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12 **ABSTRACT:** We demonstrate the use of information theory metrics, Shannon entropy and mutual  
13 information, for measuring internal and forced variability in ensemble atmosphere, ocean, or  
14 climate models. This metric delineates intrinsic and extrinsic variability reliably in a wider range  
15 of circumstances. Information entropy quantifies variability by the size of the visited probability  
16 distribution, as opposed to variance that measures only its second moment. Shannon entropy and  
17 mutual information manage correlated fields, apply to any data, and are insensitive to outliers and a  
18 change of units or scale. In the first part of this article, we use climate model ensembles to illustrate  
19 an example featuring a highly skewed probability distribution (Arctic sea surface temperature) to  
20 show that the new metric is robust even under sharp nonlinear behavior (freezing point). We apply  
21 these two metrics to quantify internal vs forced variability in (1) idealized Gaussian and uniformly  
22 distributed data, (2) an initial condition ensemble of a realistic coastal ocean model (OSOM),  
23 (3) the GFDL-ESM2M climate model large ensemble. Each case illustrates the advantages of  
24 information theory metrics over variance-based metrics. Our chosen metric can be applied to any  
25 ensemble of models where intrinsic and extrinsic factors compete to control variability and can be  
26 applied regardless of if the ensemble spread is Gaussian. In the second part of this article, mutual  
27 information and Shannon entropy are used to quantify the impact of different boundary forcing in  
28 a coastal ocean model. Information theory is useful as it enables ranking the potential impacts  
29 of improving boundary and forcing conditions across multiple predicted variables with different  
30 dimensions.

## 31 **Plain Language Summary**

32 It is important in climate and environmental modeling to distinguish variability that is caused  
33 by external forces versus variability that arises from within the system being modeled itself. In  
34 this paper, we study an ensemble of coastal ocean models, that are forced with tides, winds, and  
35 offshore and atmospheric conditions and an ensemble of climate model simulations that are forced  
36 by greenhouse gases and solar warming. Here, we propose to use information theory—a way to  
37 count the number of physical states visited by a system under study—to quantify the amount of  
38 variability in these models that results from the external forcing versus from the internal forcing.  
39 In this way, we can prioritize improvements or inclusion of the different forcings based on how  
40 large the model response to them is.

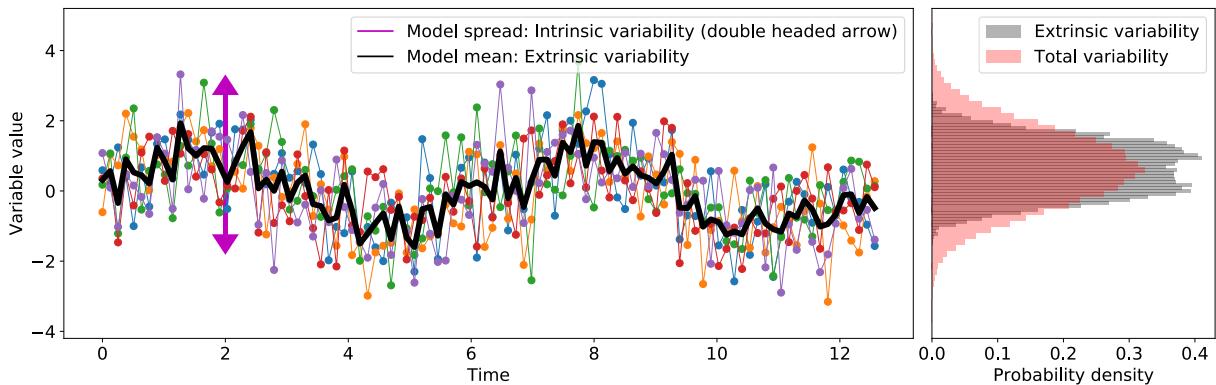
## 41 **1. Introduction**

42 In an ocean or climate model, it is pertinent to understand the cause of variability as it leads to  
43 implications for predictability, prioritization of data collections for assimilation, and provides an  
44 understanding of the dynamics at play in different regions. In a coastal model, variability can arise  
45 from extrinsic factors such as wind forcing, solar and thermal forcing, tides, rivers, evaporation and  
46 precipitation, or it can be due to internal chaos inherent to the governing fluid equations (Sane et al.  
47 2021). In a climate model, modes of variability such as El Niño, the North Atlantic Oscillation,  
48 or the Southern Annular Mode, can conceal or delay the emergence of attributable anthropogenic  
49 climate change signals (Milinski et al. 2019). In high-resolution ocean models, internal chaos or  
50 intrinsic variability can also be due to eddies (Leroux et al. 2018; Llovel et al. 2018). Accurately  
51 quantifying the relative contribution of external and internal factors can help in elucidating the  
52 causes responsible for the observed variability in models, help to identify key observable metrics,  
53 and help quantify concepts such as the time of emergence of climate signals (Hawkins and Sutton  
54 2012).

55 Numerous methods exist in the literature to quantify intrinsic and extrinsic variability using  
56 models or observations (e.g., Frankcombe et al. 2015; Schurer et al. 2013; Liang et al. 2020).  
57 Model ensembles—i.e., a set of simulations sharing the same forcing—naturally vary because each  
58 ensemble member follows the same governing equation (with same external forcings) with identical  
59 or similar parameterizations but differ due to intrinsic chaos. Two types of model ensembles are

60 common: initial condition ensembles (where the same model is used repeatedly with perturbed  
 61 initial conditions and intrinsic variability occurs via chaos), and multi-model ensembles (where a  
 62 variety of models differing in numerics and parameterizations are used to simulate change under  
 63 the same forcing—in this case “intrinsic” variability also includes aspects of model formulations).  
 64 Most of the discussion here will focus on initial condition ensembles, but the metrics proposed can  
 65 be adapted to both cases.

66 To help visualize variability, a generic output from an ocean or atmospheric model is shown in  
 67 Figure 1. Each color represents a different ensemble member and the black solid line is the mean of  
 68 those members. The black solid line is the signal mostly due to extrinsic factors (aside from finite  
 69 ensemble size limits) and the model spread (schematized by the double-headed magenta arrow in  
 70 Figure 1) can be considered due to intrinsic variability or internal chaos.



71 FIG. 1. A sketch of a typical ocean or climate model output for an arbitrary variable. Each ensemble is shown  
 72 in different color and the mean of the ensemble is shown as black line. The ensemble mean can be considered  
 73 to be the trend set by external forcings. The model spread shown by double headed magenta arrow indicates the  
 74 model chaos.

75 One method of quantifying intrinsic and extrinsic variability is to look at variances (second  
 76 central statistical moment) of model spread and model mean (Leroux et al. 2018; Llovel et al.  
 77 2018; Waldman et al. 2018; Yettella et al. 2018). Variance is sufficient to constrain all metrics of  
 78 variability about the mean when distributions are Gaussian and uncorrelated, but a single statistical  
 79 moment usually measures only part of a more complex variability. Many climatological variables  
 80 show non-Gaussian distributions (e.g., Franzke et al. 2020). In fact, generalized variance might  
 81 be misleading (e.g., Kowal 1971). Quantification of variability should be robust to or have known

82 dependence on changes in the units of the quantity or the scale (e.g., changing temperature from  
83 Celsius to Fahrenheit or Kelvin). Comparative metrics, such as intrinsic vs. extrinsic variability  
84 should not depend on these arbitrary choices of units at all.

85 Variability, in essence, is a function of the number of occurrences or frequency of occurrence  
86 (or probability  $p_i$  as a fraction over all visited system states) after appropriately binning the data  
87 (and thereby making the estimated and visited number of states finite rather than continuous).  
88 Information entropy metrics measure variability by taking into account the probability distribution  
89 of the binned data, drawing on the statistical mechanics concept of entropy in quantifying the  
90 number of microstates that a variable can occupy. The fundamental measure in information theory  
91 is the Shannon (1948) or information entropy which characterizes the amount of variability in a  
92 variable (Carcassi et al. 2019). The mutual information, another metric introduced by Shannon  
93 (1948), measures how much information one variable contains about another variable.

94 Information theory is applied in signal processing, computer science, statistical mechanics,  
95 quantum mechanics, etc. It is used to quantify amount of information, disorder, freedom, or lack  
96 of freedom (Brissaud 2005). The application of these abstract notions to geophysical flows can  
97 have immense practical benefit when information entropy is interpreted as a measure of variability,  
98 as entropy does not rely on any particular parametric probability distribution. Metrics from  
99 information theory are not new to climate sciences. They have been introduced in predictability  
100 studies, evaluating the skill of statistical models, as well as uncertainty studies (e.g., Leung and  
101 North 1990; Schneider and Griffies 1999; Kleeman 2002; DelSole and Tippett 2007; Majda and  
102 Gershgorin 2010; Stevenson et al. 2013) and recently in studying variability (Gomez 2020) and  
103 coastal predictability (Sane et al. 2021).

104 In this article we bring well-established concepts of information theory to the particular applica-  
105 tion of measuring intrinsic and extrinsic variability for ensemble model runs within atmospheric  
106 and oceanographic modeling. Our metric uses Shannon entropy and mutual information. We in-  
107 directly employ conditional entropy, which depends on Shannon entropy and mutual information.  
108 To keep the metric intuitive, we have used Shannon entropy and mutual information and not cast  
109 it using conditional entropy.

110 There are two parts to this article. In Part 1, we apply our metric to three sets of data: 1.  
111 Idealized Gaussian and uniformly distributed arrays with specified correlation 2. Ensemble output

112 of a regional coastal model (OSOM) (Sane et al. 2021) where most variables are non-Gaussian.  
113 Ensemble data for the duration of July-August of 2006 has been used. 3. The GFDL-ESM2M  
114 Large Ensemble (Rodgers et al. 2015; Deser et al. 2020), hereby referred to as GFDL-LE. This  
115 dataset contains historical and RCP 8.5 simulation data. All the monthly mean data from 1950 to  
116 2100 have been used in the analysis.

117

118 In Part 2, we use OSOM to demonstrate the use of Shannon entropy and mutual information in  
119 evaluating the effects of altered boundary forcings. In coastal and estuarine systems, it is relevant  
120 to know which forcings are dominant which could potentially lead to prioritizing data collection to  
121 improve accuracy of the forcings. For example, is wind forcing dominant over river forcing, does  
122 using temporal averaged river runoff cause any appreciable changes in the estuarine circulation, or  
123 does change in the wind product alter circulation? These questions can be tackled by switching on  
124 and off or modifying each forcing and comparing the predicted variables using information theory.

