# Semiclassical states for fractional Schrödinger equations with critical nonlinearities 

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#### Abstract

In this paper, we consider the following fractional Schrödinger equation $\varepsilon 2 \mathrm{~s}(-[?]) \mathrm{s} u+\mathrm{V}(\mathrm{x}) \mathrm{u}=\mathrm{P}(\mathrm{x}) \mathrm{f}(\mathrm{u})+\mathrm{Q}(\mathrm{x})$ $|\mathrm{u}| 2 \mathrm{~s}^{*}-2 \mathrm{u}$ in R N , where $\epsilon>0$ is a parameter, $s[?](0,1), 2 \mathrm{~s} *=2 \mathrm{~N} \mathrm{~N}-2 \mathrm{~s}, N>2 s,(-[?]) \mathrm{s}$ is the fractional Lapalacian and $f$ is a superlinear and subcritical nonlinearity. Under a local condition imposed on the potential function, combining the penalization method and the concentration-compactness principle, we prove the existence of a positive solution for the above equations.


# Semiclassical states for fractional Schrödinger equations with critical nonlinearities * 

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Abstract: In this paper, we consider the following fractional Schrödinger equation

$$
\varepsilon^{2 s}(-\Delta)^{s} u+V(x) u=P(x) f(u)+Q(x)|u|^{2_{s}^{*}-2} u \quad \text { in } \mathbb{R}^{N}
$$

where $\varepsilon>0$ is a parameter, $s \in(0,1), 2_{s}^{*}=\frac{2 N}{N-2 s}, N>2 s,(-\Delta)^{s}$ is the fractional Lapalacian and $f$ is a superlinear and subcritical nonlinearity. Under a local condition imposed on the potential function, combining the penalization method and the concentration-compactness principle, we prove the existence of a positive solution for the above equations.

Key words: Concentration-compactness principle; Penalization method; Nonautonomous nonlinearities; fractional Schrödinger equation.

## 1 Introduction

We consider the following fractional Schrödinger equation

$$
\begin{cases}\varepsilon^{2 s}(-\Delta)^{s} u+V(x) u=P(x) f(u)+Q(x)|u|^{2_{s}^{*}-2} u, & x \in \mathbb{R}^{N}, \\ u \in H^{s}\left(\mathbb{R}^{N}\right), u(x)>0 & x \in \mathbb{R}^{N},\end{cases}
$$

where $\varepsilon>0$ is a small parameter, $s \in(0,1), 2_{s}^{*}:=\frac{2 N}{N-2 s}$ is the fractional critical exponent, the function $f$ is a superlinear and subcritical nonlinearity. Here the factional Sobolev space $H^{s}\left(\mathbb{R}^{N}\right)$ is defined by

$$
H^{s}\left(\mathbb{R}^{N}\right)=\left\{u \in L^{2}\left(\mathbb{R}^{N}\right): \int_{\mathbb{R}^{2 N}} \frac{(u(x)-u(y))^{2}}{|x-y|^{N+2 s}} d x d y<\infty\right\}
$$

equipped with the norm

$$
\|u\|_{H^{s}\left(\mathbb{R}^{N}\right)}=\left(\int_{\mathbb{R}^{N}}|u(x)|^{2} d x+\int_{\mathbb{R}^{2 N}} \frac{(u(x)-u(y))^{2}}{|x-y|^{N+2 s}} d x d y\right)^{1 / 2}
$$

[^0]$(-\Delta)^{s}$ stands for the usual fractional Laplacian, $(-\Delta)^{s}$ of a smooth function $u: \mathbb{R}^{N} \rightarrow$ $\mathbb{R}$ is defined by
$$
\mathcal{F}\left((-\Delta)^{s}(u)\right)(\xi)=|\xi|^{2 s} \mathcal{F}(u)(\xi), \xi \in \mathbb{R}^{N}
$$
where $\mathcal{F}$ denotes the Fourier transform, that is,
$$
\mathcal{F}(w)(\xi)=\frac{1}{(2 \pi)^{N / 2}} \int_{\mathbb{R}^{N}} e^{-i \xi \cdot x} w(x) d x
$$
for function $w$ in the Schwartz class. Also, $(-\Delta)^{s} u$ can be equivalently represented as
$$
(-\Delta)^{s} u=-\frac{C_{N, s}}{2} \int_{\mathbb{R}^{N}} \frac{(u(x+y)+u(x-y)-2 u(x))}{|y|^{N+2 s}} d y, \forall x \in \mathbb{R}^{N}
$$
where $C_{N, s}>0$ is the normalizing constant, defined by
$$
C_{N, s}=\left(\int_{\mathbb{R}^{N}} \frac{1-\cos \xi_{1}}{|\xi|^{N+2 s}} d \xi\right)^{-1}, \xi=\left(\xi_{1}, \xi_{2}, \cdots, \xi_{N}\right)
$$

We have from [9] that

$$
\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} u(x)\right|^{2} d x=\frac{C_{N, s}}{2} \int_{\mathbb{R}^{2 N}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 s}} d x d y
$$

And by taking derivative of the above equality, for any $u, v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{2 N}} \frac{(u(x)-u(y))(v(x)-v(y))}{|x-y|^{N+2 s}} d x d y=C_{N, s} \int_{\mathbb{R}^{N}}(-\Delta)^{\frac{s}{2}} u(x)(-\Delta)^{\frac{s}{2}} v(x) d x . \tag{1.1}
\end{equation*}
$$

Moreover, for any $u, v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ we have

$$
\begin{equation*}
(-\Delta)^{\frac{s}{2}}(u v)=u(-\Delta)^{\frac{s}{2}} v+v(-\Delta)^{\frac{s}{2}} u-2 I_{\frac{s}{2}}(u, v) \tag{1.2}
\end{equation*}
$$

where $I_{\frac{s}{2}}$ is defined in the principal value sense, as follows

$$
I_{\frac{s}{2}}(u, v)(x)=P . V . \int_{\mathbb{R}^{N}} \frac{(u(x)-u(y))(v(x)-v(y))}{|x-y|^{N+s}} d y .
$$

