE-D model, which can reduce or grow the  
 410 back-barrier.

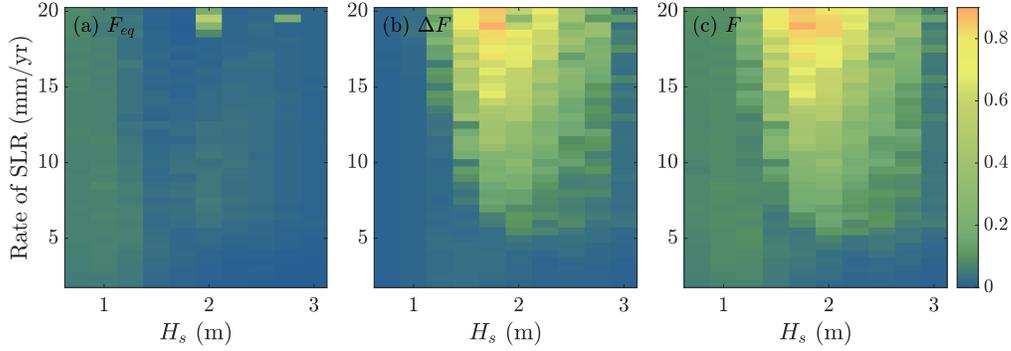
411 Both models produce roughly the same number of inlets for  $\dot{\xi} = 4$  mm/yr, because  
 412 the inlets are close to equilibrium, and in those scenarios the number of inlets is controlled  
 413 solely by the available tidal prism (i.e.,  $\Delta F \sim 0$ ). In this situation, inlet number is kept  
 414 constant at  $\sim 8$ – $9$ . Substantial differences between the two models arise for the sim-  
 415 ulations with  $\dot{\xi} = 17$  mm/yr. The number of inlets remains constant at around  $\sim 8$ –  
 416  $9$  for the BRIE model, showing no big differences with the situation with lower  $\dot{\xi}$  until  
 417 the year  $350$ . At that time, it increases up to  $\sim 11$ – $12$ . Note that the barrier behav-  
 418 ior for the BRIE simulations with  $\dot{\xi} = 17$  mm/yr starts to become irregular from the  
 419 year  $350$  as well (see Figure 4d), with inlets closing briefly after opening and sudden changes  
 420 in inlet width. In contrast, the inlet number decreases up to  $\sim 5$ – $6$  with the BRIE-D  
 421 model and stabilizes around this number from the year  $\sim 200$  onward. The reason that  
 422 for  $\dot{\xi} = 17$  mm/yr there are fewer inlets simulated by the BRIE-D model than by the  
 423 BRIE model is that inlets are wider, thus there is less subaerial portion of the barrier  
 424 where inlets can form and survive without merging with other existing inlets.

425 The change in rate of SLR that the barrier system undergoes at  $t = 0$  modifies  
 426 the barrier response. After this change, the barrier will attain a new statistical equilib-  
 427 rium state. It takes  $100$ – $150$  yrs for the barrier to adapt to the new conditions (Fig-  
 428 ure 5). This time lag in barrier response to variations in the rate of SLR is driven by the  
 429 gradual evolution imposed in inlet dynamics, which affect the dynamics of the whole bar-  
 430 rier chain.

### 431 3.3 Wave height and SLR effects on barrier drowning

432 We performed a sensitivity analysis, as described in Section 2.2. Of the paramet-  
 433 ers considered, the significant wave height  $H_s$  and the rate of SLR  $\dot{\xi}$  turned out to be  
 434 the parameters with the strongest impact on barrier drowning (see Appendix B).

435 The fraction of the barrier lying below MSL,  $F$ , may change due to the variations  
 436 induced by a new tide-wave balance (dependent on e.g., tidal prism, and wave-driven lit-  
 437 toral drift, represented by  $F_{eq}$ ), and by SLR-induced drowning (which causes tide-wave  
 438 imbalance, here represented by  $\Delta F$ ). The equilibrium fraction of the barrier lying be-  
 439 low MSL,  $F_{eq}$ , shows a dependence on significant wave height  $H_s$  and rate of SLR  $\dot{\xi}$  (see  
 440 Figure 6a). This dependence is mainly caused by variations in tidal prism and sediment  
 441 imported by waves. For example, larger significant wave heights cause a decrease in  $F_{eq}$   
 442 for low rates of SLR, as it is the case for the number of inlets. Nevertheless, the vari-  
 443 ations in the fraction of the barrier lying below MSL caused by the changes in equilib-  
 444 rium inlet widths are low compared to the effects of drowning (see Figure 6b). Thus, the  
 445 variations in the total fraction of the barrier lying below MSL,  $F$ , are mainly dominated  
 446 by the tide-wave imbalance in the inlet (i.e., the widening of inlets due to mainly SLR,



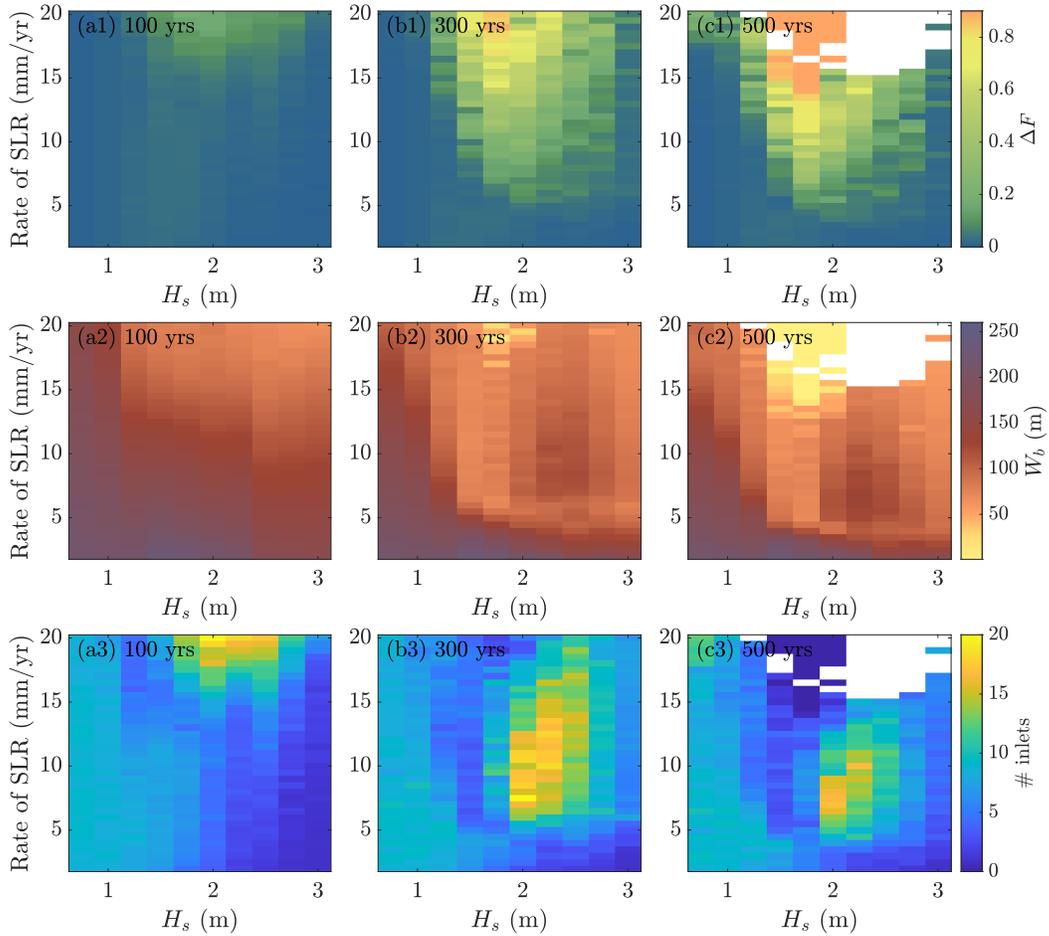
**Figure 6.** For different values of significant wave height  $H_s$  and rate of SLR  $\dot{\xi}$ : color plots of the (a) the fraction of the barrier lying below MSL assuming an equilibrium situation for the inlets ( $F_{eq}$ ), (b) the fraction of the barrier lying below MSL due to tide-wave imbalance in the inlet ( $\Delta F$ ), (c) the fraction of the barrier lying below MSL ( $F_{eq} + \Delta F = F$ ) at the year 300.

447  $\Delta F$ ). The behavior of  $F$  is only dominated by that  $F_{eq}$  for low rates of SLR ( $\dot{\xi} < 5$  mm/yr),  
 448 where the effect of SLR is lower (see Figure 6c).

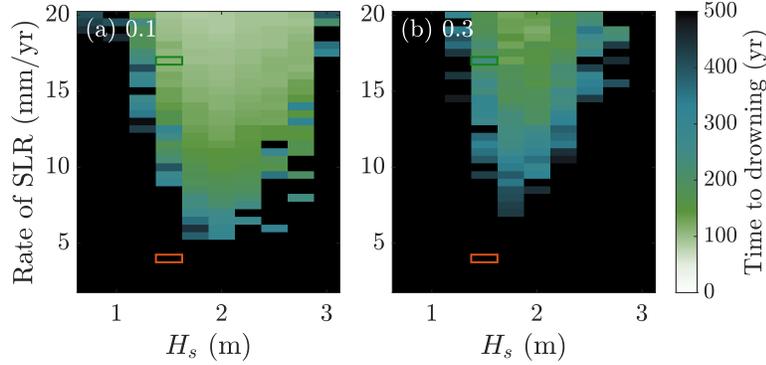
449 Color plots of fraction of drowned barrier due to wave-tide imbalance  $\Delta F$ , barrier  
 450 width  $W_b$  and number of inlets as a function of  $H_s$  and  $\dot{\xi}$  are shown in Figure 7 at years  
 451 100, 300 and 500 after the model spin-up period. The situations depicted in white in pan-  
 452 els (c2,c3) correspond to simulations that became numerically unstable while inlets were  
 453 widening due to barrier drowning and thus stopped before reaching the year 500. In some  
 454 other simulations, the barrier totally drowned before the year 500. For visualization pur-  
 455 poses, we have set the barrier width as well as the number of inlets to zero, and  $\Delta F$  to  
 456 0.9 in the latter situations. The quantities shown in all panels are computed as the mean  
 457 of five simulations.

458 The general tendency is that an increase in  $\dot{\xi}$  causes more drowning, as  $\Delta F$  eventu-  
 459 ally takes larger values (see Figure 7a1, b1, c1). The fraction of the barrier lying be-  
 460 low MSL due to tide-wave imbalance in the inlet,  $\Delta F$ , deviates from zero for rates of SLR  
 461 larger than a certain threshold ( $\dot{\xi} \sim 6$  mm/yr). For  $\dot{\xi}$  lower than 6 mm/yr, maximum  
 462 differences in  $\Delta F$  are 0.04 by the year 500. A similar general tendency can be seen for  
 463 the barrier width  $W_b$  (see Figure 7a2, b2, c2), which attains lower values at latter times  
 464 and at higher  $\dot{\xi}$ .