125 Recent theoretical advances in understanding dynamical systems through the lens of information  
126 theory relate causality analysis and information transfer (e.g., Liang 2014). Although important,  
127 the transfer of such theoretical concepts into pragmatic research applications are few. Even basic  
128 concepts of information theory (Shannon entropy and mutual information) have been adopted  
129 in a limited capacity by the oceanic and atmospheric community to address problems arising  
130 in predictability and variability. We attempt to bridge the gap using approximate but practical  
131 framework which can be easily replicated and improved upon in the future, including causality  
132 analysis and the evolution of entropy within modeling systems like those studied here.

### 133 *a. Information theory*

134 We will introduce information theory concisely assuming the reader has no background  
135 knowledge—this section contains standard definitions. Consider a probability distribution  $p_i$  ob-  
136 tained after binning data into  $N$  bins. The user chooses the appropriate number of bins or bin  
137 widths for the range of data. Shannon (1948) identified the average information content in  $N$   
138 possible outcomes, equally or not equally likely, as given by:

$$H = \sum_{i=1}^N p_i \log_2(1/p_i), \quad (1)$$

139 where  $H$  is the Shannon entropy with unit of bits when  $\log$  is base 2 and  $p_i$  is the probability of the  
 140  $i^{th}$  outcome. The factor  $\log_2(1/p_i)$  measures the information of the  $i^{th}$  outcome as proposed by  
 141 Hartley (1928) and is also a measure of uncertainty (Cover 1999), as it measures the information  
 142 gained by knowing that the  $i^{th}$  outcome has happened or equivalently that the variable falls in the  
 143  $i^{th}$  bin. The term information does not mean knowledge but it means the amount of uncertainty  
 144 shown by a variable or the freedom that a variable has in visiting different combinations of the  $N$   
 145 bins. Shannon (1948) found Equation 1 to provide the average information (or uncertainty) for  
 146 all events in a record. For the entire set of elements, a highly probable event has less uncertainty  
 147 associated with it and low probability event has high uncertainty associated with it. The prefactor  
 148  $p_i$  is thus used to weight the information over all possibilities. One way to interpret the need for  
 149 the prefactor  $p_i$  is that in repeated experiments the events with higher probability will occur more  
 150 often, hence they should contribute more to a quantification of variability than infrequent events.

151 Stone (2015) gives an intuitive way of understanding Shannon entropy using a binary tree.  
 152 A binary tree is a tree chart which starts with one node and splits to two nodes at each node.  
 153 At each node you can take a left or right turn to proceed and if there are say 3 levels in the  
 154 tree, then 8 (i.e.  $2^3$ ) outcomes or possible destinations exist. If a binary tree has  $N$  equally  
 155 probable outcomes then the set of instructions required to reach the correct destination is given by  
 156  $h = (N)(1/N) \log_2(N) = \log_2(N)$ . The *uncertainty* about reaching the correct destination will be  
 157 removed by providing  $\log_2(N)$  bits of information. In other words, if entropy is  $h$  then  $2^h$  states  
 158 are possible.

159 A second metric from Shannon (1948) which is also extensively used is now known as *mutual*  
 160 *information*. The mutual information between two signals  $x$  and  $y$  denoted by  $I(X;Y)$  is (Cover  
 161 1999)

$$I = \sum_{j=1}^N \sum_{i=1}^N p_{ij} \log_2 \left( \frac{p_{ij}}{p_i p_j} \right), \quad (2)$$

162 where  $p_{ij}$  is joint probability of  $i^{th}$  outcome of  $x$  and  $j^{th}$  outcome of  $y$ . The marginal probability of  
 163  $i^{th}$  and  $j^{th}$  outcomes of  $x$  and  $y$  respectively are  $p_i$  and  $p_j$ . The addend within the summations can  
 164 be expanded to  $p_{ij} (\log_2(p_{ij}) - \log_2(p_i) - \log_2(p_j))$ .  $I$  can be interpreted as the extra information  
 165 in entropy of marginal distributions of  $x$  and  $y$  over the joint distribution. Mutual information is  
 166 symmetric between  $x$  and  $y$  and is the measure of how much information they share. For example,

167 if the distributions are statistically independent, then  $p_{ij} = p_i p_j$  and thus  $I = 0$ . If the two records  
168  $x$  and  $y$  are identical, then  $p_{ij} = p_i = p_j$  and  $I = H$ .  $I$  is the average reduction in uncertainty in  $x$   
169 from knowing  $y$  or vice versa and denotes how much information is transmitted between the two  
170 variables.

171 In the context of ocean modeling (or in general climate modeling) entropy is used to measure  
172 variability in a model output or available data. This is in tandem with interpretation of Shannon  
173 entropy in physical sciences as given in Carcassi et al. (2019). When calculating the Shannon  
174 entropy we are concerned about the possible states (e.g. the various bins in a histogram) the vari-  
175 able can (and does) go into while the variable value and its dimensions are of lesser importance.  
176 Entropy metrics measure variability in *bits* (when logarithm is of base 2) and hence changing the  
177 scale, e.g. switching from Celsius to Fahrenheit for temperature, does not change the value of  
178 variability (under equivalent binning). Mutual information and entropy are both dimensionally  
179 agnostic. They are also not sensitive to outliers (due to the weighting prefactor) and can capture  
180 nonlinear interactions (Watanabe 1960; Correa and Lindstrom 2013) and discontinuous distribu-  
181 tions including intermittently visited states. We will present the effect of correlation and outliers  
182 by examples of idealized random vectors.

183 The following methods and results sections are divided into the two parts of the overall paper  
184 objectives. Parts A of both sections relate to evaluating intrinsic and extrinsic variability in  
185 ensemble models. Parts B describe the usage of Shannon entropy and mutual information on  
186 coastal regional modeling data to understand and compare effects of using different boundary  
187 conditions.

## 188 2. Methods

### 189 *a. Part A: Intrinsic and Extrinsic variability for ensemble data*

190 We perform analysis on each grid point at the ocean surface or ocean bottom. Let a variable  
191 from the ensemble be given by  $f(n, t, x, y)$  where  $f$  is the variable,  $n$  denotes the index of the  
192 ensemble member and goes from 1 to  $N$ ,  $t$  is the time index and goes from  $t_1$  to  $t_M$ ,  $x, y$  represents  
193 the spatial grid point at the surface or bottom. At a particular grid point  $f(n, t, x, y)$  is  $f(n, t)$ . The  
194 total number of ensemble members is  $N$  and each member has  $M$  time steps. To get the signal due  
195 to extrinsic forcings, the "differencing" approach (Frankcombe et al. 2015) has been followed to

196 estimate forced response. This approach involves averaging over the ensemble members to derive  
197 the *ensemble mean*. The ensemble mean is given by:

$$g(t) = \frac{1}{N} \sum_{n=1}^{n=N} f(n,t) \quad (3)$$

198  $g(t)$  is a single time varying signal for each grid point obtained by averaging across the ensemble  
199 members. There are potential problems with assuming the ensemble mean represents extrinsic  
200 variability only, such as if models are differently sensitive to the forcing signal based on the model's  
201 equilibrium sensitivity as elaborated in Frankcombe et al. (2015). For a first order approximation,  
202 we will assume the ensemble mean is the best estimate of the forced response. Once  $g(t)$  is  
203 obtained, the intrinsic variability can be estimated by subtracting the ensemble mean  $g(t)$  from  
204 each ensemble member. Ensemble signal, forced response and intrinsic variability are then related  
205 by:

$$f(n,t) = g(t) + \eta(n,t), \quad (4)$$

206 where  $\eta(n,t)$  is the intrinsic variability or noise which differs from ensemble member to ensemble  
207 member. Note that the decomposition above takes place at each grid point. In Figure 1 a,  $f(n,t)$   
208 are shown by multi-colored ensemble members.  $g(t)$  is shown by thick black line. As seen Figure  
209 1 b,  $g(t)$  has a probability distribution shown in gray color and subsequently has first, second,  
210 and possibly important higher statistical moments. The gray colored density histogram shows  
211 variability due to extrinsic factors and the pink colored density histogram shows total variability  
212 given by extrinsic and intrinsic factors.

## 213 1) DETRENDING AND EVALUATING ENTROPIES

214 Analysis has been done with and without detrending the data to understand its impact. For  
215 detrending, a quadratic fit using least squares was found for the ensemble mean at each grid point  
216 and subtracted from all ensemble members and ensemble mean at the same grid point to get  
217 detrended data (e.g. Frankcombe et al. 2015). Detrending will remove some non-stationarity from  
218 the data but will also remove some part of the extrinsic variability. By this method, our aim is  
219 not to determine the forced response but to estimate the degree of *variability* contributed by the  
220 forced response (extrinsic response) and intrinsic variability originating from intrinsic chaos. The

221 ensemble mean  $g(t)$  was found at each grid point after detrending. For the non-detrended case, the  
 222 raw ensemble simulation data has been used to evaluate  $g(t)$  and  $\eta(n,t)$ .

223 Usually we are limited in the number of ensemble members due to computational costs so  
 224 we perform a *jugaad* in order to use *all* the ensemble members at once to evaluate information  
 225 entropies. All the ensemble members given by  $f(n,t)$  are rearranged into a single row vector  $f$  as:

$$f = [f(1,t_1), f(1,t_2), \dots, f(1,t_M), f(2,t_1), f(2,t_2), \dots, f(N-1,t_M), f(N,t_1), \dots, f(N,t_M)], \quad (5)$$

226 and  $g$  is row vector obtained by arranging  $N$  copies of  $g(t)$  in the following fashion:

$$g = \underbrace{[g(t_1), g(t_2), \dots, g(t_M)]}_1, \underbrace{[g(t_1), g(t_2), \dots, g(t_M)]}_2, \dots, \underbrace{[g(t_1), g(t_2), \dots, g(t_M)]}_N \quad (6)$$

227 This enables wide sampling and obtains an accurate probability distribution for  $f$  (assuming  
 228 approximate stationarity, or enforcing stationarity by detrending), and enables  $g$  to be of the same  
 229 size as  $f$  and having the same probability distribution as that of  $g(t)$ . The information statistics  
 230 we get at each grid point are time invariant since the complete time series is considered. It is  
 231 the user's choice to choose either the complete time series or a section of it for analysis. We  
 232 have chosen the whole time series, as this is a sufficient demonstration of the value of information  
 233 theory metrics. A time-evolving analysis raises additional issues about causality and shifting  
 234 probabilities distributions of climate states that are not the focus here (Liang 2013; DelSole and  
 235 Tippett 2018). By using the whole time-series, we are treating all variability as drawn from the  
 236 same distribution, and seek only to associate internal (associated with each ensemble member) and  
 237 external (associated with the ensemble mean) sources of variability following Leroux et al. (2018).  
 238 The time-series  $f$  and  $g$  are both expressed as row vectors of the same size,  $N \times M$ . This step  
 239 is crucial as vectors having same number of elements are necessary to evaluate joint probability  
 240 distribution. This enables us to calculate mutual information between  $f$  and  $g$ .