Problem $\left(\widehat{P_{\varepsilon}}\right)$ describes the so called standing waves of the nonlinear, time-dependent fractional Schrödinger equation of the form

$$
\begin{equation*}
i \varepsilon \frac{\partial \psi}{\partial t}=\varepsilon^{2 s}(-\Delta)^{s} \psi+V(x) \psi-f(x, \psi) \tag{1.3}
\end{equation*}
$$

Solutions of (1.3) for sufficiently small $\varepsilon>0$ are called semiclassical states. Recently great attention has been devoted to the study of semiclassical states, see for example [1.4, 7, 12, 13, 15, 18, 19] and the references therein and most of them assume that the potential satisfies the following global condition
( $V$ ) $V \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$ and $0<\inf _{x \in \mathbb{R}^{N}} V(x)<\liminf _{|x| \rightarrow+\infty} V(x)=V_{\infty}<+\infty$,
which is first introduced by Rabinowtz in [17] in the study of a nonlinear Schrödinger equation with the nonlinear subcritical growth. There are some results for problem $\left(P_{\varepsilon}\right.$ when $V(x)$ satisfies the following local condition
$\left(V_{1}\right) V \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$ and there is constant $V_{0}>0$ such that $V_{0}:=\inf _{x \in \mathbb{R}^{N}} V(x) ;$
$\left(V_{2}\right)$ there is a bounded open set $\Omega \subset \mathbb{R}^{N}$ such that $V_{0}<\min _{\partial \Omega} V$, and $M:=\{x \in \Omega$ : $\left.V(x)=V_{0}\right\} \neq \emptyset$.

Some authors have studied the existence and concentration phenomena for potentials verifying local condition $\left(V_{1}\right)$ and $\left(V_{2}\right)$, see for example [2-4, 7, 13, 15] and the references therein. As far as we know, all of them concentrate on the problems with autonomous nonlinearities. Particularly, in [13] the authors consider the following equation

$$
\begin{equation*}
\varepsilon^{2 s}(-\Delta)^{s} u+V(x) u=f(u)+u^{2_{s}^{*}-1} \quad x \in \mathbb{R}^{N} \tag{1.4}
\end{equation*}
$$

under the conditions $\left(V_{1}\right)$ and $\left(V_{2}\right)$ and obtain the existence and concentration of multiple solutions, which concentrate on the minima of $V(x)$ as $\varepsilon \rightarrow 0$. Our aim is to study the existence and concentration of positive solutions for problem ( $\overline{P_{\varepsilon}}$ ) by combining a local assumption on $V$, and show that the penalization method introduced by del Pino and Felmer in [8] can be also applied to a general class of problems with nonautonomous nonlinearities.

Below we give some assumptions. Since we are interested in positive solutions, we assume that $f \in C^{1}(\mathbb{R}, \mathbb{R})$ vanishes in $(-\infty, 0)$ and satisfies the following conditions.
$\left(f_{1}\right) f(t)=o(t)$ as $t \rightarrow 0^{+}$.
$\left(f_{2}\right)$ There exist constants $q, \sigma \in\left(2,2_{s}^{*}\right)$, sufficiently large $C_{0}>0$ such that $f(t) \geq$ $C_{0} t^{q-1}$ for all $t \geq 0$, and $\lim _{t \rightarrow \infty} \frac{f(t)}{t^{t-1}}=0$.
$\left(f_{3}\right)$ There exists a constant $\theta \in\left(2,2_{s}^{*}\right)$ such that for all $t>0,0<\theta F(t):=$ $\theta \int_{0}^{t} f(\tau) d \tau \leq t f(t)$.
$\left(f_{4}\right)$ The function $\frac{f(t)}{t}$ is increasing on interval $(0, \infty)$.
The potential $V(x)$ satisfies $\left(V_{1}\right), P(x)$ and $Q(x)$ are assumed to satisfy the following conditions.
(P) $P \in L^{\infty}\left(\mathbb{R}^{N}\right)$ is continuous and there is a constant $\alpha>0$ such that $P(x) \geq \alpha$ for all $x \in \mathbb{R}^{N}$.
( $Q$ ) $Q \in L^{\infty}\left(\mathbb{R}^{N}\right)$ is continuous and there is a constant $\beta>0$ such that $Q(x) \geq \beta$ for all $x \in \mathbb{R}^{N}$.
$(\Omega)$ There is a bounded, nonempty domain $\Omega \subset \mathbb{R}^{N}$ such that
$\left(\Omega_{1}\right)$ there is $x_{\text {min }} \in \Omega$ such that $V\left(x_{\text {min }}\right)=V_{0}<\min _{\partial \Omega} V, P\left(x_{\text {min }}\right)=P_{0}=\sup _{\mathbb{R}^{N}} P$ and $Q\left(x_{\text {min }}\right)=Q_{0}=\sup _{\mathbb{R}^{N}} Q$,
or
$\left(\Omega_{2}\right)$ there is $x_{\max } \in \Omega$ such that $P\left(x_{\max }\right)=P_{0}>\max _{\partial \Omega} P, Q\left(x_{\max }\right)=Q_{0}>\max _{\partial \Omega} Q$ and $V\left(x_{\max }\right)=V_{0}$.