465 The sensitivity of the number of inlets on  $\dot{\xi}$  and  $H_s$  does not show such a clear pat-  
 466 tern as that of  $\Delta F$  or  $W_b$  (see Figure 7a3, b3, c3). Specifically, there are some cases with  
 467 SLR-driven drowning with a low number of inlets with (some of) them being very wide  
 468 ( $W_{inlet} \sim 10-20$  km), as it is the case for  $\dot{\xi} = 17$  mm/yr (see Figures 4c, 5). In other  
 469 cases with barrier drowning, widths of inlets overall take lower values ( $W_{inlet} \sim 1 -$   
 470 5 km). Still, the total fraction below MSL is larger than that at equilibrium, because the  
 471 number of inlets is very large ( $\sim 15-20$ ). Situations in which there is barrier drown-  
 472 ing with a large number of relatively narrow inlets are characterized by high waves ( $H_s \geq$   
 473 2 m) and rates of sea level rise generally lower than 15 mm/yr. Under these situations,  
 474 there is an important deposition of sediment by the waves, which creates narrower in-  
 475 lets. In contrast, drowning situations with few and wide inlets only take place for  $\dot{\xi} >$   
 476 15 mm/yr. In these cases, the combined effect of the deepening of the toe of the shoreface  
 477 (which depends on the wave height) and sea level rise causes a widening of the inlets which  
 478 cannot be balanced by the sediment import of waves. Thus, simulations with similar  $\Delta F$   
 479 and  $W_b$  may have a significantly different number of inlets. Still, for low  $\dot{\xi}$  (such that  $\Delta F \sim$   
 480 0, i.e.,  $\dot{\xi} < 5$  mm/yr) the number of inlets decreases significantly for  $H_s \geq 2$  m. This



**Figure 7.** For different values of significant wave height  $H_s$  and rate of SLR  $\dot{\xi}$ : color plots of  $\Delta F$  (**a1,b1,c1**), alongshore mean of barrier width  $W_b$  (**a2,b2,c2**), and number of inlets (**a3,b3,c3**). All three quantities are shown at years 100, 300 and 500 after model spin-up (first, second, and third columns, respectively) and averaged over five simulations.



**Figure 8.** Drowning timescales: time after spin-up needed to increase  $\Delta F$  by (a) 0.1 or by (b) 0.3 for different significant wave heights  $H_s$  and rate of SLR  $\dot{\xi}$ . Green and red rectangles refer to the situations shown in Figure 5.

481 is also because higher waves tend to import more sediment into the inlets, thereby clos-  
 482 ing them more often when there is no SLR-induced drowning.

483 Depending on the rate of SLR and on the wave height, the barrier starts drown-  
 484 ing (if it does) after a certain time. In all cases, this is not achieved instantly after the  
 485 rate of SLR changes, but there is a time lag for the barrier system to adapt. Situations  
 486 in which  $\Delta F$  attains a value of 0.1 or 0.3 are reached earlier for environments with higher  
 487  $\dot{\xi}$  and intermediate  $H_s$  (Figure 8). The dependence of the time lag on the rate of SLR,  
 488 arises from the gradual evolution of the inlets cross-sectional area. Still, for the same rate  
 489 of SLR, this lag in barrier response depends on the wave height as well, with interme-  
 490 diate wave heights ( $H_s \sim 2$  m) yielding the fastest barrier response. This is because  
 491 intermediate  $H_s$  cause more drowning due to the deepening of the shoreface toe, which  
 492 cannot be counteracted by the sediment imported by waves. For higher waves, sediment  
 493 imported by the littoral drift is able to counteract the effects of the deepening of the shoreface  
 494 toe, and it takes longer for a barrier to drown. For lower waves, even if the sediment im-  
 495 ported by the littoral drift is not so abundant, the toe of the shoreface is shallower, thus  
 496 the whole barrier system can adapt faster to SLR-induced drowning. Interestingly, even  
 497 if most situations deviate from equilibrium (Figure 7), not all of them reach a state of  
 498  $\Delta F = 0.1$  within 500 yrs.

499 Model simulations of barrier drowning are robust. All quantities shown in Figures 6,  
 500 7, 8 have a low standard deviation of the mean compared to their mean values. For  $\Delta F$ ,  
 501 this value takes generally values below 0.05 and only reaches 0.15 in situations where  $\Delta F$   
 502 is of the order of 0.9. The standard deviation of the mean barrier width is always be-  
 503 low 15 m, and generally around 5 m. Lastly, the standard deviation of the mean num-  
 504 ber of inlets is always below 3.

## 505 4 Discussion

### 506 4.1 Comparison with Observations and Earlier Models

507 The cross-shore dynamics reproduced with the BRIE-D model are similar to those  
 508 obtained by the BRIE model of Nienhuis and Lorenzo-Trueba (2019b) or an earlier 2D  
 509 horizontal barrier island model (which did not include inlets) of Lorenzo-Trueba and Ash-  
 510 ton (2014). For example, the width of the subaerial portion of the barrier eventually at-  
 511 tains a constant value that depends on the rate of SLR  $\dot{\xi}$ . Lorenzo-Trueba and Ashton  
 512 (2014) found the same behavior and termed this state as dynamic equilibrium, because

513 the barrier is still migrating landward, but its width does not change. Similarly, the  
 514 increase in barrier drowning found for larger wave heights due to a deepening of the toe  
 515 of the shoreface is in agreement with results of Lorenzo-Trueba and Ashton (2014). Dif-  
 516 ferences include the maximum barrier transgression rate, which can be higher in BRIE  
 517 and BRIE-D, because these also account for sources and sinks of sediment in the bar-  
 518 rier due to inlet and alongshore dynamics.

519 Compared to BRIE, the BRIE-D model computes very high inlet migration rates  
 520 when rates of SLR are higher ( $\sim 5$  km/yr, see Figure 4c). Barrier narrowing due to SLR  
 521 may increase inlet migration rates. However, a compilation of observed inlet migration  
 522 rates shows they are notably lower (e.g., highest observed rate is 700 m/yr, Nienhuis &  
 523 Ashton, 2016) than the values obtained with the BRIE-D model. The BRIE model yields  
 524 more realistic migration rates (of the order of 10 m/yr, see Figure 4d). These differences  
 525 are caused by the new dynamics in alongshore sediment transport deposition within the  
 526 inlet implemented in the BRIE-D model (first term on the right-hand side of Equation 1).  
 527 These new dynamics allow for a disconnection of the updrift and downdrift tips of the  
 528 barrier, and thus inlet widening. However, they also result in, perhaps, unrealistically  
 529 large inlet migration rates. Inlet dynamics in BRIE-D are based on Delft3D simulations  
 530 from Nienhuis and Ashton (2016), who computed the distribution of sediment transport  
 531 between the updrift and downdrift tips of the barrier. However, their experiments were  
 532 performed with barrier widths between 250 m and 800 m and inlet widths lower than  
 533 1 km. Thus, situations with SLR-driven drowning were not included. Future studies should  
 534 investigate how to better parameterize inlet sediment distributions under drowning situ-  
 535 ations in which the barrier becomes narrower, possibly inducing new inlet dynamics.

536 Observations of landward barrier migration back up our results, for example the  
 537 Isles Dernières in Louisiana have been migrating landward at a rate of  $\sim 10$  m/yr, un-  
 538 der a rate of relative SLR of 13 mm/yr (Dingler & Reiss, 1990; Dingler et al., 1993). The  
 539 BRIE-D model yields a landward migration of the order of 2–8 m/yr for rates of SLR  
 540 between 2 and 20 mm/yr. These agreements are expected on longer timescales because  
 541 barriers then follow the basement slopes, which are  $O(10^3)$  m/m.

542 We obtained good agreement with observations in terms of widening of inlets caused  
 543 by barrier drowning. For example, the Isles Dernières have experienced gradual drown-  
 544 ing during the last 200 yrs, in which  $\sim 0.7$  of their subaerial area has been lost under  
 545 a rate of relative SLR of 13 mm/yr (Dingler et al., 1993; Davis Jr. & FitzGerald, 2010).  
 546 These rates are not uncommon. Simulations performed with the BRIE-D model with a  
 547 rate of SLR of 17 mm/yr resulted in a 0.7 increase in  $F$  in 300 years, although under a  
 548 higher rate of SLR than that measured in the Isles Dernières. There could be other mech-  
 549 anisms inducing a significant land loss in the Isles Dernières, such as marsh drowning,  
 550 which are not implemented in the BRIE-D model. These comparisons are challenging  
 551 because of model sensitivities to other factors (e.g., shoreface response rate, maximum  
 552 overwash fluxes) that are difficult to retrieve from field observations.

553 Other ways to compare BRIE-D model results with observations include the frac-  
 554 tion below MSL,  $F$ . In an environment representative of the Wadden Islands (i.e.,  $\dot{\xi} =$   
 555 4 mm/yr,  $H_s = 1.5$  m), the BRIE-D model yields a constant fraction of the barrier ly-  
 556 ing below MSL of  $\sim 0.1$ . This is lower than the observed  $F$  of 0.22 (see Figure 1). Not  
 557 all model parameters have been calibrated for the Wadden Islands, such as marsh cover  
 558 or wave asymmetry, among others. Nevertheless, for higher rates of SLR ( $\dot{\xi} > 18$  mm/yr),  
 559 the BRIE-D model yields an increase in fraction of the barrier lying below MSL up to  
 560 0.6, as SLR-driven inlet widening will dominate the barrier island evolution (see Figure 6).

561 The BRIE-D model is not able to reproduce all the dynamics involved in barrier  
 562 drowning. For example, we have not modeled the curvature of barrier tips occurring in  
 563 wide inlets when bypassing diminishes (Davis Jr. & FitzGerald, 2010). Future research  
 564 should focus on finding appropriate parametrisations for these dynamics and implement-

565 ing them in the BRIE-D model such that the drowning state of a barrier is modeled as  
 566 realistically as possible.

567 Furthermore, the ebb-tidal delta is not explicitly included in the BRIE-D model  
 568 albeit it is a prominent entity in the sand balance of tidal inlets. Nevertheless, its effects  
 569 on inlet migration rate and the size of the flood-tidal delta are implicitly taken into ac-  
 570 count through its effects on waves and currents (Nienhuis & Ashton, 2016). In that sense,  
 571 the BRIE-D model, as well as the BRIE model, offers a different picture on inlet and bar-  
 572 rier dynamics than that in previous studies, such as van de Kreeke (2006).

573 Still, the BRIE-D model is a reasonable tool to understand the different mecha-  
 574 nisms involved in barrier island evolution and, particularly, drowning. Specifically, the  
 575 multiple parametrisations used in the model make it very computationally efficient, al-  
 576 lowing for an in-depth study of the effects of multiple parameters on the response of bar-  
 577 rier systems. More observations are needed to properly evaluate and compare projec-  
 578 tions from BRIE-D, also in comparison with more process-based models such as Mariotti  
 579 (2021).