241 Calculating the Shannon entropy of  $f$  and mutual information between  $f$  and  $g$  is not a trivial  
 242 task. In fact optimal binning for precise measurement of information entropies is a research topic  
 243 in itself. Multiple techniques exist such as equidistant partitioning, equi-probable partitioning,  $k$   
 244 nearest neighbor, usage of B-spline curves for binning to name a few (e.g. see Hacine-Gharbi

et al. 2012; Kowalski et al. 2012; Knuth 2019). For a comprehensive review of the methods for estimating probability distribution see Papan and Kugiumtzis (2008). We have used equidistant partitioning throughout this article. For the case of GFDL-LE data, there were 1812 time steps available as monthly averages ranging from the year 1950 to 2100. As per Rice’s rule, 25 bins are needed for the GFDL-LE data. The bin width,  $\delta w$ , was calculated by dividing the range of data (maximum minus the minimum value) at the grid point with the least spread. The same bin width was used for all the grid points for Shannon entropy and mutual information. Equal bin width was used for the two variables in the joint probability and marginal probability calculation for mutual information. Maintaining the same bin width and range for all the grid points is crucial because information entropy strongly depends on the precision with which data is binned.

## 2) PROPOSED METRIC

Using  $f$  and  $g$ , we propose the following metric  $\gamma$ , which has the same intent as metrics in (Leroux et al. 2018) to quantify the fraction of variability that is intrinsic, i.e., the typical amount that is unique to an ensemble member or statistical instance, but unlike (Leroux et al. 2018) this metric is built from standard information theory quantities:

$$\gamma = 1 - \frac{I(f;g)}{H(f)}. \quad (7)$$

$H(f)$  is the Shannon entropy of  $f$ , and  $I(f;g)$  is mutual information between  $f$  and  $g$ .  $I(f;g)$  calculates the contribution of extrinsic signal  $g$  to the whole ensemble.  $H(f)$  is the total variability in the ensemble output which is the result of extrinsic and intrinsic factors. The metric  $\gamma$  gives the *ratio of intrinsic variability to total variability*.

$H(f)$  and  $I(f;g)$  are related through conditional entropy by  $H(f) = I(f;g) + H(f|g)$  (Cover 1999).  $H(f|g)$  is the conditional entropy<sup>1</sup>, i.e., average uncertainty about the value of  $f$  after  $g$  is known. It is the uncertainty in  $f$  that is not attributed to  $g$  but is attributed to noise  $\eta$ . Hence  $H(f) - I(f;g)$  estimates variability due to intrinsic chaos, and  $\gamma$  gives the fraction of the variability due to intrinsic chaos.

Returning to the binary tree analogy,  $I(f;g)$  would be the set of instructions sent by a source to reach one among  $2^{H(f)}$  possible destinations in the presence of noise having  $H(f|g)$  entropy. To

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<sup>1</sup>Conditional entropy  $H(X|Y)$  is defined by  $H(X|Y) = \sum p(x|y) \log_2 p(x|y)$  (Cover 1999). It is not necessary to calculate conditional entropy to arrive at  $\gamma$ , but understanding is aided by the expected relation between entropy and mutual information.

271 capture the entropy in the noisy binary tree, to each of the  $2^{I(f;g)}$  micro state possibilities noise  
272  $(2^{H(f|g)})$  gets multiplied and the relation becomes  $2^{H(f)} = 2^{I(f;g)}2^{H(f|g)}$ .

273  $I(f;g)$  takes into account any correlation or information shared between  $f$  and  $g$ . This is vital  
274 because even though the model spread  $\eta$  is being treated similarly to noise added to the mean  
275 signal, it might be that model spread depends on the mean signal. A simple example is if the  
276 model spread is relative (e.g., 10% of the mean signal), rather than absolute (e.g., 2 units), then  
277 there is information about the model spread contained in the ensemble mean signal. This situation  
278 is sometimes called multiplicative noise in contrast to additive noise. The nonlinear and chaotic  
279 nature of fluid mechanics often leads the mean flow to amplify the chaotic signal (e.g., eddies) and  
280 thereby result in altered variability statistics. When  $f \rightarrow g$ , then  $I(f;g) \rightarrow H(f) = H(g)$  from  
281 (2). This makes  $\gamma = 0$  when there is no intrinsic variability or chaos. When intrinsic chaos fully  
282 dominates the ensemble output, i.e.  $f$  and  $g$  are fully decorrelated, then  $I(f;g) = 0$  yielding  $\gamma = 1$ .  
283 We see that  $\gamma$  satisfies the extremes of zero noise as well as total chaos.

284 Another analogue for a climate system component is a noisy communication channel as given  
285 in Leung and North (1990), where the governing equations of ocean (atmosphere) modeling are  
286 taken to communicate from forcing to response. The extrinsic forcings are inputs to the channel,  
287 the intrinsic chaos is the noise created because of channel's inherent mechanisms while the outputs  
288 are the ensemble members. A noiseless channel will give  $\gamma$  as zero and completely noisy channel  
289 where output is independent of input will yield  $\gamma$  as 1.

290 A seemingly enticing and simpler alternative is  $\gamma = 1 - \frac{H(g)}{H(f)}$ , i.e. just the difference between  
291 ensemble entropy and mean entropy as a ratio with the ensemble entropy. However, this formulation  
292 is incorrect because  $H(g)$  does not quantify the contribution of extrinsic factors to the variability  
293 in the ensemble, it only quantifies the variability of the mean. Relatedly,  $H(f) - H(g)$  does not  
294 correctly manage mutual information between the ensemble members and their mean in estimating  
295 the intrinsic variability.

296 Recently, another alternative was proposed by Gomez (2020): using Shannon entropy directly as  
297 a measure of intrinsic variability. They propose using Shannon entropy of model spread  $\eta(n, t)$  at  
298 each time step normalized by the logarithm of the number of bins utilized. Their metric has a lower  
299 limit of 0 and an upper limit of 1, where 0 denotes zero noise and hence zero intrinsic variability  
300 and 1 denotes complete intrinsic variability. Again, this metric is similar to  $\gamma$  in building upon

301 information theory, but  $\gamma$  takes into account the variability of the ensemble mean, correlations  
 302 between the ensemble mean and the intrinsic variability, and it is time invariant. A time dependent  
 303 version of  $\gamma$  can be made using running time windows instead of the whole time series, but care  
 304 in quantifying or controlling for lack of stationarity is needed in this interpretation (DelSole and  
 305 Tippet 2018). The Gomez (2020) metric uses the spread of the ensemble members similar to  
 306 measuring Shannon entropy whereas  $\gamma$  utilizes, in an abstract sense, the set of instructions required  
 307 to choose a destination for the particular variable among the possible model states.

### 308 3) VARIANCE BASED METRIC

309 A variance based metric as given in (Leroux et al. 2018) has been utilized to compare to our  
 310 information based metric. The variance based metric measures intrinsic and extrinsic variability  
 311 using the second moment, variance. It involves calculation of the following terms  $\sigma_g$  and  $\sigma_\eta$  given  
 312 by:

$$\sigma_g^2 = \frac{1}{M} \sum_{t=1}^{t=M} \left( g(t) - \overline{g(t)} \right)^2, \quad (8)$$

$$\sigma_\eta^2(t) = \frac{1}{N} \sum_{n=1}^N \eta(n,t)^2, \quad (9)$$

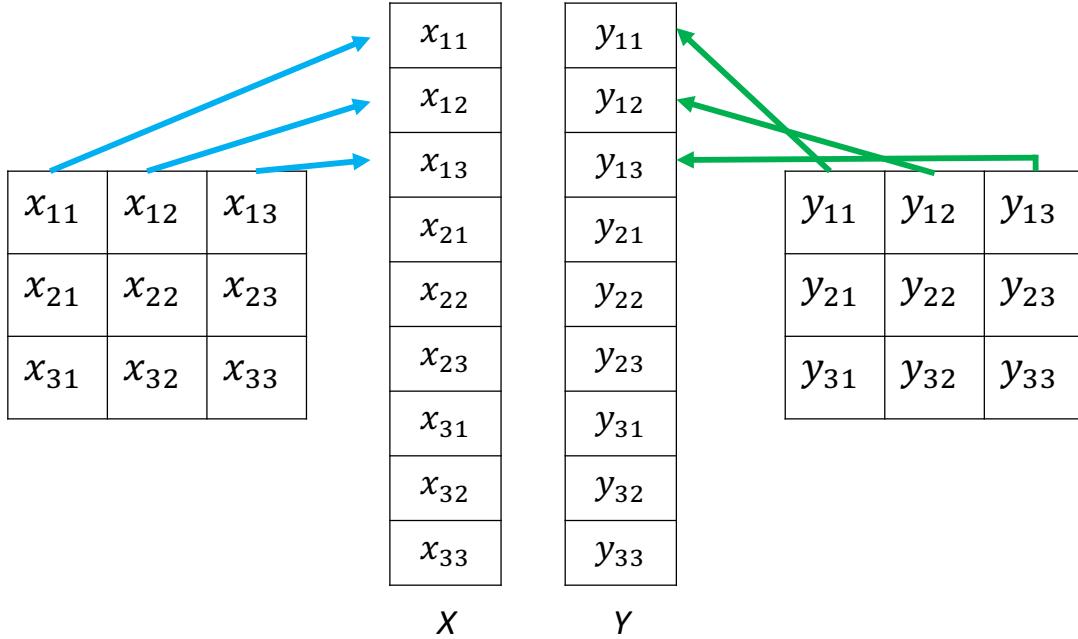
314 where the overbar denotes temporal averaging. The total variability has been estimated  
 315 as  $\left( \sigma_g^2 + \overline{\sigma_\eta^2(t)} \right)^{1/2}$ . The forced variability  $\sigma_g$  is equivalent to  $I(f;g)$ , and total variability  
 316  $\left( \sigma_g^2 + \overline{\sigma_\eta^2(t)} \right)^{1/2}$  is equivalent to  $H(f)$ . Hence,  $\gamma$  is compared with  $\gamma_{std}$  given by

$$\gamma_{std} = \frac{\left( \overline{\sigma_\eta^2(t)} \right)^{1/2}}{\left( \sigma_g^2 + \overline{\sigma_\eta^2(t)} \right)^{1/2}} \quad (10)$$

### 317 *b. Part B*

#### 318 1) IMPACT OF CHANGES IN BOUNDARY FORCINGS IN COASTAL MODELS

323 Here instead of using the new metric  $\gamma$ , we use its components: Shannon entropy and mutual  
 324 information individually to compare variability between different simulations. Quantifying differ-  
 325 ences because of modifications in the extrinsic forcings may be required for coastal applications



319 FIG. 2. Flattening process for comparing two dimensional fields using Shannon entropy and mutual information.  
 320 As the flattened arrays  $x_1, x_2, \dots$  and  $y_1, y_2, \dots$  might not have linear dependence on each other, using linear  
 321 dependence measure such as Pearson correlation will yield incorrect results. Mutual information measures  
 322 nonlinear correlations and hence captures all linear and non-linear dependence.