The main result of this paper is stated as follows:
Theorem 1.1. Assume that $\left(V_{1}\right),(P),(Q),(\Omega)$ and $\left(f_{1}\right)-\left(f_{4}\right)$ hold. Then there exists $\varepsilon_{0}>0$ for any $\varepsilon \in\left(0, \varepsilon_{0}\right)$, problem ( $P_{\varepsilon}$ has a positive solution $u_{\varepsilon}$. Furthermore, if $\eta_{\varepsilon} \in \mathbb{R}^{N}$ denotes its global maximum point, then

$$
u_{\varepsilon}(x) \leq \frac{C \varepsilon^{N+2 s}}{\varepsilon^{N+2 s}+\left|x-\eta_{\varepsilon}\right|^{N+2 s}}
$$

Remark 1.2. To our best knowledge, the existence and qualitative properties of solutions for problem $\left(P_{\varepsilon}\right)$ have been extensively studied when $V(x)$ satisfies the global condition $(V)$. There are few results for problem $\left(\overline{P_{\varepsilon}}\right)$ when $V(x)$ satisfies a local condition as above, even in the $P=Q=1$ case. Under a local condition imposed on $V$, it is necessary to create a penalization function. If $P \neq 1, Q \neq 1$, the construction of penalization function is more complicated. Especially after adding the critical nonlinearities, the problem is more difficult, so far, no one has studied this aspect. Motivated by the penalization approach used in [8], we will investigate the existence of positive solution for problem $\left(P_{\varepsilon}\right.$ by supposing that $V$ satisfies a local assumption as above. Hence, our results can be seen as an improvement and supplement to [2, 4, 7, 13, 15].

Remark 1.3. Compared with the previous works, the main difficulty lies in the nonautonomous nonlinearity with the critical Sobolev growth and the potential $V$ with a local assumption, which makes it more complicated to recover the compactness. So we focus on the essential difficulty of the problem under the assumption that $f \in C^{1}\left(\mathbb{R}^{N}, \mathbb{R}\right)$. We note that if $f \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$, Theorem 1.1 also holds true by using the method of [20] under the condition of this paper. The related specific proof can be found in [13, 21].

To establish the existence of positive solution, we will use the penalization method introduced by Del Pino and Felmer [8]. First, we need to fix some notations.

Let $K, L>0, \frac{P_{0}}{K}+\frac{Q_{0}}{L}<\min \left\{\frac{1}{2}, \frac{\theta-2}{\theta}\right\}$, and $a>0$ such that $f(a)=\frac{V_{0}}{K} a$ and $a^{2_{s}^{*}-1}=\frac{V_{0}}{L} a$, where $\theta$ and $V_{0}$ are introduced in $\left(f_{3}\right)$ and $\left(V_{1}\right)$ respectively. We set

$$
\tilde{f}(t)=\left\{\begin{array}{ll}
f(t), & \text { if } t \leq a \\
\frac{V_{0}}{K} t, & \text { if } t>a
\end{array}, \quad \tilde{g}(t)= \begin{cases}t^{2 s}-1, & \text { if } t \leq a \\
\frac{V_{0}}{L} t, & \text { if } t>a,\end{cases}\right.
$$

and

$$
\begin{equation*}
g(x, t)=\chi_{\Omega}(x)\left(P(x) f(t)+Q(x) t^{2_{s}^{*}-1}\right)+\left(1-\chi_{\Omega}(x)\right)(P(x) \tilde{f}(t)+Q(x) \tilde{g}(t)) \tag{1.5}
\end{equation*}
$$

where $\chi_{\Omega}$ is the charateristic function of the set $\Omega$. Form $\left(f_{1}\right)-\left(f_{4}\right)$, it is easy to check that $g$ satisfies the following properties,
$\left(g_{1}\right) \lim _{t \rightarrow 0^{+}} \frac{g(x, t)}{t}=0$ uniformly in $x \in \mathbb{R}^{N}$.
$\left(g_{2}\right) g(x, t) \leq P_{0} f(t)+Q_{0} t^{2_{s}^{*}-1}$ for all $x \in \mathbb{R}^{N}, t>0$.
$\left(g_{3}\right)$ There is $\theta \in\left(2,2_{s}^{*}\right)$ such that

$$
0 \leq \theta G(x, t):=\theta \int_{0}^{t} g(x, s) d s<g(x, t) t \text { for } \forall x \in \Omega \text { and } \forall t>0
$$

and

$$
0 \leq 2 G(x, t)<g(x, t) t \leq\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right) V(x) t^{2} \text { for } \forall x \in \mathbb{R}^{N} \backslash \Omega, \text { and } \forall t>0 .
$$

$\left(g_{4}\right)$ For each $x \in \Omega$, the function $t \rightarrow \frac{g(x, t)}{t}$ is increasing in interval $(0, \infty)$ and for each $x \in \mathbb{R}^{N} \backslash \Omega$, the function $t \rightarrow \frac{g(x, t)}{t}$ is increasing in $(0, a)$.

Now we study the modified problem

$$
\left\{\begin{array}{ll}
\varepsilon^{2 s}(-\Delta)^{s} u+V(x) u=g(x, u), & x \in \mathbb{R}^{N}  \tag{*}\\
u \in H^{s}\left(\mathbb{R}^{N}\right), u(x)>0, & x \in \mathbb{R}^{N}
\end{array} .\right.
$$

Note that the positive solution of $\left(P_{\varepsilon}^{*}\right)$ with $u(x) \leq a$ for each $x \in \mathbb{R}^{N} \backslash \Omega$ is also the positive solution of $\left(P_{\varepsilon}\right)$.