## 580 4.2 Choice of Parameters

581 The main objective of this study was to gain insight on the different dynamics re-  
 582 lated to barrier drowning rather than mimicking real situations. We have thereby kept  
 583 wave height, tidal range and storm return period constant through the simulations, al-  
 584 beit they are expected to change as  $\dot{\xi}$  increases (Bricheno & Wolf, 2018; Pickering et al.,  
 585 2012). Yet, we have selected their values such that different barrier systems in the world  
 586 are represented by our simulations (Mulhern et al., 2017; Nienhuis & van de Wal, 2021).

587 Furthermore, global projections of SLR may not be representative of barrier sys-  
 588 tems, given the high subsidence rates present in deltas. Consequently, we have chosen  
 589 constant  $\dot{\xi}$  through all simulations. We can thus have a broader range of scenarios and  
 590 we can apply them in longer-term situations. The only drawback of applying constant  
 591 rates of SLR is the abrupt change that the system experiences between the spin-up pe-  
 592 riod and the rest of the simulation, which causes an adaptation period of  $\sim 100$  yrs (see  
 593 Figures 4, 5). Nevertheless, these 100 yrs of adaptation do not seem to alter the grad-  
 594 ual path toward drowning of the barriers. Similarly, Mariotti and Hein (2022) found a  
 595 lag of hundreds of years in barrier response to abrupt changes in rates of SLR. Another  
 596 consequence of the abrupt change in rate of SLR are the irregularities in the backbar-  
 597 rier shoreline of the barrier just after spin-up (see Figure 3a). Still, these irregularities  
 598 are smoothed with time and end up disappearing, hence we do not consider them to be  
 599 a sign of model instability.

600 Simulations performed with an increasing rate of SLR (based on RCP scenarios)  
 601 showed the same tendency as the respective simulations with equivalent constant rates  
 602 of SLR (i.e., RCP2.6 and  $\dot{\xi} = 5$  mm/yr, RCP4.5 and  $\dot{\xi} = 6$  mm/yr, RCP8.5 and  $\dot{\xi} =$   
 603 10 mm/yr; see Figure S8). Future studies could, however, study in further detail the ef-  
 604 fects of a gradual increase in rate of SLR by varying the increase in sea level as well as  
 605 the timescale involved in this gradual evolution.

## 606 5 Conclusions

607 Here we aimed to (1) understand how the fraction of the barrier lying below MSL  
 608 is affected by SLR, (2) examine the temporal evolution of the barrier island drowning,  
 609 as well as quantifying drowning timescales, and (3) explore its dependence on model pa-  
 610 rameters. With our new model (BRIE-D), we have performed simulations with a wide  
 611 range of values for significant wave height  $H_s$  and rate of SLR  $\dot{\xi}$ . Using the model out-  
 612 put, we have quantified barrier island drowning by computing the total fraction of the

613 barrier lying below MSL  $F$ , that caused by tide-wave imbalance  $\Delta F$ , the alongshore mean  
614 of the barrier width, and the number of inlets.

615 Effects of SLR on barrier islands manifest as an increase of inlet width and/or an  
616 increase in inlet number. Barriers drown faster for higher  $\xi$ . Barrier response to changes  
617 in rates of SLR takes place at timescales of the order of hundreds of years, and occurs  
618 in a gradual manner. It takes  $\sim 100$  yrs for a barrier to adapt to a different rate of SLR.  
619 After this period, drowning may occur under high rates of SLR within the following cen-  
620 turies.

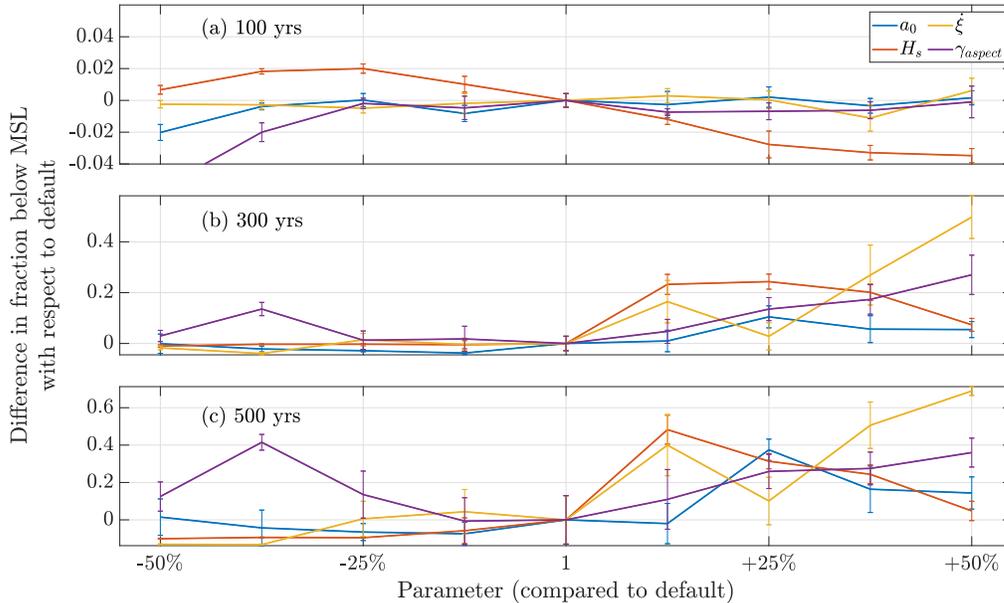
621 We expect environments with intermediate wave heights ( $\sim 2$  m) to be most sensi-  
622 tive to SLR-induced drowning. Lower wave environments have shallower depth of clo-  
623 sure and can respond faster to SLR. Higher waves trigger two opposed mechanisms: a  
624 more frequent inlet closure, and a more intense barrier drowning. The former is caused  
625 by the larger amount of sediment imported into the inlet system, whereas the latter is  
626 a result of the deeper shoreface toe, which makes a barrier system more prone to drown-  
627 ing.

## 628 Appendix A Default model parameters

629 Unless stated otherwise model parameters take their default values, given in Ta-  
630 ble A1.

**Table A1.** Default values of model parameters. Shortened references are as follows: LTA14 (Lorenzo-Trueba & Ashton, 2014), B80 (Bowen, 1980), M17 (Mulhern et al., 2017), AM06 (Ashton & Murray, 2006), SZ09 (de Swart & Zimmerman, 2009), R13 (Roos et al., 2013), N15 (Nienhuis et al., 2015).

Name	Value	Units	Explanation
$\rho_w$	1025	kg m <sup>-3</sup>	Density of water
$\omega$	$1.4 \cdot 10^{-4}$	s <sup>-1</sup>	Offshore tidal radial frequency
$g$	9.81	m s <sup>-2</sup>	Gravitational acceleration
$R$	1.65	–	Submerged specific gravity of sediment
$e_s$	0.01	–	Suspended sediment transport efficiency factor (LTA14)
$c_s$	0.01	–	Friction factor (B80)
$n$	0.05	s m <sup>-1/3</sup>	Manning roughness coefficient
$\dot{\xi}$	10	m yr <sup>-1</sup>	Rate of SLR
$H_s$	1.5	m	Significant wave height in deepwater (M17)
$a_0$	0.8	m	Offshore tidal amplitude (M17)
$T_{storm}$	10	yr	Minimum period between inlet forming storms
$T_p$	10	s	Peak wave period
$a$	0.8	–	Wave asymmetry (AM06)
$h$	0.2	–	Wave highness (AM06)
$\gamma_{aspect}$	0.0707	–	Inlet aspect ratio ( $\gamma_{aspect}^2 = D_{inlet}/W_{inlet}$ )
$u_e$	1	m/s	Tidal inlet equilibrium velocity (SZ09)
$H_{crit}$	2	m	Critical barrier height (LTA14)
$W_{b,crit}$	200	m	Critical barrier width (LTA14)
$Q_{ow,max}$	50	m <sup>3</sup> m <sup>-1</sup> yr <sup>-1</sup>	Maximum overwash flux (LTA14)
$L_{min}$	5	km	Minimum distance between tidal inlets (R13)
$s_{background}$	$10^{-3}$	–	Background slope (LTA14)
$k$	0.06	m <sup>3/5</sup> s <sup>-6/5</sup>	Alongshore sediment transport constant (N15)
$\Delta y$	100	m	Alongshore grid spacing
$\Delta t$	0.05	yr	Time step



**Figure B1.** Differences in fraction below MSL with respect to the default case when varying different morphodynamic parameters at (a) 100 yrs, (b) 300 yrs and (c) 500 yrs after model spin up. Note the different scales in the vertical axis.

## Appendix B Sensitivity analysis

We performed a sensitivity analysis for the main parameters that control the system: tidal range  $a_0$ , significant wave height  $H_s$ , wave period  $T_p$ , rate of SLR  $\dot{\xi}$ , wave asymmetry  $a$ , inlet aspect ratio  $\gamma_{aspect}$ , maximum overwash transport  $Q_{ow,max}$  and the suspended sediment transport efficiency factor  $e_s$ , which controls the shoreface transport. We varied each of the parameters around  $\pm 50\%$  of their default values and computed the alongshore fraction below MSL at three different stages: at years 100, 300 and 500 after model spin-up. For each set of parameters we created five realizations, from which we computed the mean fraction below MSL and the standard error of the mean. We found clear patterns and deviations from the default case for only four of the eight parameters:  $a_0$ ,  $H_s$ ,  $\dot{\xi}$  and  $\gamma_{aspect}$  (see Figure B1). Among these four, largest variations were observed for the significant wave height  $H_s$  and the rate of SLR  $\dot{\xi}$ . Thus, we decided to study the dependence of the model on these parameters in more detail (see Section 3.3).

Increasing the tidal range,  $a_0$ , results in a generally larger fraction below MSL due to a gain in tidal prism, which increases the amount of sediment exported by tidal currents (de Swart & Zimmerman, 2009). Lower tidal ranges cause a lower fraction below MSL  $F$  due to less sediment being exported by tidal currents.