326 where systems vary predominantly due to external forcings. For these forcing significance ex-  
 327 periments, OSOM was run after modifying the external forcings (Table 1). OSOM is forced by  
 328 tides, river runoff, atmospheric winds and air-sea fluxes, etc. (Full details of the model can be  
 329 found in Sane et al. 2021). For this comparison, we quantify the effects of altering forcing on 4  
 330 modeled variables: sea surface temperature and salinity, and bottom temperature and salinity. Four  
 331 altered forcing sets were utilized, beyond set (1) Full set of atmospheric forcings using the North  
 332 American Mesoscale (NAM) analyses, a data-assimilating, high resolution (12 km) meteorologi-  
 333 cal simulation ([https://www.ncei.noaa.gov/data/north-american-mesoscale-model/](https://www.ncei.noaa.gov/data/north-american-mesoscale-model/access/historical/analysis)  
 334 [access/historical/analysis](https://www.ncei.noaa.gov/data/north-american-mesoscale-model/access/historical/analysis)) denoted as FF. FF stands for full forcing. (2) Full set of at-  
 335 mospheric forcings but using the Northeast Coastal Ocean Forecast System (NECOFS) winds  
 336 (Beardsley and Chen 2014) instead of NAM, denoted as NECOFS. (3) River flows are replaced  
 337 with their monthly-averaged flow, other forcing as in FF (4) River flows set to zero, other forcing

338 as in FF. (5) Wind forcing set to zero, other forcing as in FF. These forcings have been tabulated in  
 339 Table 1. The aim is to quantify the effect on total variability by removing or altering one of many  
 340 processes which might contribute.

Forcing Set	Wind forcing	River forcing
FF	NAM	As Observed
NECOFS	NECOFS	As Observed
MR	NAM	Time-averaged
ZR	NAM	Zero river input
ZW	Zero winds	As Observed

341 TABLE 1. Different types of forcing combinations employed to test their effect on variability. FF stands for full  
 342 forcing: winds, tides, rivers, etc. For more details see Sane et al. (2021). MR: mean rives; ZR: zero rivers; ZW:  
 343 zero wind.

344 To evaluate Shannon entropy, the spatial output at a particular instant of time was rearranged into  
 345 a row vector by a process called 'flattening' as shown in Figure 2. Land mask points were removed.  
 346 A variable  $x$  which is a two-dimensional variable was converted to one-dimension (flattened) by  
 347 concatenation. Shannon entropy was found out for the flattened variable at each time step to obtain  
 348 time varying entropy of the surface or bottom variable.

349 Mutual information was applied between the flattened row vectors. Our focus is towards a  
 350 pragmatic approach on using information theory for simulation comparisons, as opposed to an  
 351 equation for the evolution of Shannon entropy and mutual information with respect to time (see  
 352 Liang and Kleeman 2005). Relative comparison between mutual information values is what we  
 353 seek. For example, if mutual information of surface salinity between FF and MR is higher than  
 354 between FF and ZR, this implies the penalty for using time-averaged river runoff is not as severe  
 355 as using zero river runoff. Replacing FF with MR will give better results than ZR. Small errors in  
 356 river runoff flow rates won't cause appreciable changes to surface salinity than using zero rivers.

### 357 3. Results

#### 358 a. Part A

##### 359 1) IDEALIZED GAUSSIAN ARRAYS

360 We test our metric,  $\gamma$ , equation (7) on synthetic data consisting of idealized arrays of Gaussian  
361 data:  $\mathcal{N}(0, 1)$ . For a normal Gaussian distribution Shannon entropy depends<sup>2</sup> only on the standard  
362 deviation  $\sigma$  i.e.  $H = \log_2(2\pi e\sigma^2)$ . The variability in a Gaussian distribution can be increased or  
363 decreased by changing its standard deviation. Our goal is to compare  $\gamma$  and  $\gamma_{std}$ . We set out our  
364 numerical experiment as follows: we create 10 arrays, each having 10,000 elements drawn from a  
365 Gaussian distribution. Any two arrays from those 10 have a prescribed linear Pearson correlation  
366 coefficient from 0 to 1.

367 Thus, the 10 arrays covary linearly with a specified correlation coefficient. These 10 arrays  
368 represent ensemble members from climate simulations. The mean of 10 members gives us the  
369 synthetic forced variability signal as would be determined from the model output; averaging over  
370 the 10 ensemble members reduces the contribution from uncorrelated variability and reaffirms the  
371 covarying component into the forced variability. We apply  $\gamma$  and  $\gamma_{std}$  on this synthetic ensemble  
372 by varying the prescribed correlation coefficient from 0 to 1. Figure 3 shows that as expected  
373 both metrics increase as the correlation decreases, i.e., as internal variability dominates forced.  
374 Both metrics behave similarly when correlation decreases, i.e. noise increases but  $\gamma$  is more  
375 sensitive as correlation tends to 1. This distinction is due to the logarithmic nature of Shannon  
376 entropy for Gaussian distributions—in essence, information measured in bits is not proportional to  
377 distance measured between distributions in terms of summed variance—in the examples following  
378 the consequences of this distinction will become clearer. Critically both functions are monotonic  
379 with correlation, however so relative comparisons (more intrinsic fraction in this region vs. that  
380 region) are preserved.

381 A second related experiment was derived from the first is also shown in Figure 3: adding outliers  
382 outside of the Gaussian distribution. 50 out of 10000 elements of each individual member were  
383 artificially corrupted (values were set to a constant value of 5) to test the sensitivity of both the  
384 metrics. Figure 3 shows that  $\gamma$  is insensitive to outliers while  $\gamma_{std}$  is not.  $\gamma$  is not sensitive because

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<sup>2</sup> $H = \log_2 2\pi e\sigma^2$  is the Shannon entropy of a Gaussian distribution when probability density is continuous with  $\sigma$  as standard deviation. The Shannon entropy of a discrete probability distribution differs, which is inconsequential here but the reader is encouraged to read Jaynes (1962). Consistently here discretely sampled and binned probability distributions are obtained directly from data without any further parameterization.

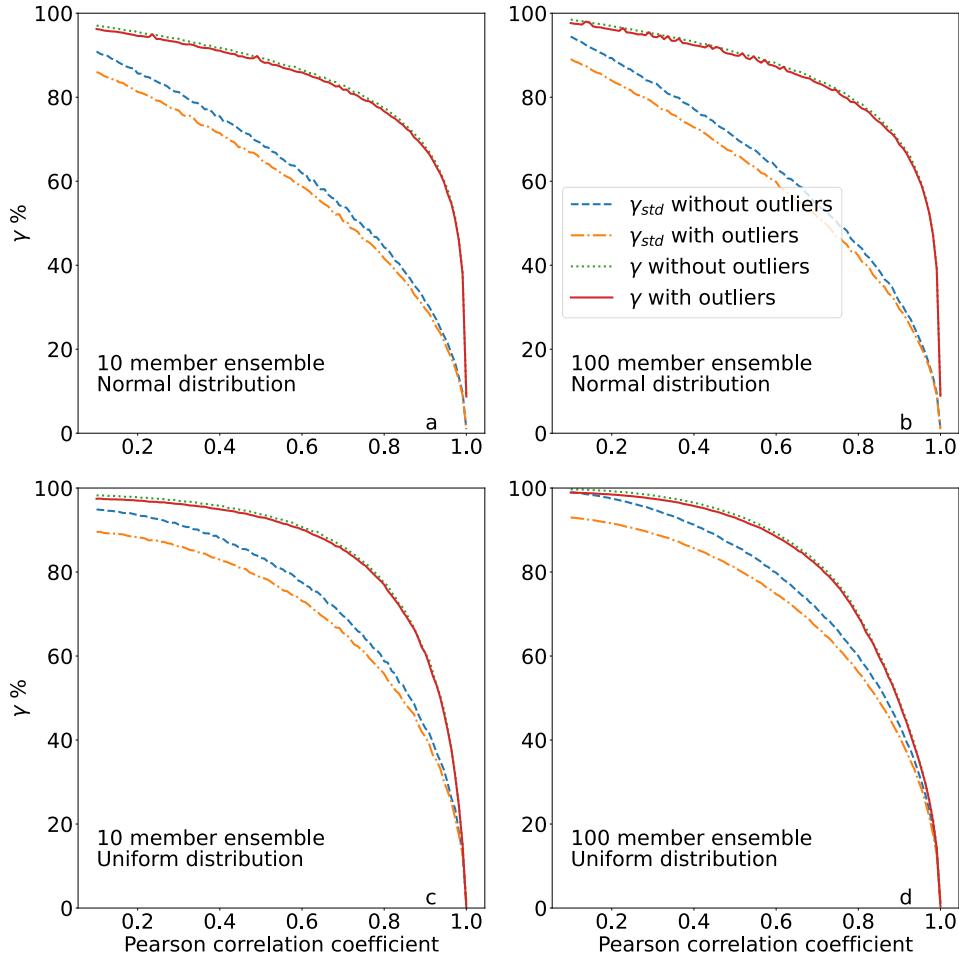
385 outliers occur less frequently and hence do not affect the probability distribution much, especially  
386 with the prefactor in (1) and (2). Hence information theory metrics are robust in comparison to  
387 using standard deviation (or variance). If the outliers (extreme events) occur at higher frequencies,  
388 information metrics will naturally start sensing them even if they are discontinuous from the typical  
389 conditions (e.g., multimodal distributions). The above process was repeated for 100 ensemble  
390 members each sampled from Gaussian distributions. Increasing the number of ensemble members  
391 does not change the result qualitatively for both the experiments. The results for 10 member  
392 Gaussian ensemble is shown in Figure 3 a and 100 member in Figure 3 b.