In view of the presence of potential $V(x)$, we introduce the following fractional Sobolev space

$$
H_{\varepsilon}:=\left\{u \in H^{s}\left(\mathbb{R}^{N}\right): \int_{\mathbb{R}^{N}} V(x) u^{2} d x<\infty\right\}
$$

endowed with the norm

$$
\|u\|_{\varepsilon}^{2}=\int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+V(x) u^{2}\right) d x .
$$

Consider the energy functional $J_{\varepsilon}: H_{\varepsilon} \rightarrow \mathbb{R}$ associated to ( $\left(P_{\varepsilon}^{*}\right)$ given by

$$
J_{\varepsilon}(u)=\frac{\varepsilon^{2 s}}{2} \int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} u\right|^{2} d x+\frac{1}{2} \int_{\mathbb{R}^{N}} V(x) u^{2} d x-\int_{\mathbb{R}^{N}} G(x, u) d x,
$$

and its Nehari manifold is defined by

$$
\mathcal{N}_{\varepsilon}:=\left\{u \in H_{\varepsilon} \backslash\{0\}:\left\langle J_{\varepsilon}^{\prime}(u), u\right\rangle=0\right\} .
$$

## 2 Proof of the main result

We only discuss the case that $V, P, Q$ satisfy $\left(\Omega_{1}\right)$, and when $V, P, Q$ satisfy $\left(\Omega_{2}\right)$ the proof of the conclusion is similar to that of the case $\left(\Omega_{1}\right)$.

### 2.1 The autonomous problem $P_{V_{0}}$

We start by considering the autonomous problem associated to $P_{\varepsilon}$, namely,

$$
\begin{equation*}
(-\Delta)^{s} u+V_{0} u=P_{0} f(u)+Q_{0}|u|^{2_{s}^{*}-2} u . \tag{0}
\end{equation*}
$$

The solutions of problem $\left(P_{V_{0}}\right)$ are precisely the positive critical points of the functional defined by

$$
J_{0}(u)=\frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+V_{0} u^{2}\right) d x-\int_{\mathbb{R}^{N}} P_{0} F(u) d x-\frac{Q_{0}}{2_{s}^{*}} \int_{\mathbb{R}^{N}}|u|^{2_{s}^{*}} d x,
$$

and its Nehari manifold is defined by

$$
\mathcal{N}_{0}:=\left\{u \in H_{0} \backslash\{0\}:\left\langle J_{0}^{\prime}(u), u\right\rangle=0\right\} .
$$

where $H_{0}:=\left\{u \in H^{s}\left(\mathbb{R}^{N}\right): \int_{\mathbb{R}^{N}} V_{0} u^{2} d x<\infty\right\}$,
The functional $J_{0}$ satisfies the mountain pass geometry, the proof is standard, and hence, it is omitted. By using a version of the mountain pass theorem without (PS) condition [22], it follows that there exists a sequence $\left\{u_{n}\right\} \subset H_{0}$ such that $J_{0}\left(u_{n}\right) \rightarrow c_{0}$ and $J_{0}^{\prime}\left(u_{n}\right) \rightarrow 0$, and $c_{0}:=\inf _{\gamma \in \Gamma_{0}} \max _{t \in[0,1]} J_{0}(\gamma(t))>0$, where $\Gamma_{0}:=\left\{\gamma \in C\left([0,1], H_{0}\right): \gamma(0)=0\right.$ and $\left.J_{0}(\gamma(1))<0\right\}$. Similarly to the arguments in [17], by $\left(f_{4}\right)$, the equivalent characterization of $c_{0}$ is given by

$$
c_{0}=\inf _{u \in H_{0} \backslash\{0\}} \sup _{t \geq 0} J_{0}(t u)=\inf _{u \in \mathcal{N}_{0}} J_{0}(u) .
$$

The following lemma gives the estimate of the critical value $c_{0}$.
Lemma 2.1. Suppose that $\left(f_{1}\right)-\left(f_{4}\right)$ hold, then

$$
0<c_{0}<\frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}},
$$

where $S_{*}$ is the best Sobolev constant

$$
S_{*}=\inf _{u \in H^{s}\left(\mathbb{R}^{N}\right) \backslash\{0\}} \frac{\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} u\right|^{2} d x}{\left(\int_{\mathbb{R}^{N}}|u|^{2_{s}^{*}} d x\right)^{\frac{2}{2_{s}^{*}}}} .
$$

Proof. We define

$$
\tilde{u}_{\varepsilon}(x)=\psi(x) U_{\varepsilon}(x), x \in \mathbb{R}^{N},
$$

where $U_{\varepsilon}(x)=\varepsilon^{-\frac{N-2 s}{2}} u^{*}\left(\frac{x}{\varepsilon}\right), u^{*}(x)=\frac{\bar{u}\left(x / S_{2}^{\frac{1}{2 s}}\right)}{|\bar{u}|_{2 s}^{*}}$, where $\bar{u}(x)=\kappa\left(\mu^{2}+\left|x-x_{0}\right|^{2}\right)^{-\frac{N-2 s}{2}}$, with $\kappa \in \mathbb{R} \backslash\{0\}, \mu>0$, and $x_{0} \in \mathbb{R}^{N}$, and $\psi \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ such that $0 \leq \psi \leq 1$ in $\mathbb{R}^{N}$, $\psi(x) \equiv 1$ in $B_{r}(0)$, and $\psi \equiv 0$ in $\mathbb{R}^{N} \backslash B_{2 r}(0)$. Define $v_{\varepsilon}(x)=\frac{\tilde{u}_{\varepsilon}(x)}{\mid\left\{\left.\tilde{u}_{\varepsilon}(x)\right|_{2_{\varepsilon}^{*}}\right.}$, from [13], we have the following estimates:

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}(x)\right|^{2} d x \leq S_{*}+O\left(\varepsilon^{N-2 s}\right) \tag{2.1}
\end{equation*}
$$

and

$$
\begin{gather*}
\int_{\mathbb{R}^{N}}\left|v_{\varepsilon}(x)\right|^{2} d x= \begin{cases}O\left(\varepsilon^{2 s}\right), & N>4 s \\
O\left(\varepsilon^{2 s} \ln \frac{1}{\varepsilon}\right), & N=4 s, \\
O\left(\varepsilon^{N-2 s}\right), & N<4 s\end{cases}  \tag{2.2}\\
\int_{\mathbb{R}^{N}}\left|v_{\varepsilon}(x)\right|^{q} d x \geq C \varepsilon^{\frac{2 N-(N-2 s) q}{2}} . \tag{2.3}
\end{gather*}
$$