Regarding the significant wave height, we observe two opposite behaviors in terms of  $F$ . Depending on the time after model spin-up, higher waves may produce a decrease or an increase in  $F$ . This is explained by distinguishing two processes caused by high waves: (1) higher waves tend to import more sediment into an inlet, thereby favoring its closure (Ecoffier, 1940), and (2) higher waves affect the sediment at deeper bed levels, causing a larger depth of closure (Houston, 1995). A larger depth of closure means that a larger volume of sand responds to sea level variations, yielding a system that is more prone to drowning. After 100 yrs, an increase in significant wave height causes a decrease in fraction below MSL of up to  $-0.03$ , while a decrease in  $H_s$  increases the fraction below

MSL up to +0.02. This is because at this stage the first mechanism dominates. Nevertheless, after 300 or 500 yrs of model evolution, when the effects of SLR-induced drowning are more prevalent, a decrease in  $H_s$  causes a decrease in the fraction of the barrier lying below MSL. On the other hand, there is a larger fraction of the barrier lying below MSL (up to +0.5) for slightly larger  $H_s$  (+25%). With higher waves (+50%) the fraction corresponds to 0.5 after 500 yrs. The peak in  $F$  attained when increasing  $H_s$  by -25% is due to the second mechanism dominating the barrier evolution. The decrease in  $F$  (even if still higher than in the default situation) when increasing  $H_s$  by +50% is because the first process is more important than in situations with lower waves. The effects of the deepening of the shoreface toe are only visible from year 300 onwards, because the barrier system needs time to adapt and to be affected by SLR.

Increasing the rate of SLR results in more drowning, inducing a fraction below MSL of up to +0.68 by the year 500 for the most extreme case. Note that as explained for significant wave heights, effects of drowning are only visible from year 300 onwards. In contrast, decreasing the rate of SLR decreases the fraction of the barrier lying below MSL up to -0.13 because there is less drowning.

An increase in inlet aspect ratio creates narrower inlets for the same cross-sectional area, thereby yielding a slightly lower fraction below MSL at year 100 (up to -0.01). However, in the years 300 and 500 an increase in inlet aspect ratio results in the opposite effect, yielding an increase in  $F$  of up to +0.36. Lowering the inlet aspect ratio makes shallower inlets, increasing the bottom friction. This causes the inlets to be more susceptible to closing, decreasing thus the fraction of the barrier lying below MSL at the year 100 up to -0.05. However, at 300 or 500 yrs after model spin-up,  $F$  increases for lower values of the inlet aspect ratio. These differences in behavior between earlier and latter times suggest that the dependence of the barrier evolution on the inlet aspect ratio is susceptible to SLR-driven drowning, similarly to the situation obtained when varying  $H_s$ .

## Open Research

The code for the BRIE-D model is accessible from <https://doi.org/10.5281/zenodo.7353693>.

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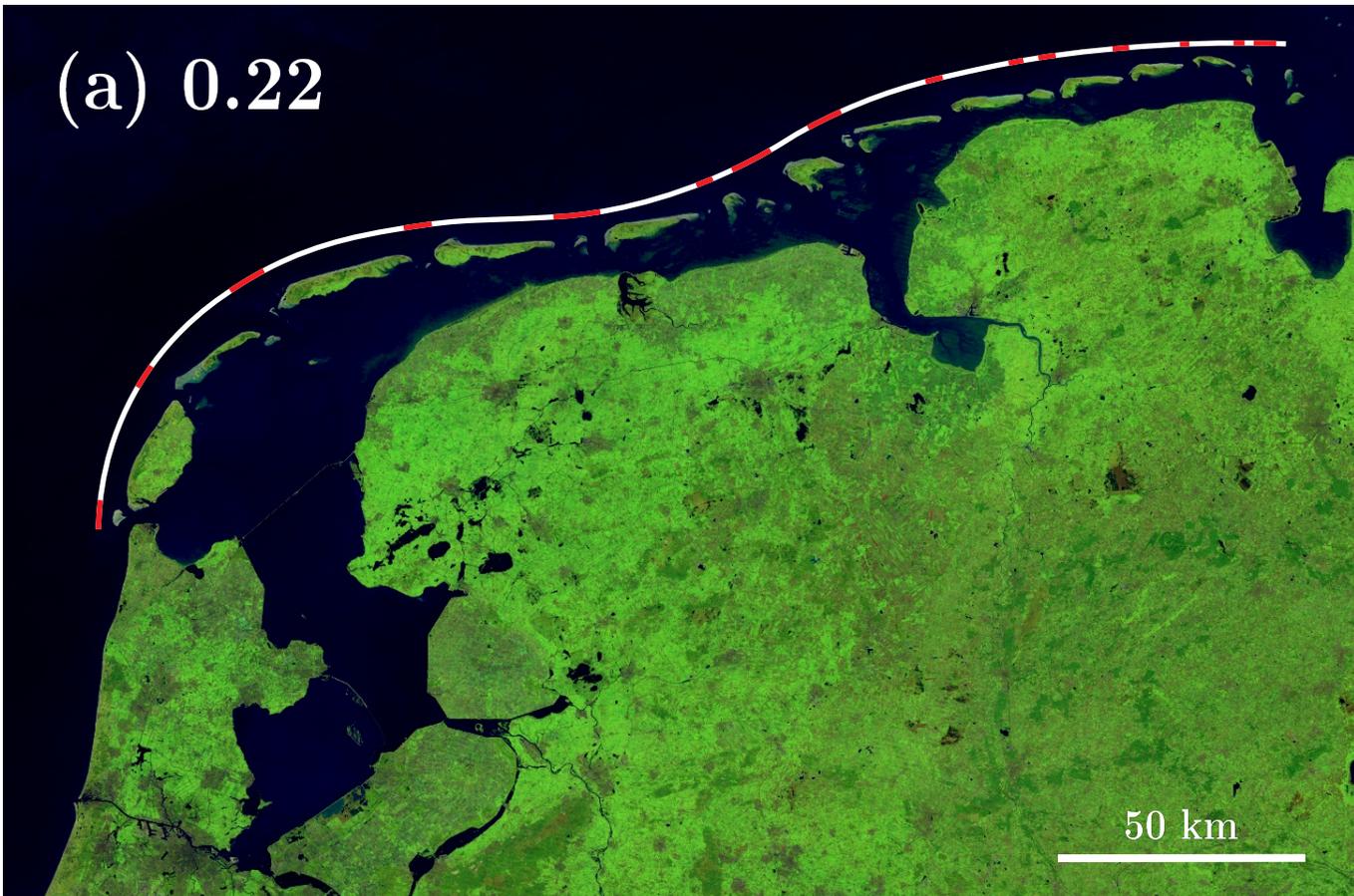
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**Figure 1.**