393 Additionally, a set of experiment was done by using uniformly distributed data  $U(-1, 1)$ . The  
394 prescribed correlated vectors were created using the procedure outlined in Demirtas (2014). 10  
395 and 100 ensemble members were created and  $\gamma$  and  $\gamma_{std}$  was found between the members and their  
396 mean. Results are shown in Figure 3 c, d respectively. The outlier had a value of 1.5. In all the  
397 cases,  $\gamma$  was less sensitive to outliers than  $\gamma_{std}$ .

## 406 2) REGIONAL COASTAL MODEL OUTPUT

407 In this section we show the results of applying  $\gamma$  and  $\gamma_{std}$  on realistic simulation data from the  
408 Ocean State Ocean Model, hereafter OSOM (Sane et al. 2021). OSOM uses the Regional Ocean  
409 Modeling System (ROMS) (Shchepetkin and McWilliams 2005) to model Narragansett Bay and  
410 surrounding coastal oceanic regions and waterways. OSOM's primary purpose is for understanding  
411 and predictive modeling and forecasting of the estuarine state and climate of this Rhode Island  
412 body. Sane et al. (2021) gives more details about the model.

413 Using OSOM, an ensemble of simulations have been performed using perturbed initial (ocean)  
414 conditions under the same atmospheric and tidal forcing for the months July - August of 2006. This  
415 ensemble consists of 10 members. The data during the first predictability window ( 20 days) that is  
416 sensitive to initial conditions has been ignored and the remaining simulation has been used to look  
417 at variability within the "climate projection" of the model beyond when forecasts sensitive to initial  
418 conditions are possible (see the related application of information theory to assess predictability  
419 in Sane et al. 2021). We examine whether the modeled temperature and salinity at each grid point  
420 follow normal distributions by evaluating the skewness and kurtosis of the ensemble mean at each  
421 grid point. Figure 4 shows skewness and kurtosis for sea surface salinity and temperature as well



398 FIG. 3. Information theory metric of intrinsic vs. extrinsic variability  $\gamma$  as a function of correlation coefficient  
 399 in idealized Gaussian correlated arrays (a and b) and idealized uniformly distributed arrays (c and d). The  
 400 horizontal axis is the correlation coefficient between mean member and ensemble members. The vertical axis  
 401 shows the information theory metric  $\gamma$  from (7) and the traditional metric  $\gamma_{std}$  from Equation (10). A second  
 402 related experiment adding (50 out of 10,000) “corrupted” outliers to each individual member is also shown. The  
 403 information theory metric  $\gamma$  does not change for these outliers which shows its robustness while  $\gamma_{std}$  is highly  
 404 sensitive. Results are similar for Gaussian distribution members and uniformly distributed members.  $\gamma$  is more  
 405 sensitive towards linear correlation of 1. This is due to the logarithmic nature of  $\gamma$ .

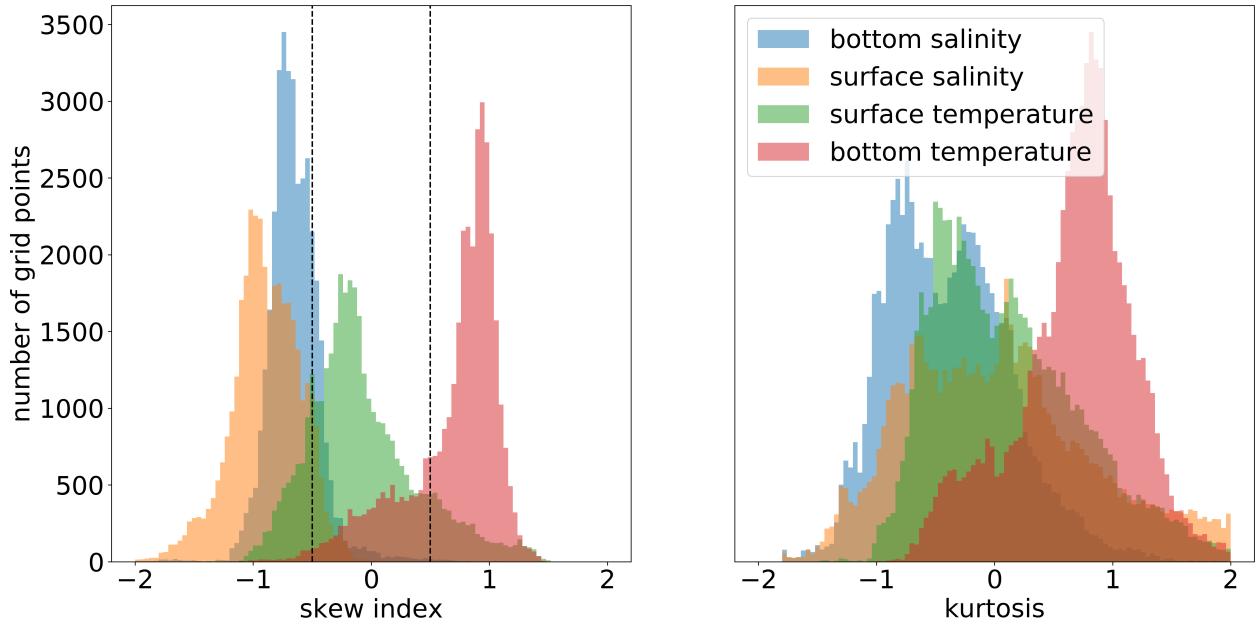
422 as bottom salinity and temperature for the Narragansett Bay region. The horizontal axis shows  
 423 skewness and excess kurtosis, which are the third and fourth statistical moments respectively,  
 424 normalized by powers of the standard deviation to dimensionless ratio and in the case of excess

425 kurtosis a constant value of 3 is subtracted. For Gaussian distributions, skewness and excess  
426 kurtosis both should be close to zero. The vertical axis denotes the number of occurrences at a  
427 grid point. Observe that the majority of grid point values are away from zero. These variables  
428 are considerably non-Gaussian in OSOM. Thus, Equation (10) is at a disadvantage, because the  
429 prevalence of higher statistical moments implies that the variance does not contain a complete  
430 description of the variability. The information theory metric (7) is suitable for such data as it takes  
431 into account higher moments and does not rely on Gaussian distributions.

432 Figure 5 shows the ratio of intrinsic variability to total variability applied on every grid point  
433 for OSOM.  $\gamma$  is displayed on left whereas  $\gamma_{std}$  is shown on right for comparison. The features  
434 highlighted by both metrics are qualitatively different. The contribution of intrinsic chaos to total  
435 variability is more uniform using the  $\gamma$  metric than using  $\gamma_{std}$ . The intrinsic chaos displayed using  
436  $\gamma_{std}$  might be misleading because the probability distributions are non-Gaussian. Furthermore,  
437 where the  $\gamma$  metric highlights internal variability tends to agree in similar dynamical locations—all  
438 river mouths show high surface salinity intrinsic variability. While surface temperature intrinsic  
439 variability is higher in more open regions of the Bay where eddies form intermittently due to  
440 varying topography. Also note that the ranges are quite different between  $\gamma$  and  $\gamma_{std}$ , but this is to  
441 be expected from the different rate of increase with correlation seen in Figure 3.

### 464 3) COMMUNITY EARTH SYSTEM MODEL LARGE ENSEMBLE

465 A complementary experiment was performed by using  $\gamma$  to evaluate internal vs. forced variability  
466 in the global climate simulation output for climate change scenario RCP8.5 using the (randomly  
467 selected among the models compared) GFDL-LE model. All the 40 members from the ensemble  
468 were utilized. Variability of sea surface temperature (Figures 6) as well as sea surface salinity  
469 (Figures 7) were estimated using both  $\gamma$  and  $\gamma_{std}$  (upper left and upper right). Similar results  
470 were obtained for the detrended data for temperature (Figures 8) and salinity (Figures 9) The  
471 skewness and excess kurtosis of the ensemble mean were also plotted to find the deviation of  
472 variables away from Gaussian distributions (lower). Regions shaded in purple have low values  
473 of excess kurtosis and skewness and might be considered Gaussian. The detrended data shows a  
474 higher percentage of intrinsic variability than non-detrended data which suggests that detrending