By the definition of $v_{\varepsilon}$ and $\left(f_{2}\right)$, we have

$$
\begin{aligned}
J_{0}\left(t v_{\varepsilon}\right) & =\frac{t^{2}}{2} \int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\frac{t^{2}}{2} \int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x-\int_{\mathbb{R}^{N}} P_{0} F\left(t v_{\varepsilon}\right) d x-\frac{Q_{0}}{2_{s}^{*}} t^{2_{s}^{*}} \\
& \leq \frac{t^{2}}{2} \int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\frac{t^{2}}{2} \int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x-C_{0} P_{0} t^{q} \int_{\mathbb{R}^{N}}\left|v_{\varepsilon}\right|^{q} d x-\frac{Q_{0}}{2_{s}^{*}} t^{2_{s}^{*}} .
\end{aligned}
$$

We consider the function

$$
\begin{equation*}
g(t)=\frac{t^{2}}{2} \int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\frac{t^{2}}{2} \int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x-C_{0} P_{0} t^{q} \int_{\mathbb{R}^{N}}\left|v_{\varepsilon}\right|^{q} d x-\frac{Q_{0}}{2_{s}^{*}} t^{2_{s}^{*}} . \tag{2.4}
\end{equation*}
$$

It is clear that $g(t)>0$ for $t>0$ small enough, and $g(t) \rightarrow-\infty$ as $t \rightarrow+\infty$. Hence there exists $t_{\varepsilon}>0$ such that $\max _{t \geq 0} g(t)=g\left(t_{\varepsilon}\right)$, and
$0=g^{\prime}\left(t_{\varepsilon}\right)=t_{\varepsilon}\left(\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x-q C_{0} P_{0} t_{\varepsilon}^{q-2} \int_{\mathbb{R}^{N}}\left|v_{\varepsilon}\right|^{q} d x-Q_{0} t_{\varepsilon}^{t_{s}^{*}-2}\right)$.
Therefore,

$$
\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x=q C_{0} P_{0} t_{\varepsilon}^{q-2} \int_{\mathbb{R}^{N}}\left|v_{\varepsilon}\right|^{q} d x+Q_{0} t_{\varepsilon}^{2_{s}^{*}-2},
$$

which implies

$$
0<t_{\varepsilon} \leq \frac{1}{Q_{0}^{\frac{1}{2_{s}^{s}-2}}}\left(\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x\right)^{\frac{1}{2_{s}^{\frac{1}{s}-2}}}
$$

It is easy to verifies that $J_{0}$ satisfies the mountain-pass geometry conditions, and we get

$$
0<\delta \leq J_{0}\left(t_{\varepsilon} v_{\varepsilon}\right) \leq \frac{t_{\varepsilon}^{2}}{2}\left(\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x\right) .
$$

Hence we have a lower bound and a upper bound for $t_{\varepsilon}$, independent of $\varepsilon$, let

$$
\bar{g}(t)=\frac{t^{2}}{2}\left(\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x\right)-\frac{Q_{0}}{2_{s}^{*}} t^{2_{s}^{*}},
$$

then $t_{\varepsilon}=\frac{1}{Q_{0}^{\frac{1}{2_{s}^{x}-2}}}\left(\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2}+V_{0}\left|v_{\varepsilon}\right|^{2} d x\right)^{\frac{1}{2_{s}^{\frac{1}{s}-2}}}$ is the maximum point of $\bar{g}(t)$. Hence, by (2.1)-(2.3), and the elementary inequality $(a+b)^{p} \leq a^{p}+p(a+b)^{p-1} b$ for $a, b>0$ and $p \geq 1$, we obtain

$$
g\left(t_{\varepsilon}\right)=\bar{g}\left(t_{\varepsilon}\right)-C_{0} P_{0} t_{\varepsilon}^{q} \int_{\mathbb{R}^{N}}\left|v_{\varepsilon}\right|^{q} d x
$$

$$
\begin{aligned}
& \leq \bar{g}\left(\frac{1}{Q_{0}^{\frac{1}{2 s-2}}}\left(\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2}+V_{0}\left|v_{\varepsilon}\right|^{2}\right) d x\right)^{\frac{1}{2 s-2}}\right)-C_{0} P_{0} t_{\varepsilon}^{q} \int_{\mathbb{R}^{N}}\left|v_{\varepsilon}\right|^{q} d x \\
& \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}}\left(\int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x\right)^{\frac{N}{2 s}}-C_{0} C\left|v_{\varepsilon}\right|_{q}^{q} \\
& \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}}\left(S_{*}+O\left(\varepsilon^{N-2 s}\right)+\int_{\mathbb{R}^{N}} V_{0}\left|v_{\varepsilon}\right|^{2} d x\right)^{\frac{N}{2 s}}-C_{0} C\left|v_{\varepsilon}\right|_{q}^{q} \\
& \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}+O\left(\varepsilon^{N-2 s}\right)+C\left|v_{\varepsilon}\right|_{2}^{2}-C_{0} C\left|v_{\varepsilon}\right|_{q}^{q} .
\end{aligned}
$$