(a) 0.22



(b) 0.03

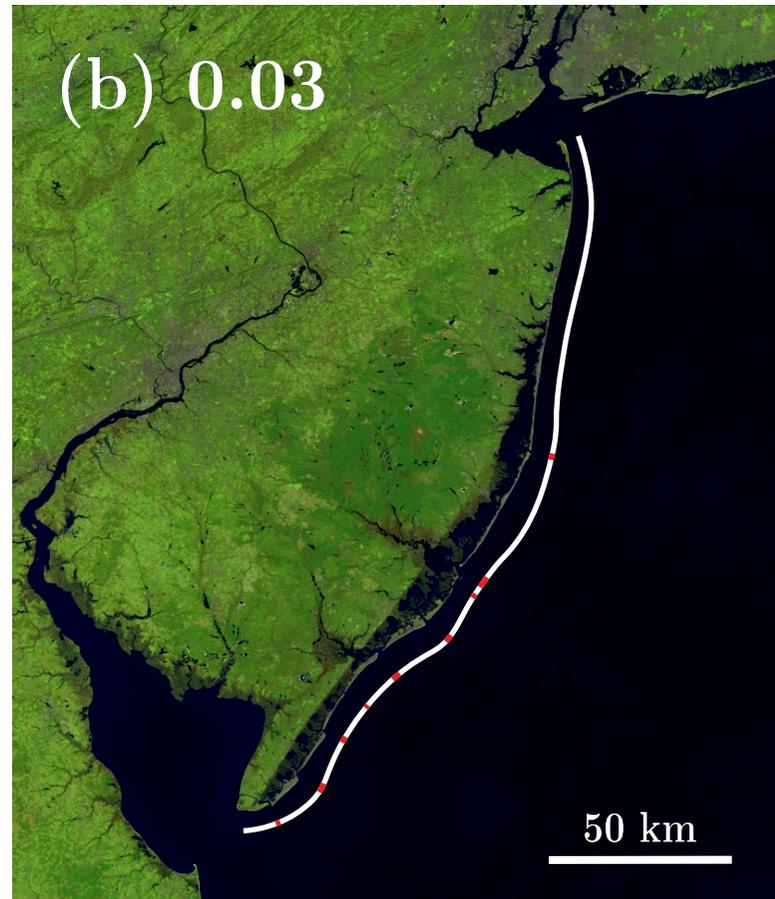


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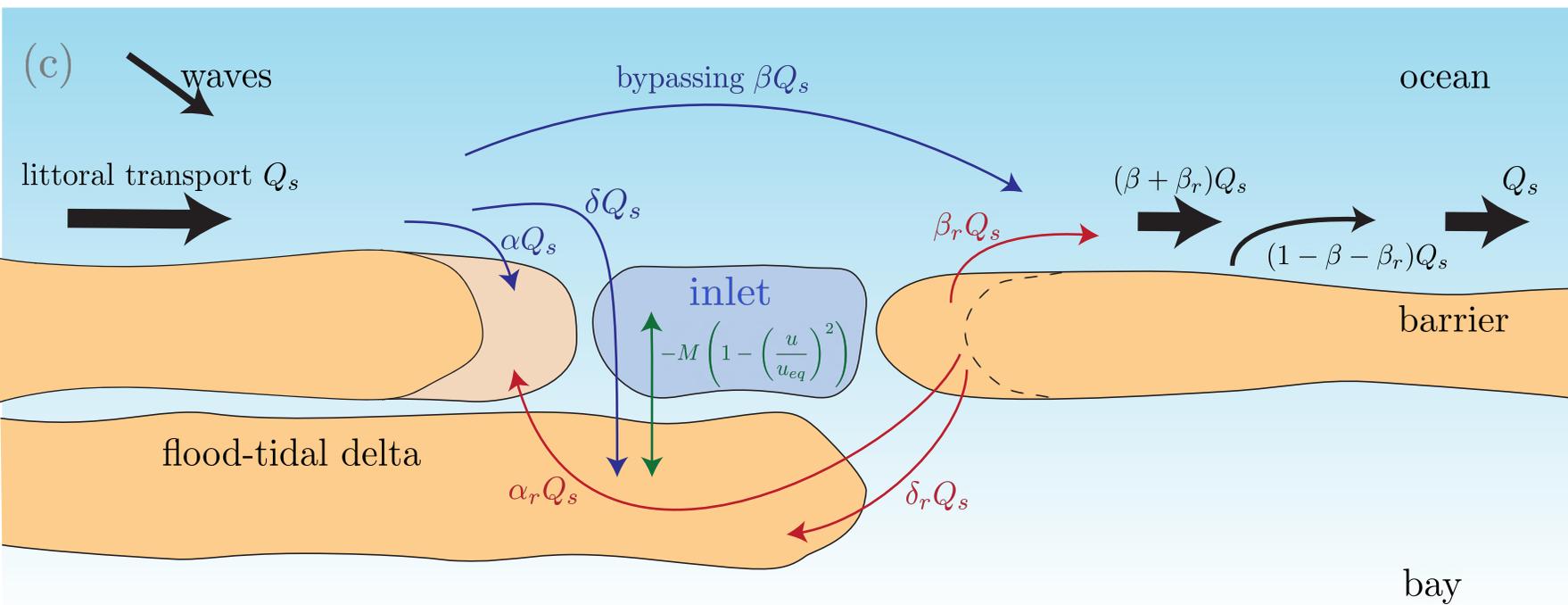
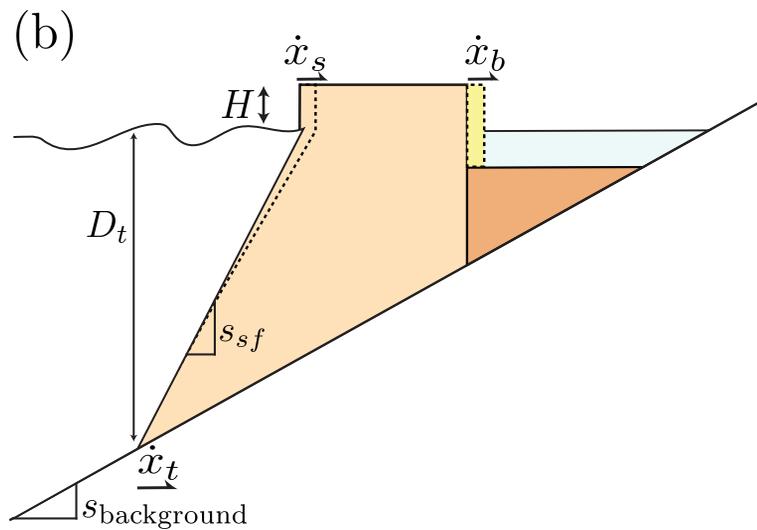
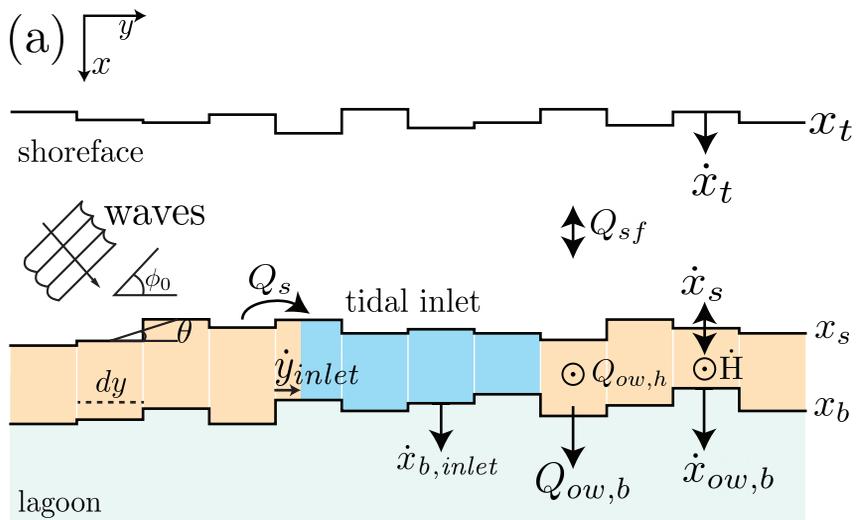


Figure 3.



Figure 4.