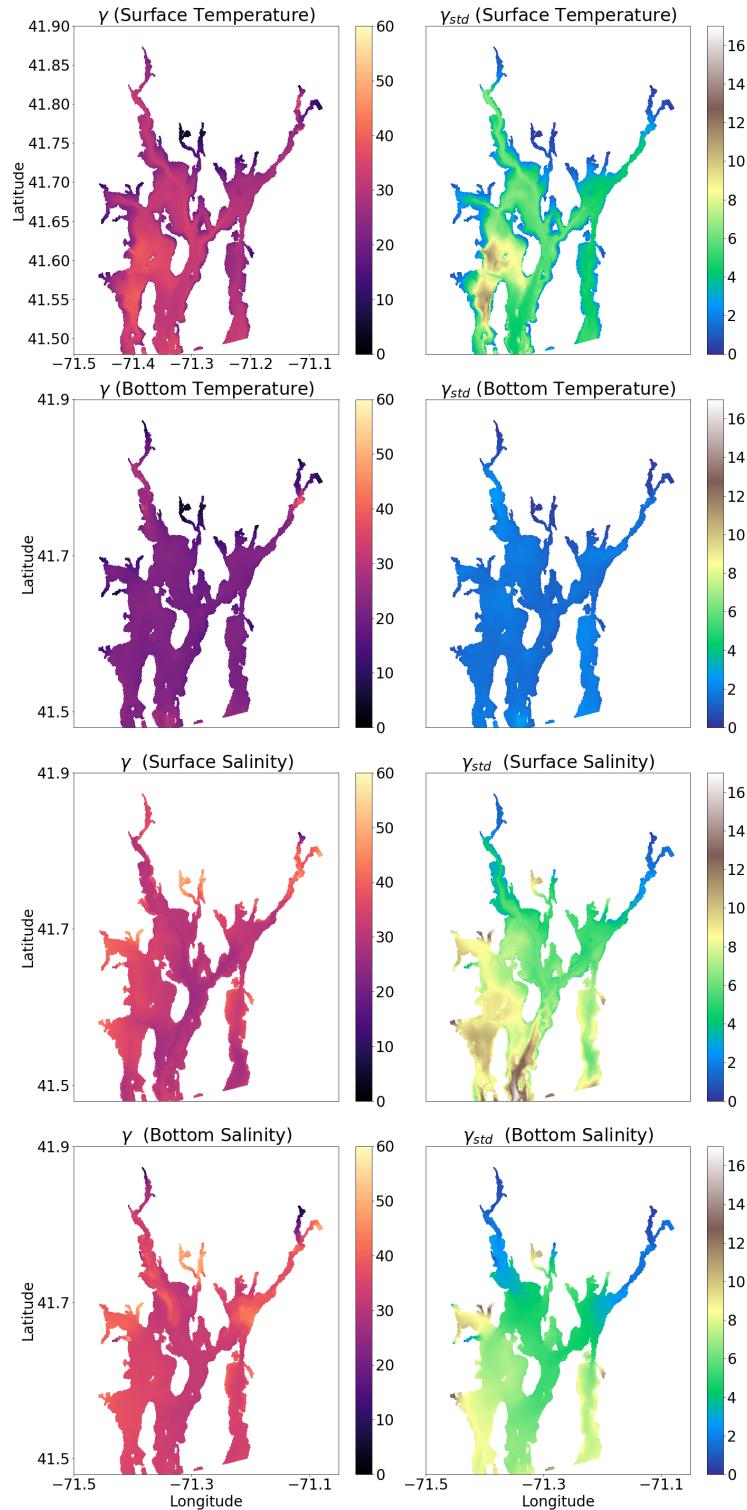


442 FIG. 4. Grid point wise kurtosis for OSOM output. Kurtosis is not closer to zero within (-0.5, 0.5) suggesting  
 443 the data distribution is non Gaussian.

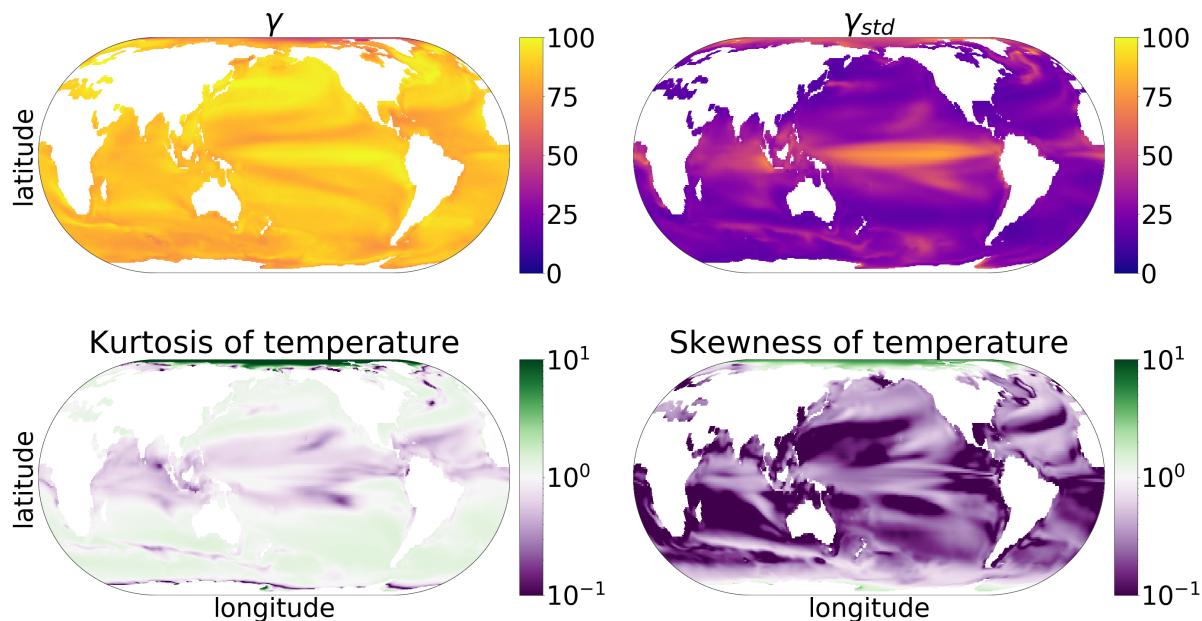
475 has removed some proportion of extrinsic variability, presumably the climate change signal present  
 476 in these simulations.

477 Note in particular the Arctic sea surface temperatures, which have a highly skewed and excessive  
 478 kurtosis distribution due to the freezing point of seawater. The standard metric ( $\gamma_{std}$ ) deems this  
 479 region to be among the most intrinsically variable in the world, while the information theory metric  
 480 has it as a low intrinsic variability region. It is clear that a Gaussian metric should not be applied  
 481 to this region due to the skewness and excess kurtosis, and in this case the inference is opposite  
 482 using the two metrics. In the equatorial Pacific where Gaussian statistics are more reliable, the two  
 483 metrics agree that internal variability is high.

484 A less drastic failure occurs from the modest excess kurtosis in extra-tropical temperatures and  
 485 in a few isolated regions in surface salinity. These regions are also non-Gaussian, but also are  
 486 not heavily skewed (i.e., they are more long-tailed and intermittent than Gaussian). These regions  
 487 differ in relative estimation of intrinsic versus total variability. It is also the case that the  $\gamma$  metric  
 488 is closer to one in most regions than  $\gamma_{std}$ , which is to be expected when the correlation coefficients  
 489 are low from Figure 3.



444 FIG. 5. Metrics  $\gamma$  vs  $\gamma_{std}$  for OSOM output. Both metrics show different contribution of intrinsic variability  
 445 to total variability.  $\gamma$  is more uniform throughout the domain than  $\gamma_{std}$ . Colormaps for  $\gamma$  and  $\gamma_{std}$  are different  
 446 to highlight the different ranges each of them have.  $\gamma_{std}$  for bottom temperature has maximum value of 5%, and  
 447 pattern is almost uniform except at the river sources where values are on the lower side (less than 1%).

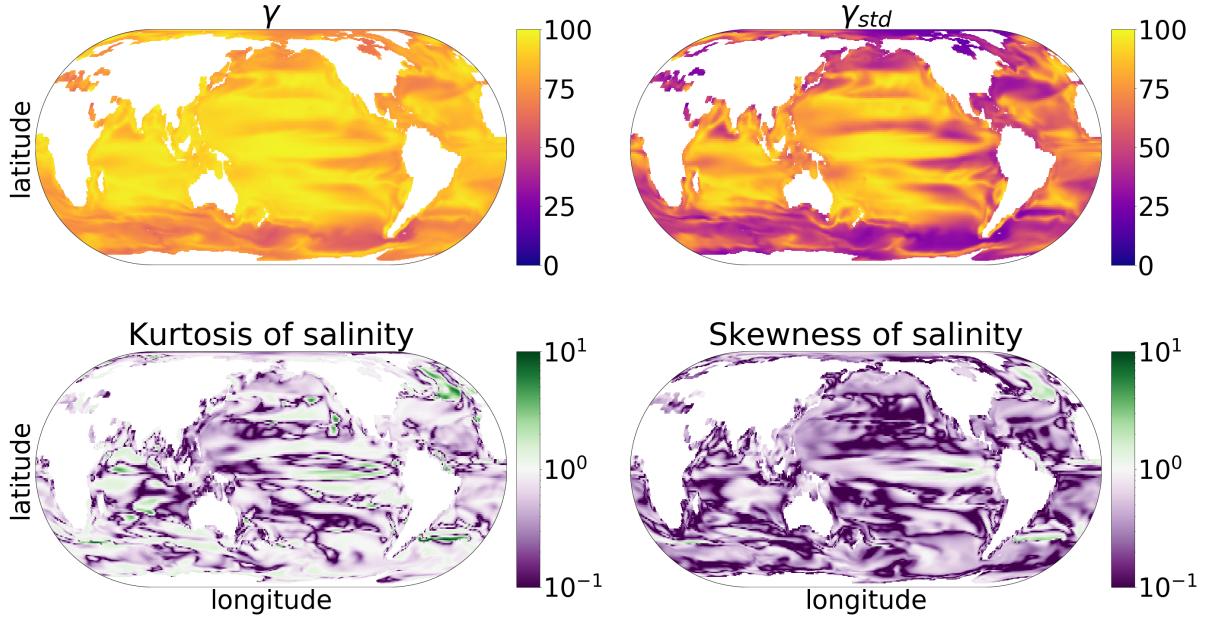


448 FIG. 6. Top: Intrinsic to total variability percentage for sea surface temperature. Bottom: Excess kurtosis  
 449 and skewness of the ensemble mean of temperature at each grid point. Values closer to zero (within 0.5 of  
 450 zero, purple shades) are considered approximately Gaussian. The deviation of ensemble mean away from non  
 451 normality implies that the ensemble members are also non normal. The Arctic regions have the most skewness  
 452 and excess kurtosis implying non-Gaussian distributions.