Next we distinguish the following cases.
(i) If $N>4 s$, then $\frac{N}{N-2 s}<2$, we have $q>\frac{N}{N-2 s}$, by (2.2) and (2.3) we get

$$
\sup _{t>0} g(t) \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}+O\left(\varepsilon^{N-2 s}\right)+O\left(\varepsilon^{2 s}\right)-O\left(\varepsilon^{\frac{2 N-(N-2 s) q}{2}}\right)
$$

since $\frac{2 N-(N-2 s) q}{2}<2 s<N-2 s$, we get the conclusion for $\varepsilon$ sufficiently small.
(ii) If $N=4 s$, then $2<q<2_{s}^{*}=4$, by (2.2) and (2.3) we obtain

$$
\begin{aligned}
\sup _{t>0} g(t) & \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}+O\left(\varepsilon^{N-2 s}\right)+O\left(\varepsilon^{2 s} \ln \frac{1}{\varepsilon}\right)-O\left(\varepsilon^{4 s-s q}\right) \\
& \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}} S_{*}^{\frac{N}{2 s}}+O\left(\varepsilon^{2 s}\left(1+\ln \frac{1}{\varepsilon}\right)\right)-O\left(\varepsilon^{4 s-s q}\right)} \\
& <\frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}} S_{*}^{\frac{N}{2 s}} .}
\end{aligned}
$$

Since

$$
\lim _{\varepsilon \rightarrow 0^{+}} \frac{\varepsilon^{4 s-s q}}{\varepsilon^{2 s}\left(1+\ln \frac{1}{\varepsilon}\right)}=+\infty
$$

we get the conclusion for $\varepsilon$ sufficiently small.
(iii) If $2 s<N<4 s$ and $\frac{N}{N-2 s}<q<2_{s}^{*}$, by (2.2) and (2.3) we have

$$
\sup _{t>0} g(t) \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}+O\left(\varepsilon^{N-2 s}\right)-O\left(\varepsilon^{\frac{2 N-(N-2 s) q}{2}}\right)
$$

In view of $\frac{2 N-(N-2 s) q}{2}<N-2 s$, we get the conclusion for $\varepsilon$ sufficiently small.
(iv) If $2 s<N<4 s$ and $2<q \leq \frac{N}{N-2 s}$, from (2.2) and (2.3) we obtain

$$
\sup _{t>0} g(t) \leq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}+O\left(\varepsilon^{N-2 s}\right)-C_{0} O\left(\varepsilon^{\frac{2 N-(N-2 s) q}{2}}\right),
$$

and for $C_{0}=\varepsilon^{-\theta}$ with $\theta>\frac{4 s-(N-2 s) q}{2}$, we also get the conclusion. Hence, $c_{0}<$ $\frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}$.

Lemma 2.2. Assume that $\left\{u_{n}\right\} \subset H_{0}$ is a $(P S)_{c}$ sequence for $J_{0}$ with $c<\frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}$ and such that $u_{n} \rightharpoonup 0$. Then, one of the following alternatives occurs:
(a) $u_{n} \rightarrow 0$ in $H_{0}$, or
(b) There exists a sequence $\left\{z_{n}\right\} \subset \mathbb{R}^{N}$ and constants $R, \eta>0$ such that

$$
\liminf _{n \rightarrow \infty} \int_{B_{R}\left(z_{n}\right)}\left|u_{n}\right|^{2} d x \geq \eta>0
$$

Proof. Suppose (b) is not satisfied. Then for any $R>0$, we have

$$
\lim _{n \rightarrow \infty} \sup _{z \in \mathbb{R}^{N}} \int_{B_{R}(z)}\left|u_{n}\right|^{2} d x=0
$$

Since $\left\{u_{n}\right\}$ is bounded, from [11], it follows that $u_{n} \rightarrow 0$ in $L^{p}\left(\mathbb{R}^{N}\right), \forall p \in\left(2,2_{s}^{*}\right)$. Hence,

$$
\int_{\mathbb{R}^{N}} f\left(u_{n}\right) u_{n} d x=\int_{\mathbb{R}^{N}} F\left(u_{n}\right) d x=o_{n}(1) .
$$

Moreover, from $J_{0}\left(u_{n}\right) \rightarrow c>0$ and $\left\langle J_{0}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \rightarrow 0$, we have that

$$
\begin{equation*}
\frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u_{n}\right|^{2}+V_{0}\left|u_{n}\right|^{2}\right) d x-\frac{Q_{0}}{2_{s}^{*}} \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2_{s}^{*}} d x \rightarrow c, \tag{2.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u_{n}\right|^{2}+V_{0}\left|u_{n}\right|^{2}\right) d x=Q_{0} \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2_{s}^{*}} d x+o_{n}(1) . \tag{2.6}
\end{equation*}
$$

Since $\left\{u_{n}\right\}$ is bounded, up to a subsequence, we get

$$
\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u_{n}\right|^{2}+V_{0}\left|u_{n}\right|^{2}\right) d x \rightarrow l \geq 0 .
$$

If $l>0$, by (2.5) and 2.6), we obtain $c=\frac{s}{N} l$. On the other hand,

$$
l \geq \int_{\mathbb{R}^{N}}\left|(-\Delta)^{\frac{s}{2}} u_{n}\right|^{2} d x \geq S_{*}\left(\left|u_{n}\right|_{2_{s}^{*}}^{2_{s}^{*}}\right)^{2 / 2_{s}^{*}}
$$

taking the limit as $n \rightarrow \infty$, it follows that $l \geq \frac{1}{Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}$. Hence,

$$
c=\frac{s}{N} l \geq \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}},
$$

which is a contradiction. Therefore, $l=0$. It implies that $u_{n} \rightarrow 0$ in $H_{0}$.
Lemma 2.3. Suppose that $\left(f_{1}\right)-\left(f_{4}\right)$ hold, then problem $P_{V_{0}}$ has a positive ground state solution.