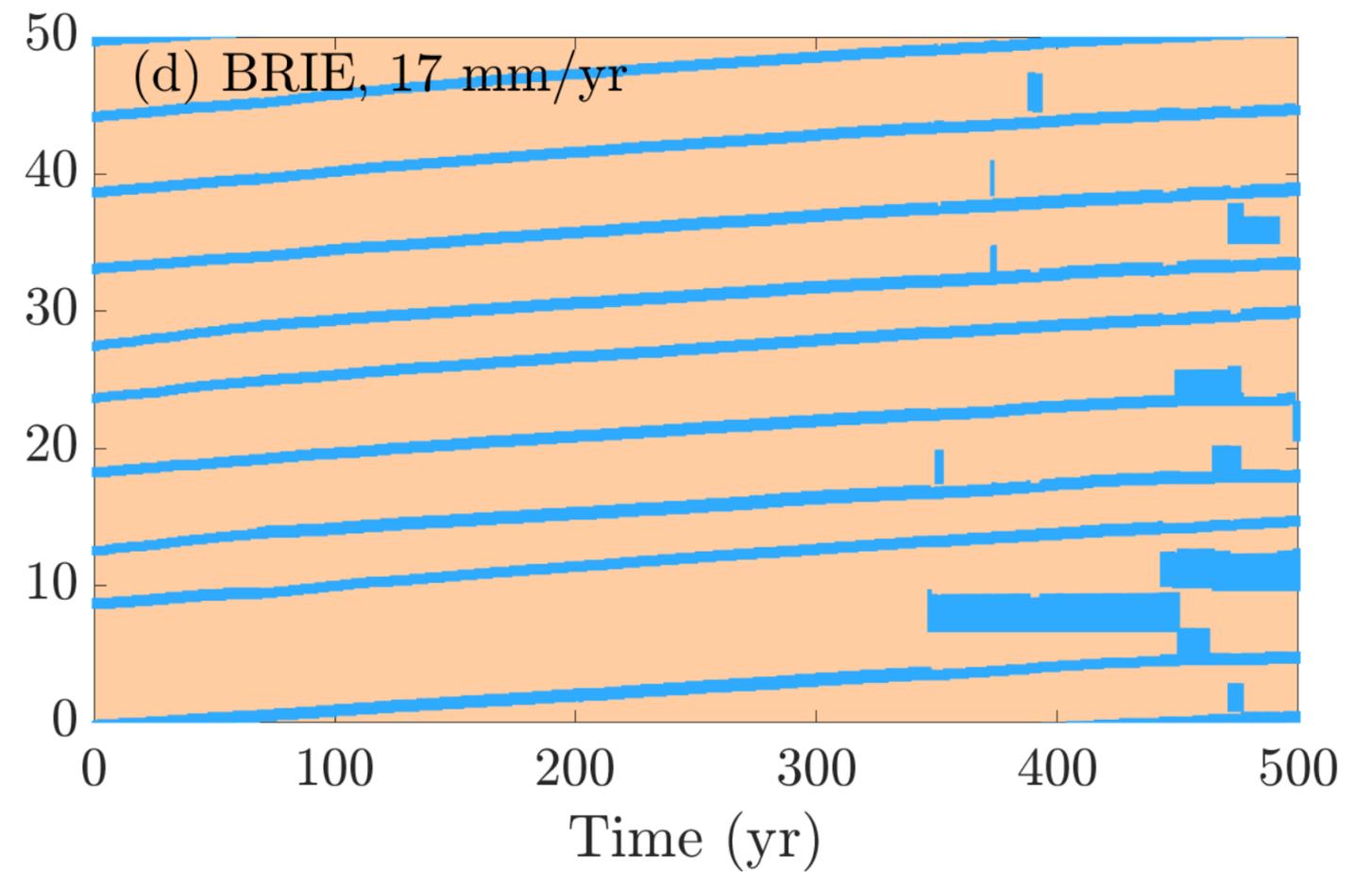
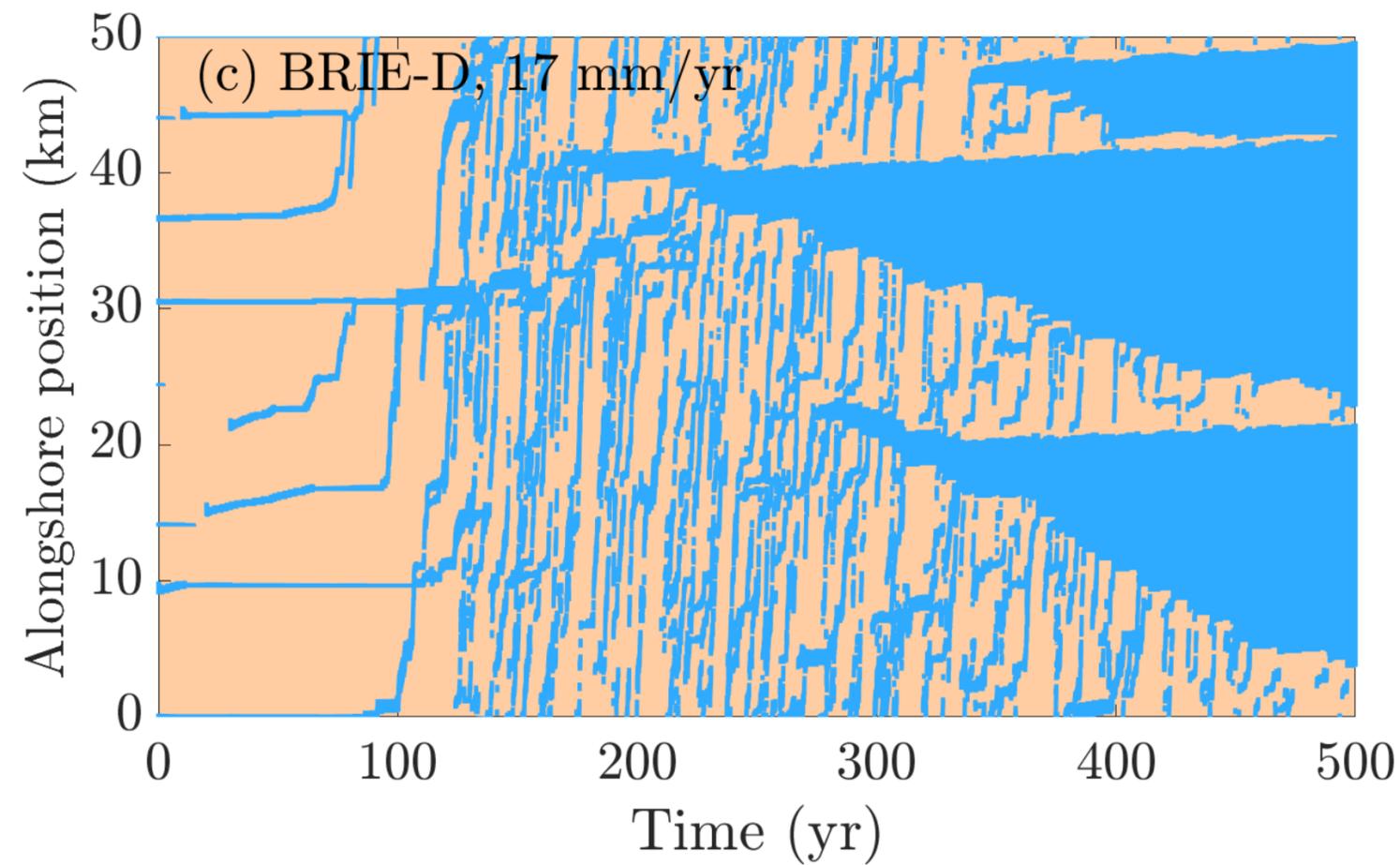
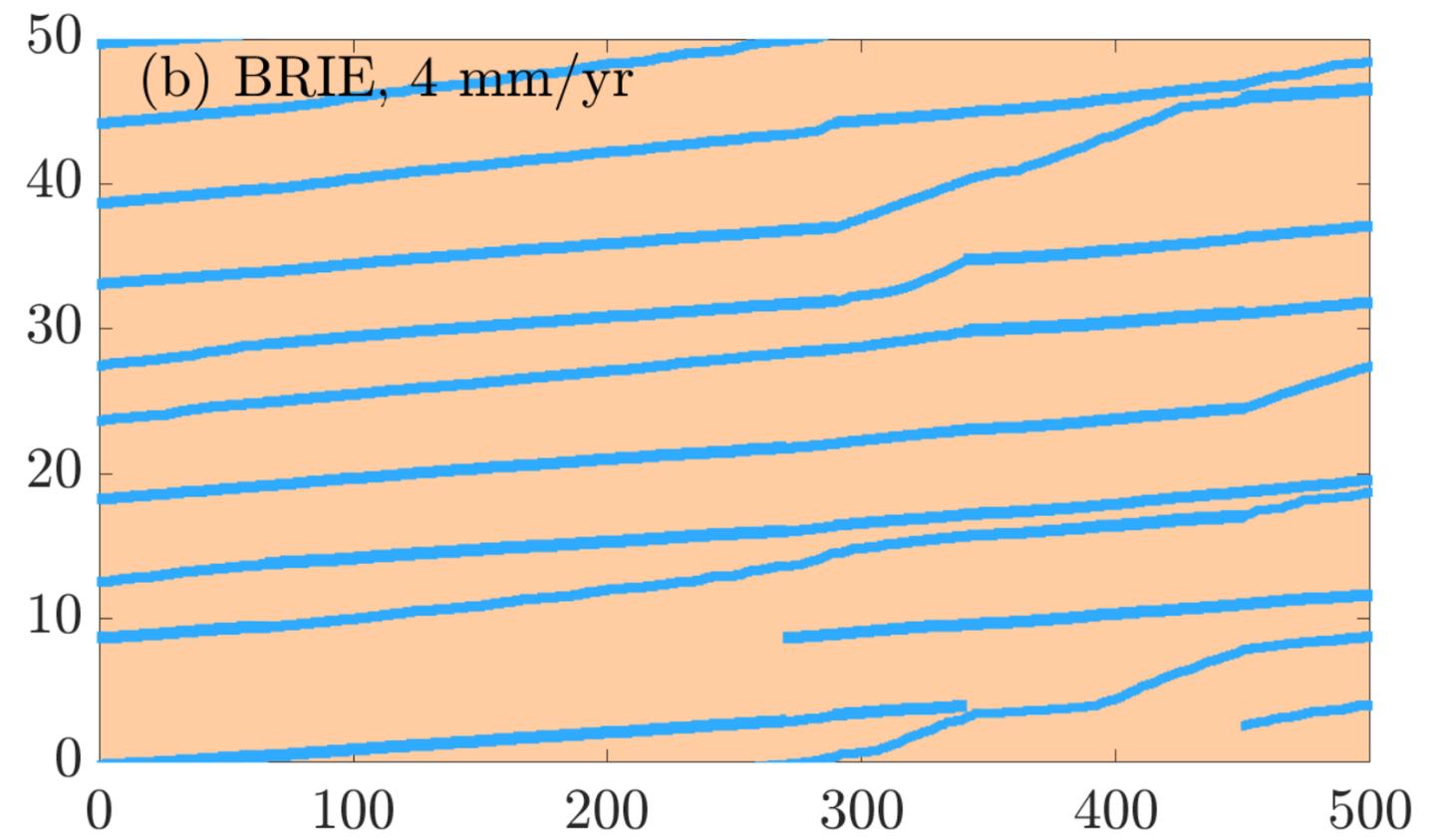
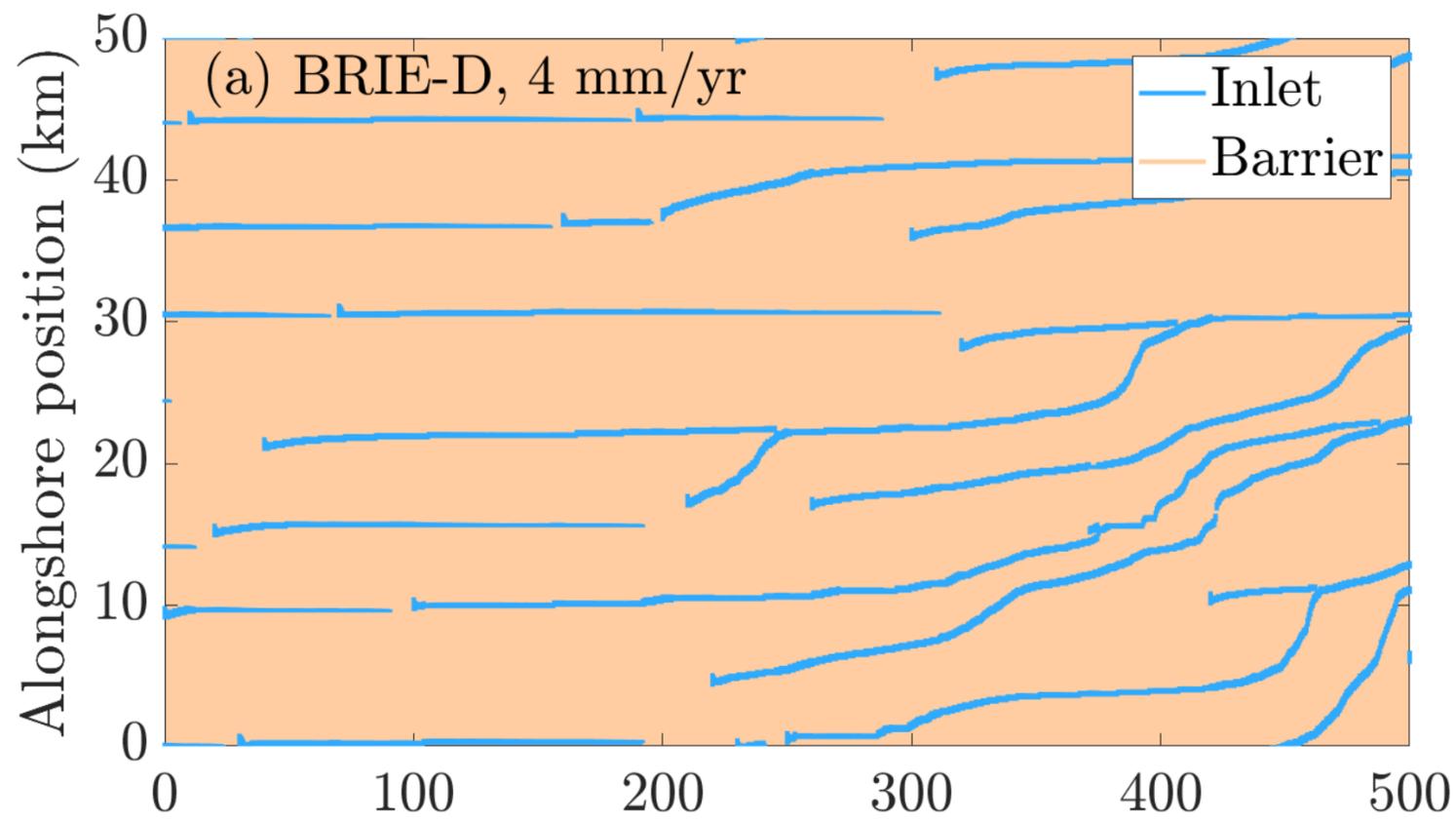


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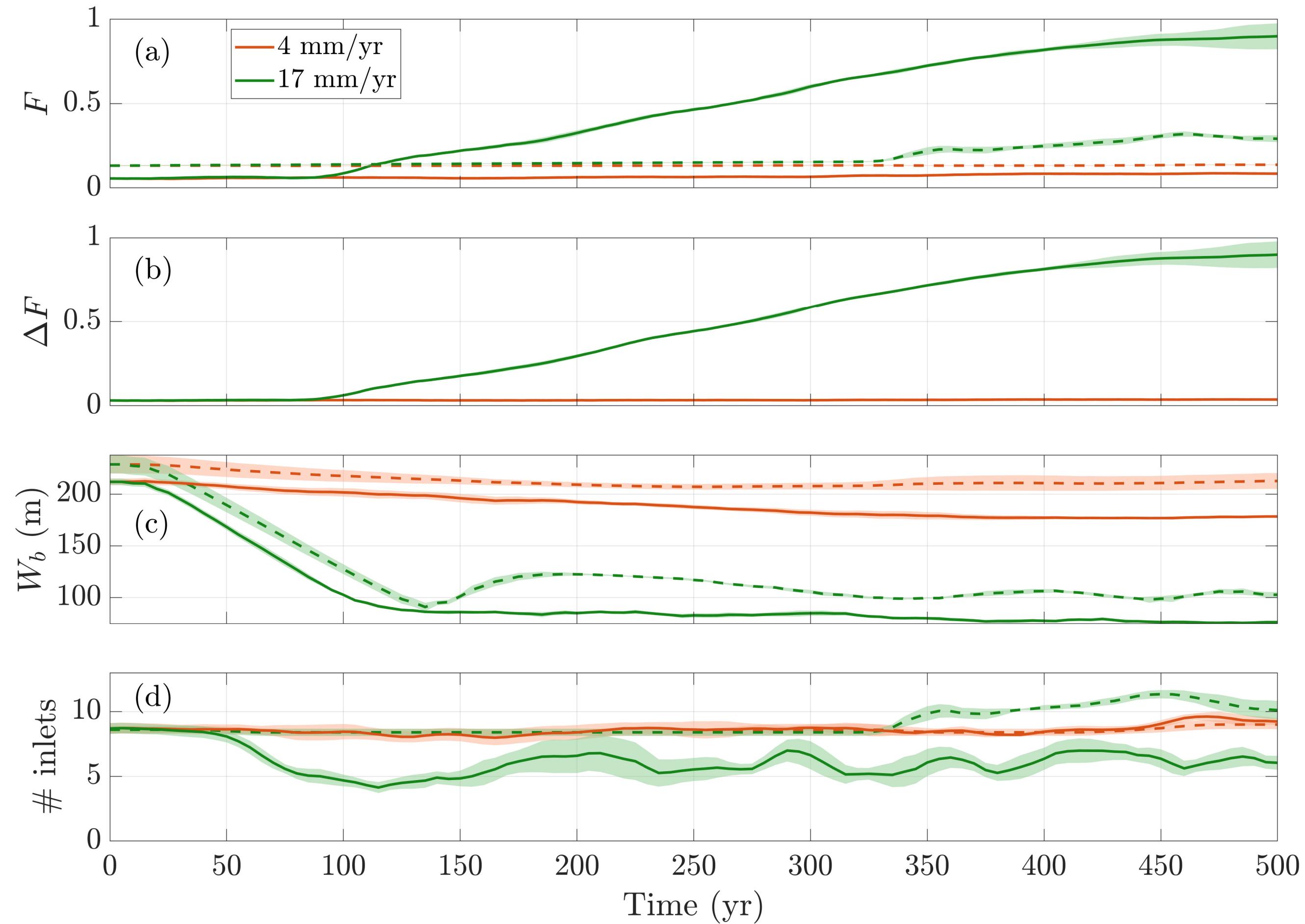


Figure 6.