490 *b. Part B*

491 1) IMPACT DUE TO CHANGES IN BOUNDARY CONDITIONS IN COASTAL MODELS:

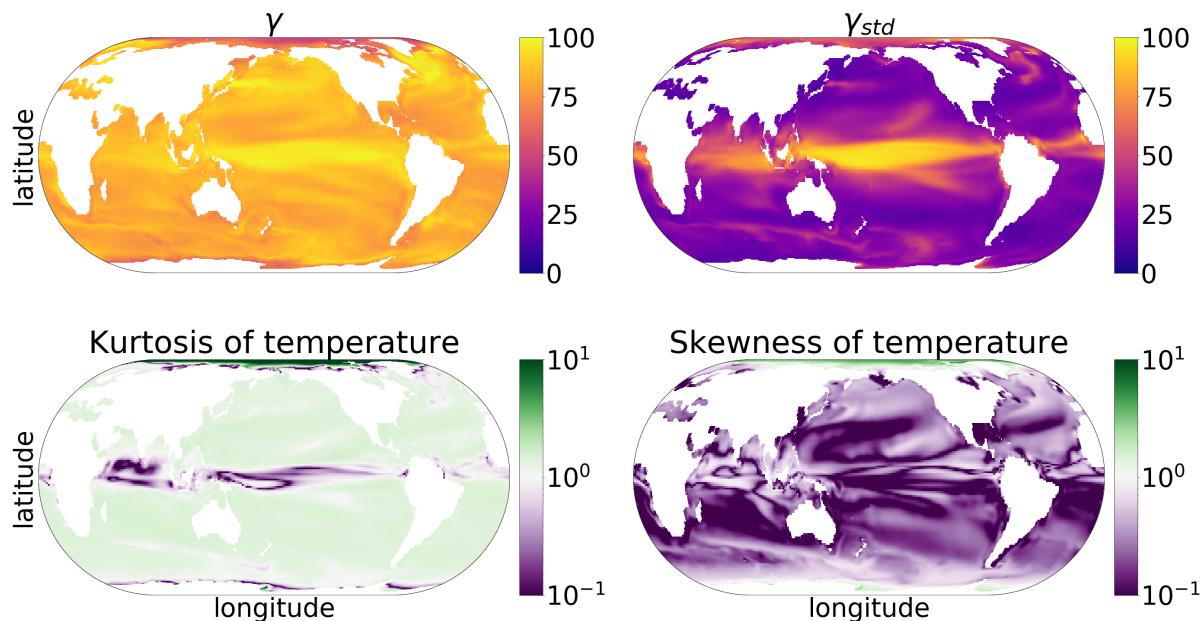
499 We show results of the coastal model analysis under different forcing in Figures 10 and 11.  
 500 Entropy has been plotted with respect to time to aid understanding. In Figure 10, Shannon entropy  
 501 is plotted for spatial quantities. For example, for surface salinity, all the surface values have been  
 502 considered to find Shannon entropy using the flattening approach. Figure 11 displays mutual  
 503 information. It is user's choice to choose the type of domain, here we have chosen the same domain  
 504 of OSOM as shown in Figure 5. If Shannon entropy is more or less equal for two forcings, it implies  
 505 they similarly affect variability. Mutual information should be compared for two pairs of forcings.  
 506 Greater mutual information implies the two pairs share more *bits* of information, suggesting one  
 507 of the forcing in that pair can be replaced with the other without significantly affecting variability.



453 FIG. 7. Top: Intrinsic to total variability percentage for sea surface salinity. Bottom: Kurtosis and skewness  
 454 of the ensemble mean of salinity at each grid point. Values closer to zero (within 0.5 of zero, purple shades) are  
 455 considered approximately Gaussian.

#### 508 4. Discussion

509 Our numerical experiments performed using  $\gamma$  on idealized Gaussian arrays show that  $\gamma$  is  
 510 monotonic and decreases as the linear Pearson correlation coefficient increases. Thus aside from  
 511 the qualitative differences the new metric finds when the data are non-Gaussian, the ranges of  
 512 intrinsic versus total variability are quite different between  $\gamma$  and  $\gamma_{std}$ . This is to be expected from  
 513 the different rates of increase with correlation seen in Figure 3. The traditional metric ( $\gamma_{std}$ ) falls  
 514 approximately linearly as the correlation coefficient increases, so that a correlation coefficient of  
 515 0.5 gives a  $\gamma_{std}$  just above 0.5. The new metric  $\gamma$  agrees with  $\gamma_{std}$  that correlation of 0 implies  
 516  $\gamma = 1$ , and correlation of 1 implies  $\gamma = 0$ , but for a correlation of 0.5 is closer to  $\gamma = 0.9$ . Only very  
 517 near correlation coefficients of 1 does  $\gamma$  fall below 0.5. If roughly linear dependence on correlation  
 518 coefficient is desired,  $\gamma$  can be raised to a power— $\gamma^3$  resembles  $\gamma_{std}$  and  $\gamma^6$  resembles the correlation  
 519 coefficient. These higher powers do not lose the ability to apply to non-Gaussian data nor become  
 520 non-monotonic, but they will lose their interpretation as a ratio of bits of information entropy, and

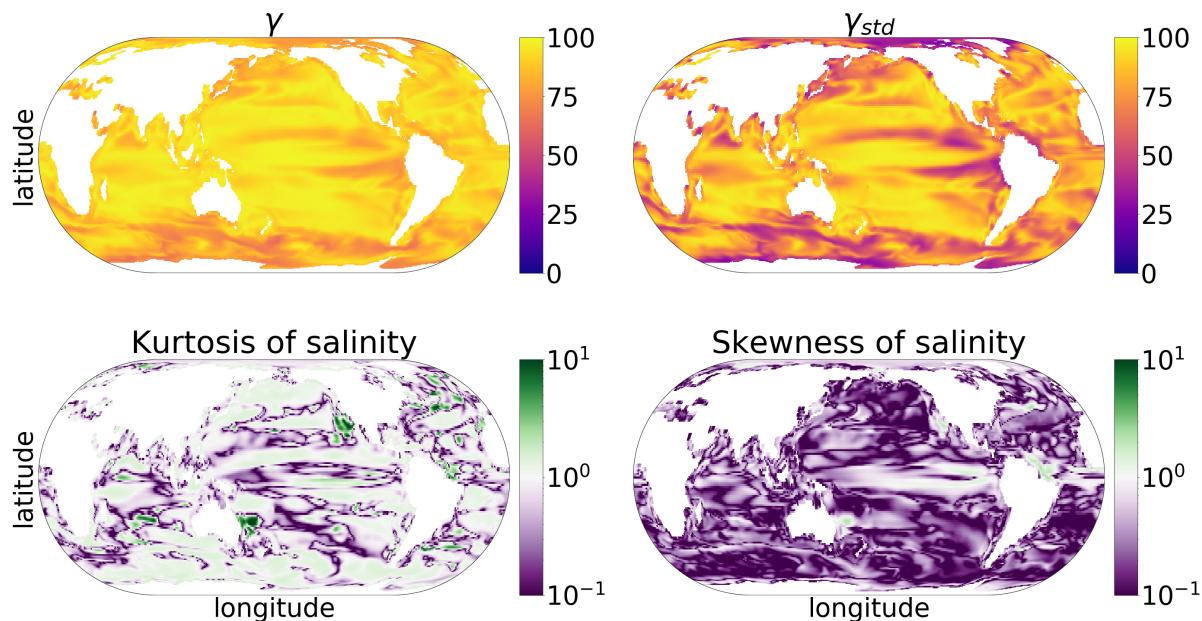


456 FIG. 8. Top: Intrinsic to total variability percentage for detrended sea surface temperature. Bottom: Excess  
 457 kurtosis and skewness of the ensemble mean of temperature at each grid point. Values closer to zero (within 0.5  
 458 of zero, purple shades) are considered approximately Gaussian. The deviation of ensemble mean away from non  
 459 normality implies that the ensemble members are also non normal. The Arctic regions have the most skewness  
 460 and excess kurtosis implying non-Gaussian distributions.

521 instead reflect ratios of bits cubed of information entropy, etc. An alternative is to take  $\gamma_{std}$  raised  
 522 to a different power:  $\gamma_{std}^{1/3}$  is roughly similar to  $\gamma$ .

523 To check for sensitivity due to our binning choice, the endpoints of each bin were shifted by  
 524  $\delta w/2$  and the results were compared. In theory, such a shift should not meaningfully affect the  
 525 outcome, so this comparison gives a sense of how sensitive the results are to binning choices. For  
 526 temperature, the raw GFDL-LE data (without detrending) gave an error of 4.7% in  $\gamma$  (see next  
 527 section for definition) and detrended data gave an error of 11%. For salinity, the error in  $\gamma$  for  
 528 raw data was 1.7% and for detrended data was 2%. Similar analysis for ROMS-OSOM coastal  
 529 ensemble data gave negligible error for shifting the bin endpoints. Different binning strategies will  
 530 be left to be explored for future research.

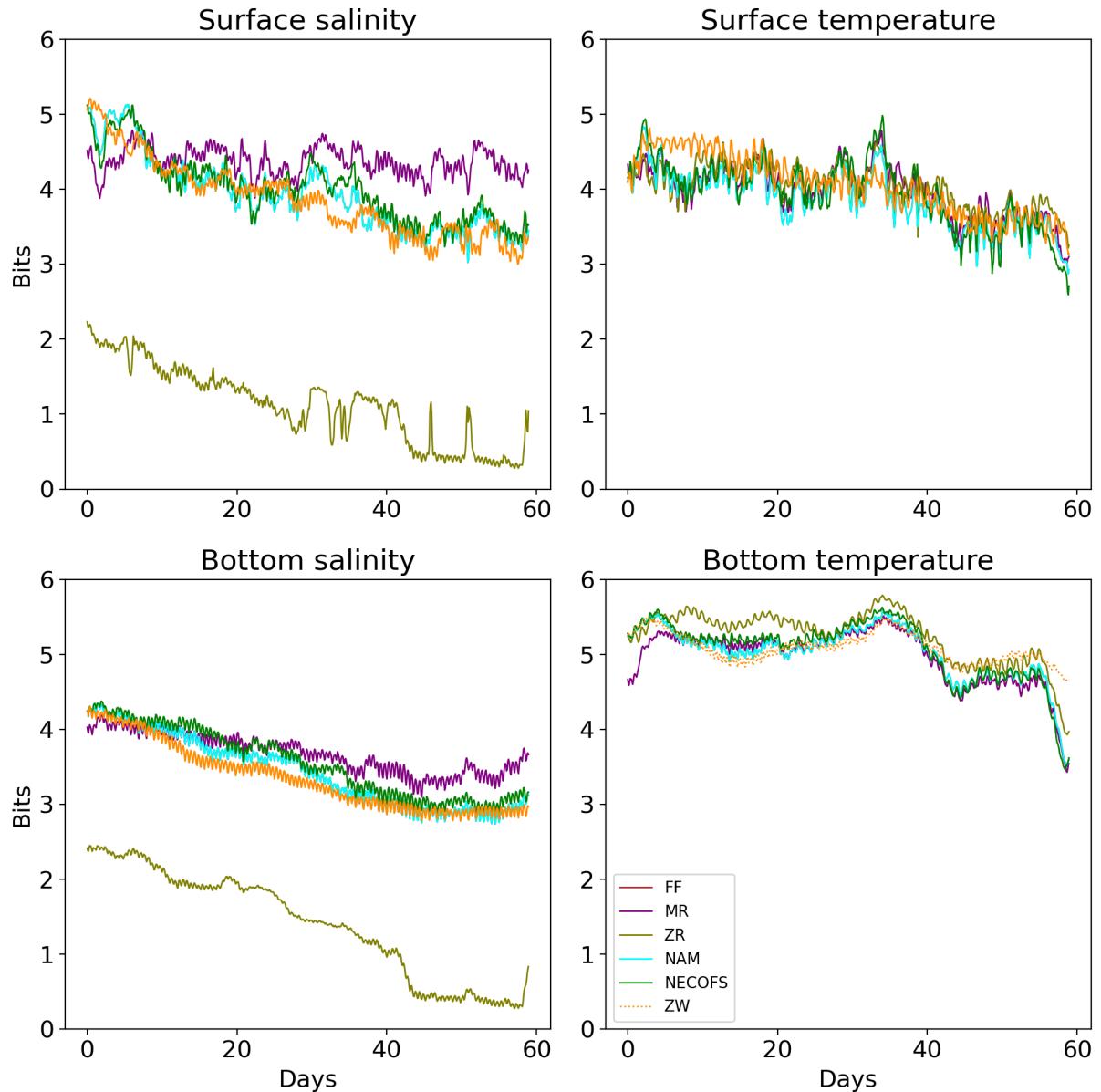
531 As can be seen in Figures 5, 6, and 7, information theory metrics show different patterns when  
 532 compared to variance. Information theory metrics, especially mutual information, account for



461 FIG. 9. Top: Intrinsic to total variability percentage for detrended sea surface salinity. Bottom: Kurtosis and  
 462 skewness of the ensemble mean of salinity at each grid point. Values closer to zero (within 0.5 of zero, purple  
 463 shades) are considered approximately Gaussian.