Proof. Assume $\left\{u_{n}\right\} \subset H_{0}$ is a $(P S)_{c_{0}}$ sequence, from the condition of nonlinearity $f$, we can easily check that $\left\{u_{n}\right\}$ is bounded in $H_{0}$. Thus, up to a subsequence, $u_{n} \rightharpoonup u$ weakly in $H_{0}$. Moreover, $u$ is a critical point of $J_{0}$. If $u \neq 0$, it remains to show that $J_{0}(u)=c_{0}$. By Fatou's Lemma, we have that

$$
\begin{aligned}
c_{0} \leq & J_{0}(u)-\frac{1}{\theta}\left\langle J_{0}^{\prime}(u), u\right\rangle \\
= & \left(\frac{1}{2}-\frac{1}{\theta}\right) \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+V_{0} u^{2}\right) d x \\
& +\frac{1}{\theta} \int_{\mathbb{R}^{N}} P_{0}(f(u) u-\theta F(u)) d x+\left(\frac{1}{\theta}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{N}} Q_{0}|u|^{2_{s}^{*}} d x \\
\leq & \liminf _{n \rightarrow \infty}\left\{\left(\frac{1}{2}-\frac{1}{\theta}\right) I R\left(\left|(-\Delta)^{\frac{s}{2}} u_{n}\right|^{2}+V_{0} u_{n}^{2}\right) d x\right. \\
& \left.+\frac{1}{\theta} \int_{\mathbb{R}^{N}} P_{0}\left(f\left(u_{n}\right) u_{n}-\theta F\left(u_{n}\right)\right) d x+\left(\frac{1}{\theta}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{N}} Q_{0}\left|u_{n}\right|^{2_{s}^{*}} d x\right\} \\
= & \liminf _{n \rightarrow \infty}\left(J_{0}\left(u_{n}\right)-\frac{1}{\theta}\left\langle J_{0}^{\prime}\left(u_{n}\right) u_{n}\right\rangle\right) \\
= & c_{0} .
\end{aligned}
$$

Hence, we proved that $J_{0}(u)=c_{0}$.
Now, we consider the case $u=0$. In fact, $u_{n} \nrightarrow 0$ in $H_{0}$. If $u_{n} \rightarrow 0$ in $H_{0}$, we have $J_{0}\left(u_{n}\right) \rightarrow 0$, this is contradicts to $c_{0}>0$. By Lemma 2.2 , there exists a sequence $\left\{z_{n}\right\} \subset \mathbb{R}^{N}$ and constants $R, \eta>0$ such that

$$
\liminf _{n \rightarrow \infty} \int_{B_{R}\left(z_{n}\right)}\left|u_{n}\right|^{2} d x \geq \eta>0
$$

Set $w_{n}(x)=u_{n}\left(x+z_{n}\right)$, again $\left\{w_{n}\right\}$ is a $(P S)_{c_{0}}$ sequence of $J_{0}$. Thus $\left\{w_{n}\right\}$ is bounded in $H_{0}$, and there exists $w \in H_{0}$ such that $w_{n} \rightharpoonup w$ in $H_{0}$ with $w \neq 0$, then the conclusion follows as in the first case.

Using $-u^{-}$as a test function, we have

$$
0=\left\langle\left(J_{0}^{\prime}(u),-u^{-}\right\rangle=\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u^{-}\right|^{2}+V_{0}\left|u^{-}\right|^{2}\right) d x\right.
$$

This implies that $u^{-}=0$. Noting that $u \not \equiv 0$, we claim that $u>0$ in $\mathbb{R}^{N}$. In fact, if $u\left(x_{0}\right)=0$ for some $x_{0} \in \mathbb{R}^{N}$, then $(-\Delta)^{s} u\left(x_{0}\right)=0$ and by the representation formula

$$
(-\Delta)^{s} u(x)=-\frac{C_{N, s}}{2} \int_{\mathbb{R}^{N}} \frac{u(x+y)+u(x-y)-2 u(x)}{|y|^{N+2 s}} d y,
$$

we obtain that, at $x_{0}$,

$$
\int_{\mathbb{R}^{N}} \frac{u\left(x_{0}+y\right)+u\left(x_{0}-y\right)}{|y|^{N+2 s}} d y=0,
$$

yielding $u \equiv 0$, which leads a contradiction. The proof is completed.

### 2.2 The modified problem ( $P_{\varepsilon}^{*}$ )

In this section, we shall prove the existence of positive solution for the modified problem ( $P_{\varepsilon}^{*}$ ).

Similar to section 2.1, by the condition of nonlinearity $g$, we can easily prove that the functional $J_{\varepsilon}$ verifies the mountain pass geometry. Thus there exists a sequence $\left\{u_{n}\right\} \subset H_{\varepsilon}$ such that $J_{\varepsilon}\left(u_{n}\right) \rightarrow c_{\varepsilon}$ and $J_{\varepsilon}^{\prime}\left(u_{n}\right) \rightarrow 0$, where $c_{\varepsilon}:=\inf _{\gamma \in \Gamma_{\varepsilon}} \max _{t \in[0,1]} J_{\varepsilon}(\gamma(t))>$ 0 , where $\Gamma_{\varepsilon}:=\left\{\gamma \in C\left([0,1], H_{\varepsilon}\right): \gamma(0)=0\right.$ and $\left.J_{\varepsilon}(\gamma(1))<0\right\}$. As in the previous section 2.1, we have the equivalent characterization of $c_{\varepsilon}$

$$
c_{\varepsilon}=\inf _{u \in H_{\varepsilon} \backslash\{0\}} \sup _{t \geq 0} J_{\varepsilon}(t u)=\inf _{u \in \mathcal{N}_{\varepsilon}} J_{\varepsilon}(u) .
$$

The main feature of the modified functional is that it satisfies the local compactness condition, which can be stated in the following results.