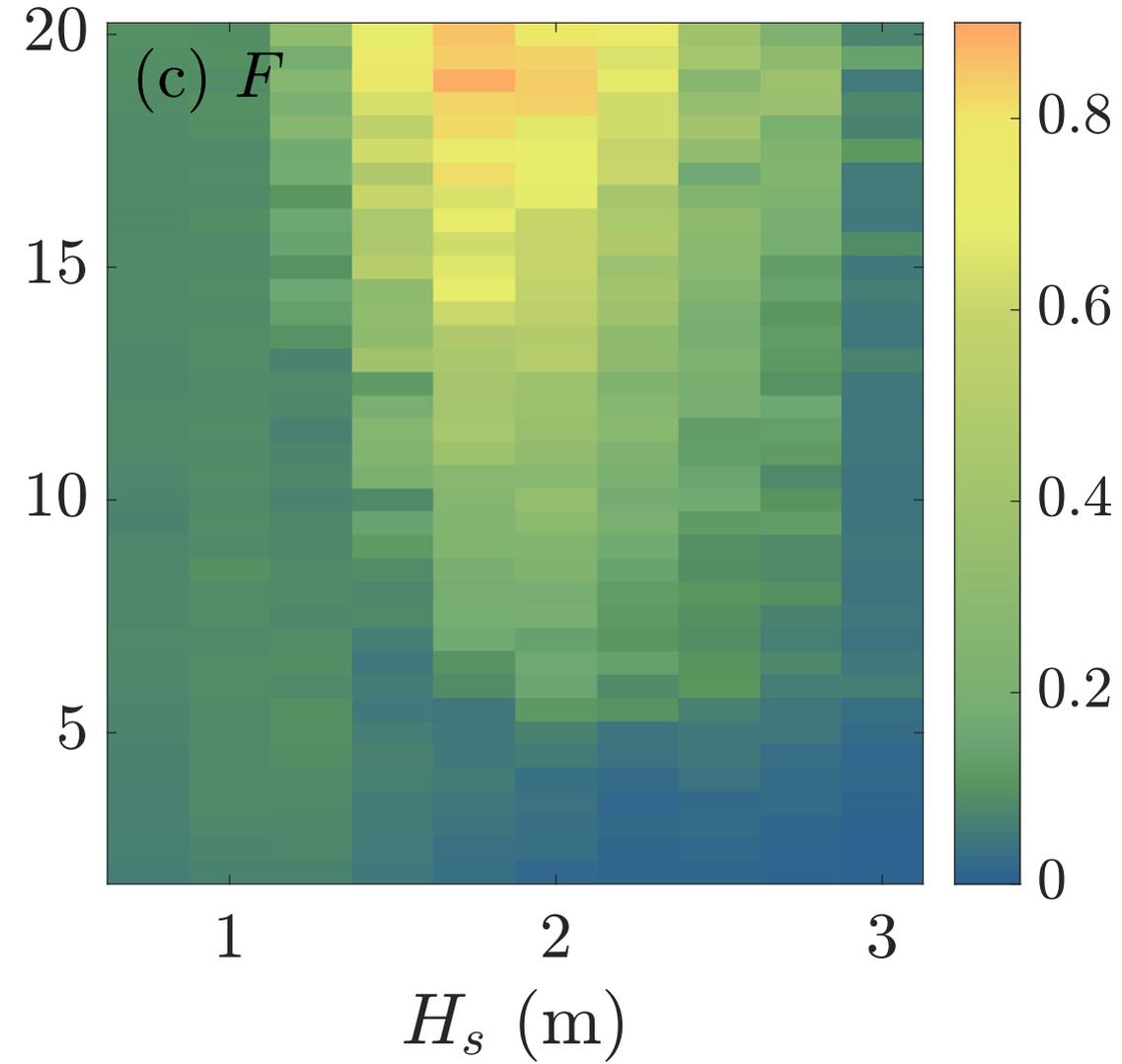
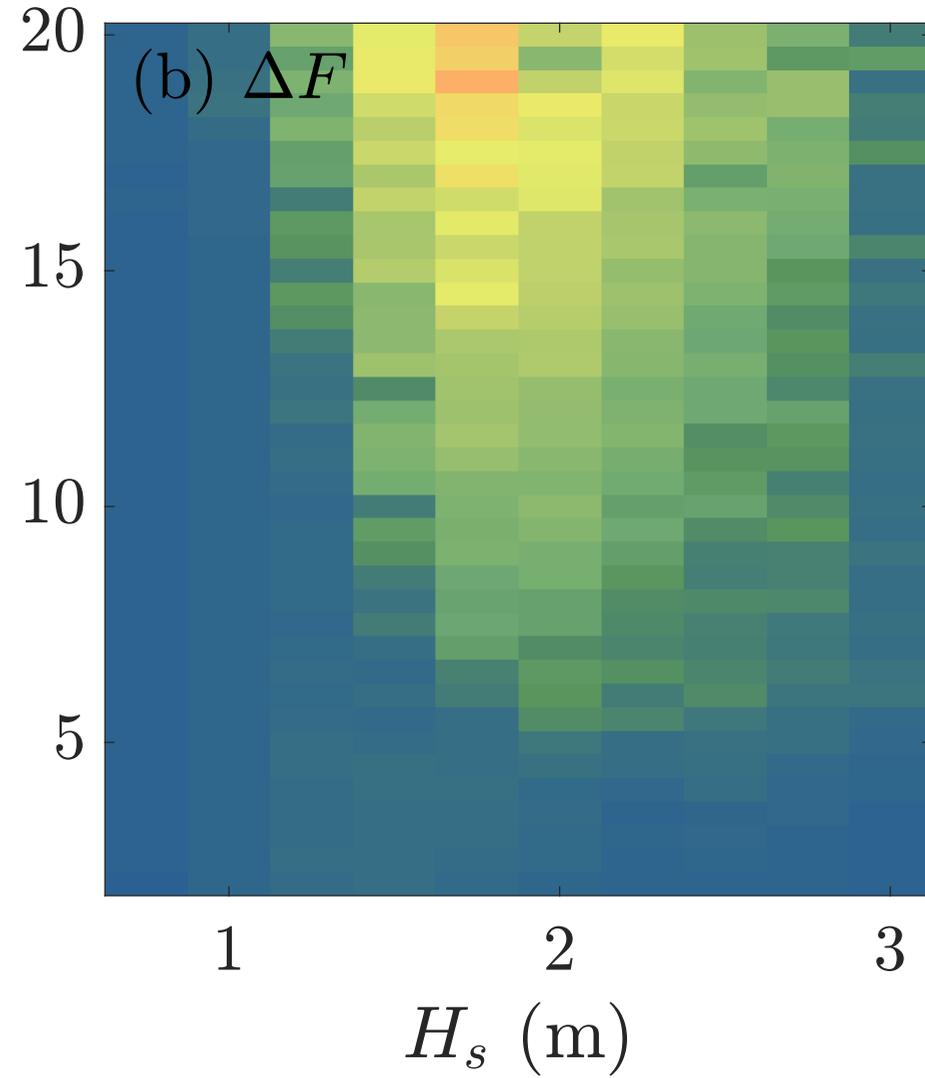
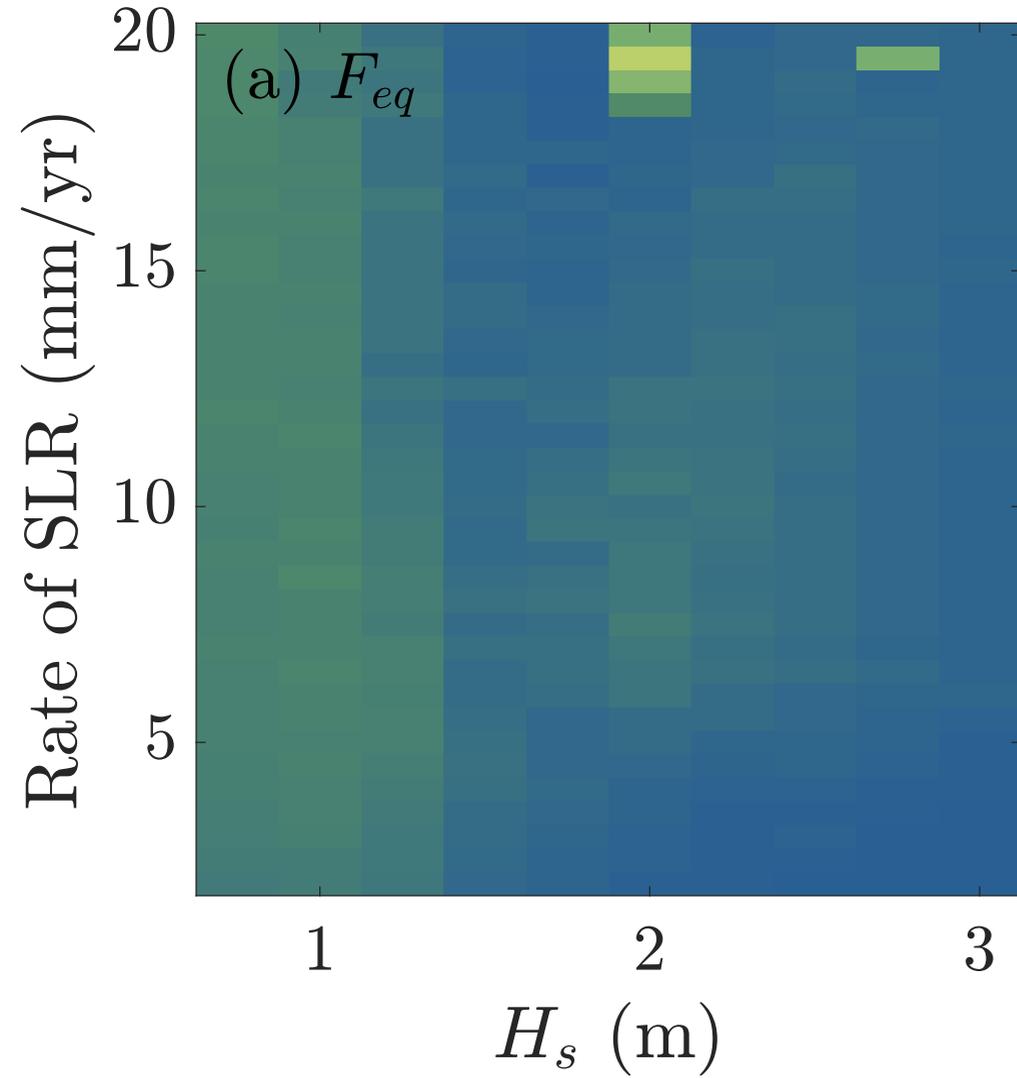