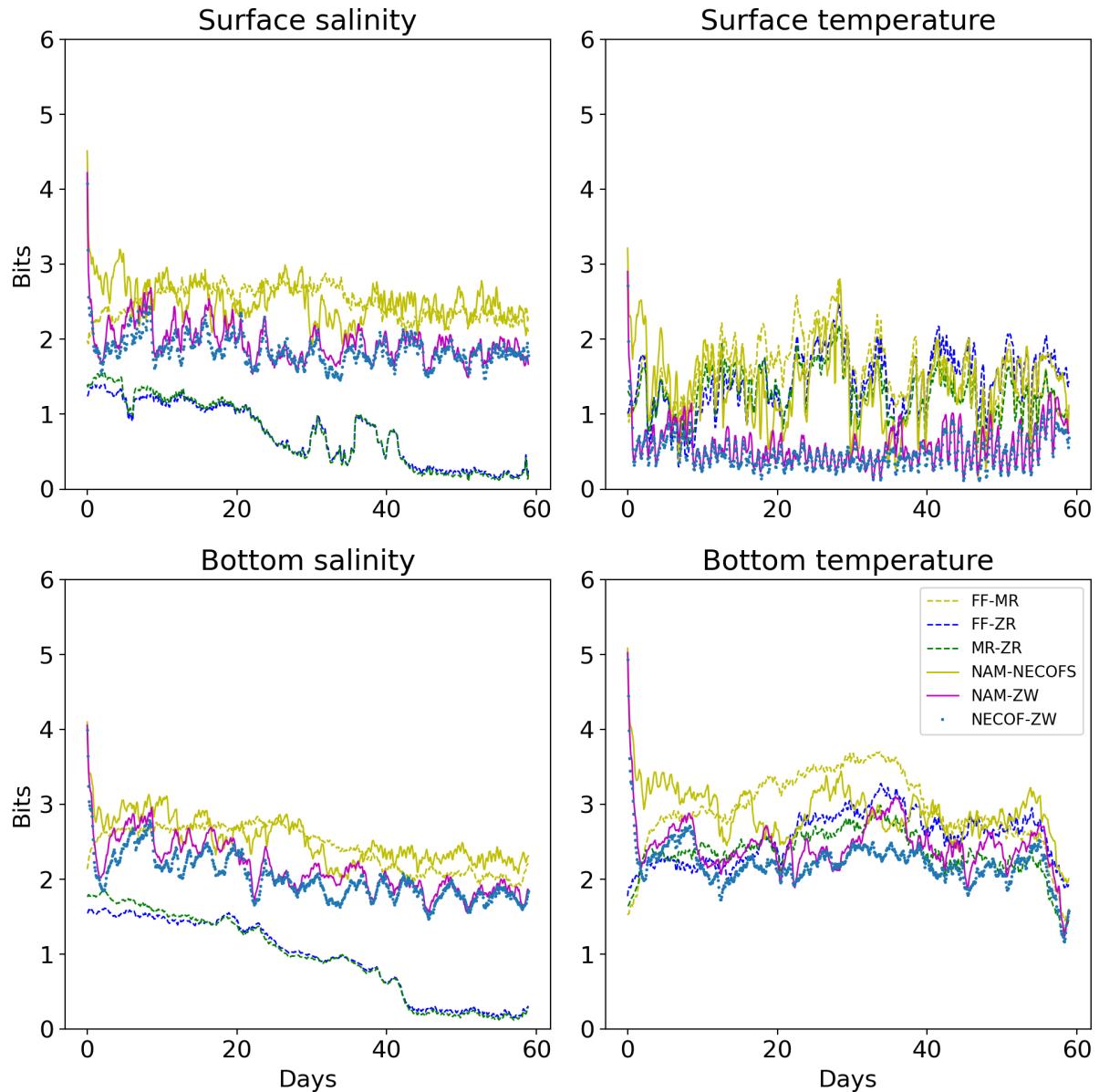
533 *all* non-linear shared information between the ensemble members and the mean including linear  
 534 correlation, and this is one reason for the differences. We have argued that non-Gaussian statistics  
 535 are another (which is not wholly independent of non-linear shared relationships). There are likely  
 536 other aspects of differences between these metrics, but the management of these two expected  
 537 aspects of geophysical fluids—nonlinear relationships and non-Gaussian distributions—justify the  
 538 introduction of the new metric.

539 For the regional coastal model OSOM, forcings differ as to how they affect different variables.  
 540 As might be expected, river runoff is more important for salinity than for temperature. However,  
 541 for July-August, replacing rivers with the monthly-mean river flow gives nearly the same result  
 542 (in terms of variability) as fully time-varying rivers. For the duration considered (July-August),  
 543 averaging the river runoff gives similar effect for salinity as compared to giving the observed  
 544 river runoff in the simulations, see Figure 10. Temperature is less sensitive to any of the forcing  
 545 alterations, because although temperature and salinity are passive tracers they have different sources  
 546 and sinks. Switching the wind product from NAM to NECOFS does not have any significant effect



492 FIG. 10. Shannon entropy applied to temperature and salinity. Replacing fully time varying rivers with  
 493 monthly-mean river flow gives almost the same result for salinity. Same is true by replacing wind product with a  
 494 different one. Rivers set to zero affects salinity but not temperature. Winds are important in terms of variability  
 495 but different wind products do not noticeably alter variability.

547 on the sources or sinks of temperature or salinity, but switching the wind off definitely affects  
 548 the parameters by eliminating wind-driven mixing altogether. Figure 11 shows that zero wind  
 549 (ZW) simulations are markedly different than the rest in terms of *mutual information* (i.e., they



496 FIG. 11. Mutual information applied to simulations from different forcings. Higher mutual information implies  
 497 higher similarity in terms of variability. For example NAM-NECOFS values are higher than NAM-ZW implying  
 498 that NAM and NECOFS are significantly different than having no wind.

500 do not covary), although very similar in terms of amount of spatial variability (Shannon entropy,  
 501 Figure 10), because even without winds tides, fluxes, and rivers still vary. The zero river case tends  
 502 to eliminate both variability and mutual information (ZR). Please note that our simulations are for  
 503 July-August, and results might be different for different season.

554 If we were to prioritize improvements based on Shannon entropy and mutual information, note  
555 that the two highest mutual information cases are where NAM is substituted with NECOFS and  
556 where mean rivers are substituted for varying rivers. The first observation is important from a  
557 forecast perspective, because it means that we can not easily tell the difference between different  
558 wind products, although something rather than zero winds should be used if the estuary needs to be  
559 forecasted for the full 20 day predictability range (weather forecasts are reliable for about 7 days in  
560 this region). Similarly, knowing that substituting the mean of the rivers for the fully varying rivers  
561 has little impact implies that rivers can be fixed in time for forecasts beyond where they might be  
562 predicted based on expected weather and precipitation. Finally, despite the fact that Narragansett  
563 Bay is a dominantly tidally-mixed estuary, among the sources of overall variability (i.e., sources  
564 of information entropy) considered, preserving an inflow of fresh water is key, even though that  
565 inflow can be steady. Winds do not appreciably increase information entropy of the Bay, but they  
566 are an important source of forced co-variation, and so are important for predictions but do not raise  
567 the overall level of variability.

## 568 **5. Conclusion**

569 We have proposed an information theory metric to determine contribution of intrinsic chaos and  
570 external variability to total variability in ensemble model simulations. Our metric uses Shannon  
571 entropy and mutual information and has several advantages over using only standard deviation (or  
572 variance). We have applied our metric on idealized Gaussian arrays as well as realistic coastal  
573 ocean and global climate model. We conclude that:

- 574 1. The new information theory metric is more reliable when outliers are present, because out-  
575 liers get assigned less probability and because Gaussian distributions have a difficult time  
576 approximating long-tailed (i.e., outlier prone) distributions.
- 577 2. The new information theory metric is more reliable when variability is non-Gaussian because  
578 it is based on non-parametric measures of the probability distributions.
- 579 3. The new information theory metric varies monotonically with ensemble member to ensemble  
580 mean correlation, but is quantified in fraction of bits required to capture internal variability  
581 versus bits required to capture of total variability.

- 582 4. The use of the information theory metric in a coastal ocean model ensemble and a climate  
583 model ensemble qualitatively changes the focus to regions that were previously erroneously  
584 labeled as having high or low internal variability.
- 585 5. In this case, the coastal ensemble had a much smaller intrinsic (chaotic) proportion of its  
586 total variability in comparison to the climate ensemble had more intrinsic (weather, climate  
587 oscillations, etc.) as a proportion of its total. Importantly, the resolution of the models helps  
588 determine the proportion of intrinsic variability, so such comparisons are model-specific:  
589 a higher resolution coastal model might well have a larger intrinsic fraction than a coarser  
590 climate model.

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594 for Computation and Visualization (Brown University), supported in part by the National Science  
595 Foundation EPSCoR Cooperative Agreement #OIA-1655221.

596 *Data availability statement.* All the data and the codes used to plot results can be downloaded  
597 via Brown University's digital archive DOI: [urlplaceholder](#).

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