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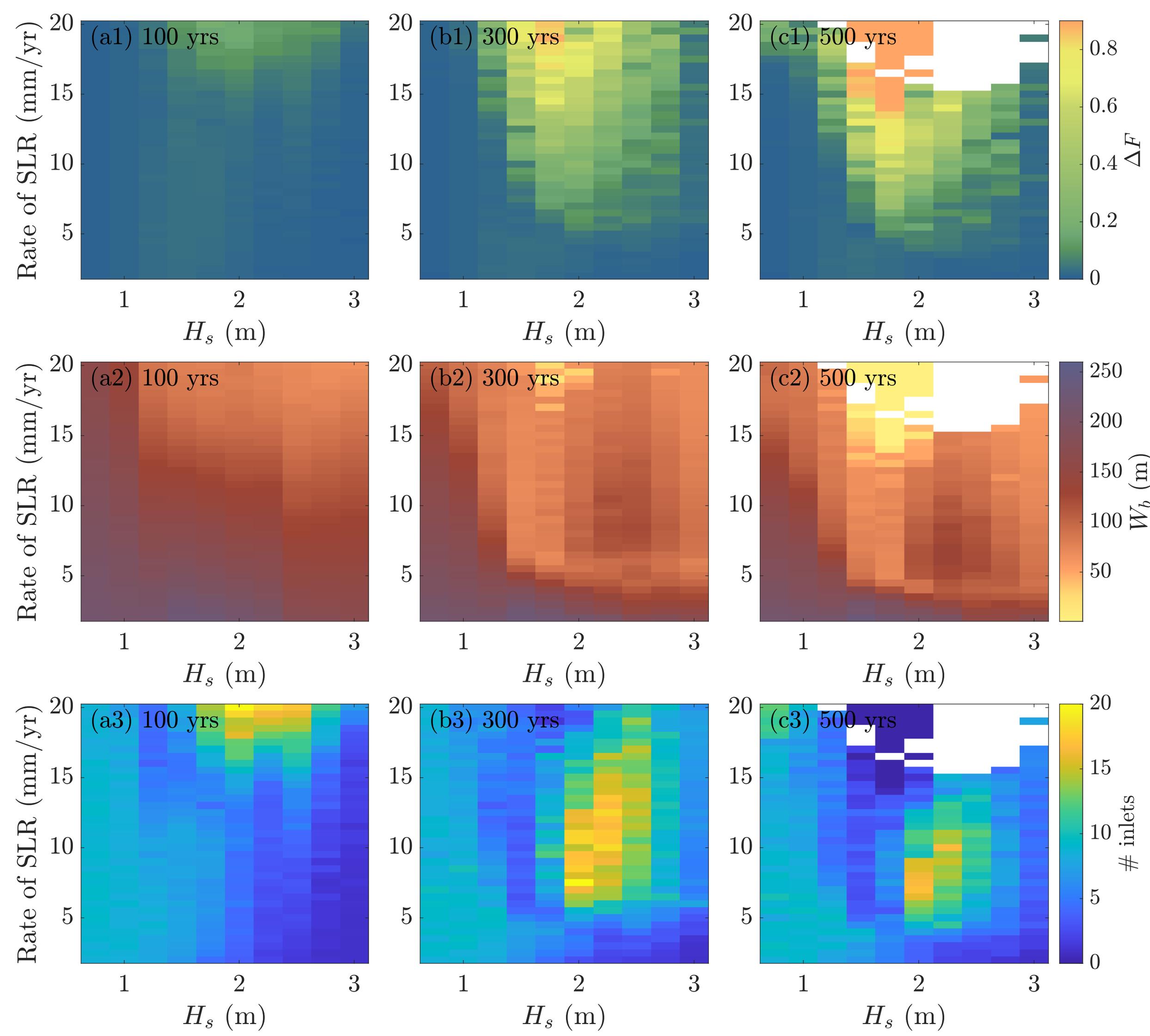


Figure 8.

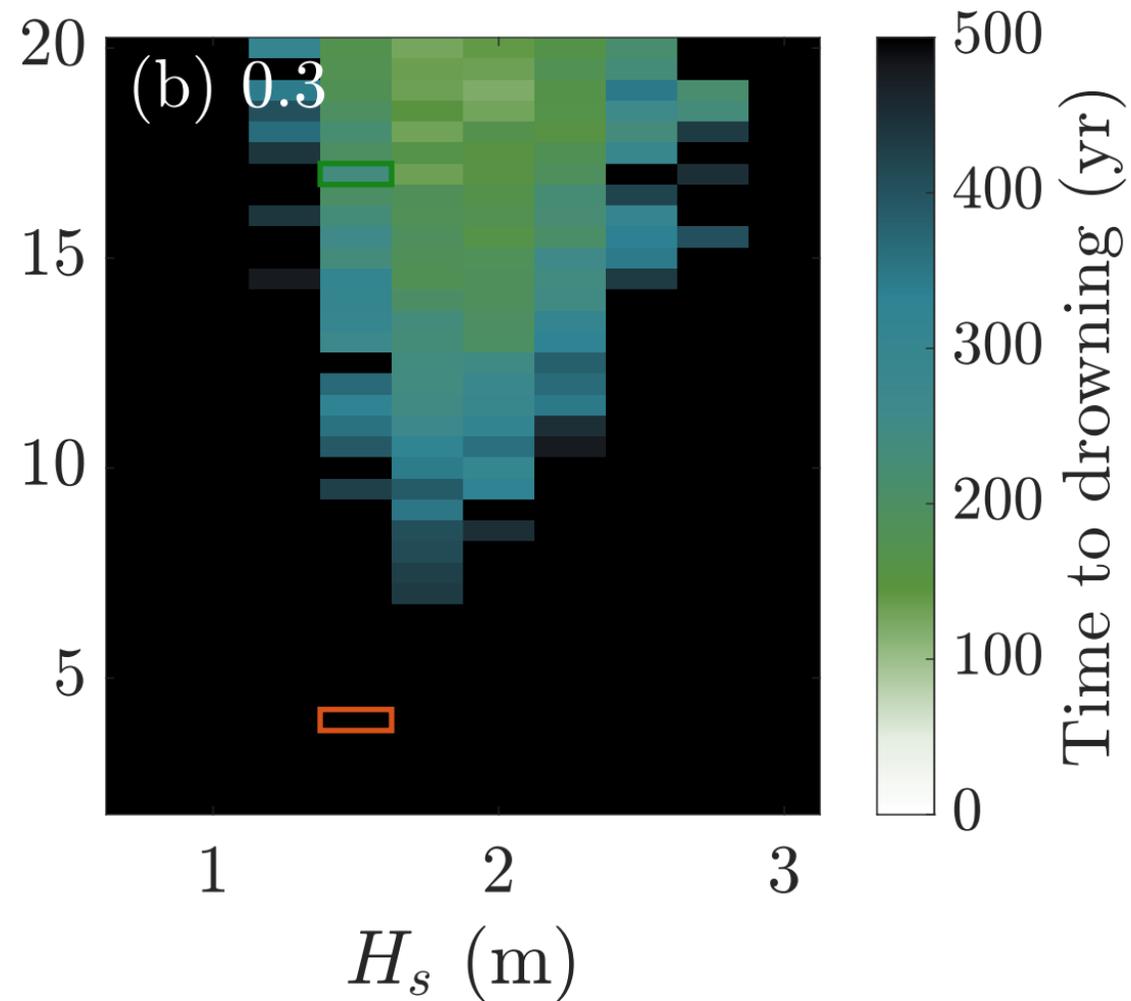
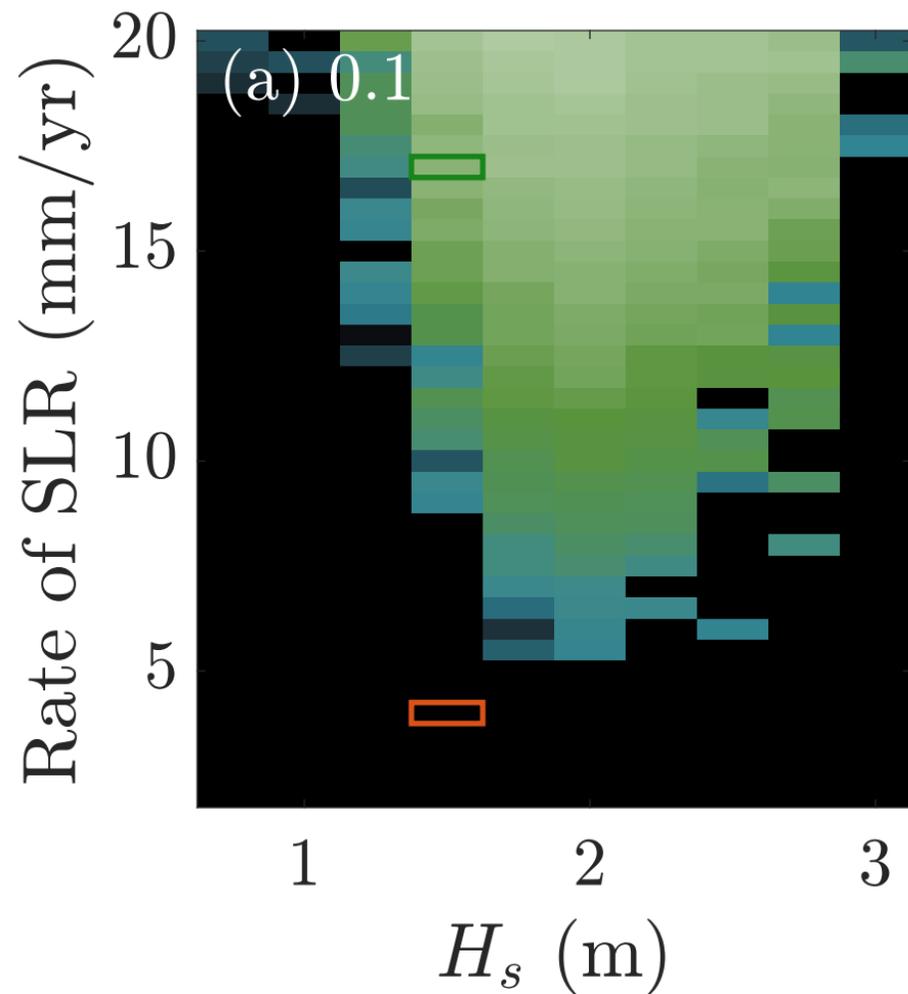


Figure 9.

Difference in fraction below MSL  
with respect to default