Lemma 2.4. Let $c>0$ and $\left\{u_{n}\right\}$ be a $(P S)_{c}$ sequence for $J_{\varepsilon}$, then $\left\{u_{n}\right\}$ is bounded.
Proof. Assume that $\left\{u_{n}\right\}$ is a $(P S)_{c}$ sequence for $J_{\varepsilon}$, then $J_{\varepsilon}\left(u_{n}\right) \rightarrow c$, and $J_{\varepsilon}^{\prime}\left(u_{n}\right) \rightarrow 0$. Therefore, we have

$$
\begin{aligned}
c+o_{n}(1)= & J_{\varepsilon}\left(u_{n}\right)-\frac{1}{\theta}\left\langle J_{\varepsilon}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
= & \left(\frac{1}{2}-\frac{1}{\theta}\right) \int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{n}\right|^{2}+V(x) u_{n}^{2}\right) d x \\
& +\frac{1}{\theta} \int_{\mathbb{R}^{N}}\left[g\left(x, u_{n}\right) u_{n}-\theta G\left(x, u_{n}\right)\right] d x \\
\geq & \frac{\theta-2}{2 \theta}\left\|u_{n}\right\|_{\varepsilon}^{2}+\frac{1}{\theta} \int_{\mathbb{R}^{N} \backslash \Omega}\left[g\left(x, u_{n}\right) u_{n}-\theta G\left(x, u_{n}\right)\right] d x \\
\geq & \frac{\theta-2}{2 \theta}\left\|u_{n}\right\|_{\varepsilon}^{2}-\frac{\theta-2}{2 \theta} \int_{\mathbb{R}^{N} \backslash \Omega}\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right) V(x) u_{n}^{2} d x \\
\geq & \frac{\theta-2}{2 \theta}\left\|u_{n}\right\|_{\varepsilon}^{2}-\frac{\theta-2}{2 \theta}\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right)\left\|u_{n}\right\|_{\varepsilon}^{2} \\
= & \left(1-\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right)\right) \frac{\theta-2}{2 \theta}\left\|u_{n}\right\|_{\varepsilon}^{2} .
\end{aligned}
$$

From $\theta>2$ and $\frac{P_{0}}{K}+\frac{Q_{0}}{L}<\min \left\{\frac{1}{2}, \frac{\theta-2}{\theta}\right\}$, we have that $\left\{u_{n}\right\}$ is bounded.
Lemma 2.5. ([14]) Let $\left\{u_{n}\right\}$ be a bounded $(P S)_{c}$ sequence for $J_{\varepsilon}$, then for each $\zeta>0$, there is a number $R=R(\zeta)>0$ such that

$$
\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{N} \backslash B_{R}(0)}\left(\varepsilon^{2 s} \int_{\mathbb{R}^{N}} \frac{\left|u_{n}(x)-u_{n}(y)\right|^{2}}{|x-y|^{N+2 s}} d y+V(x)\left|u_{n}\right|^{2}\right) d x<\zeta .
$$

To establish the local compactness results for $J_{\varepsilon}$, we need an extension of a concentration compactness result proposed by Lions.

Lemma 2.6. ([16]) Let $\Omega \subseteq \mathbb{R}^{N}$ be an open subset and let $\left\{u_{n}\right\}$ be a sequence in $H^{s}(\Omega)$ weakly converging to $u$ as $n \rightarrow \infty$ and such that

$$
\left|(-\Delta)^{\frac{s}{2}} u_{n}(x)\right|^{2} d x \stackrel{*}{\rightharpoonup} \mu,\left|u_{n}(x)\right|^{2_{s}^{*}} d x \stackrel{*}{\rightharpoonup} \nu \text { in } \mathcal{M}\left(\mathbb{R}^{N}\right) .
$$

Then, either $u_{n} \rightarrow u$ in $L_{\text {loc }}^{2_{s}^{*}}\left(\mathbb{R}^{N}\right)$ or there exists a (at most countable) set of distinct points $\left\{x_{j}\right\}_{j \in J} \subset \bar{\Omega}$ and positive numbers $\left\{\nu_{j}\right\}_{j \in J}$ such that

$$
\begin{equation*}
\nu=|u(x)|^{2_{s}^{*}} d x+\sum_{j \in J} \nu_{j} \delta_{x_{j}} \tag{2.7}
\end{equation*}
$$

Moreover, if $\Omega$ is bounded, then there exists a positive measure $\tilde{\mu} \in \mathcal{M}\left(\mathbb{R}^{N}\right)$ with spt $\tilde{\mu} \subset \bar{\Omega}$ and positive numbers $\left\{\mu_{j}\right\}_{j \in J}$ such that

$$
\begin{equation*}
\mu=\left|(-\Delta)^{\frac{s}{2}} u(x)\right|^{2} d x+\tilde{\mu}+\sum_{j \in J} \mu_{j} \delta_{x_{j}}, \quad S_{*} \nu_{j}^{2 / 2_{s}^{*}} \leq \mu_{j} \tag{2.8}
\end{equation*}
$$

Lemma 2.7. The functional $J_{\varepsilon}$ satisfies the Palais-Smale condition at any level $c<$ $\varepsilon^{N} \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} \frac{N}{*}_{\frac{N}{2 s}}$.
Proof. Let $\left\{u_{n}\right\} \subset H_{\varepsilon}$ be such that $J_{\varepsilon}\left(u_{n}\right) \rightarrow c<\varepsilon^{N} \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}$, and $J_{\varepsilon}^{\prime}\left(u_{n}\right) \rightarrow 0$. Then from Lemma 2.4, we have that $\left\{u_{n}\right\}$ is bounded in $H_{\varepsilon}$. Since $\left\langle J_{\varepsilon}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \rightarrow 0$, we get

$$
\begin{equation*}
\left\|u_{n}\right\|_{\varepsilon}^{2}=\int_{\mathbb{R}^{N}} g\left(x, u_{n}\right) u_{n} d x+o_{n}(1) . \tag{2.9}
\end{equation*}
$$

Up to a subsquence, we may assume that

$$
\begin{cases}u_{n} \rightharpoonup u & \text { weakly in } H_{\varepsilon}  \tag{2.10}\\ u_{n} \rightarrow u & \text { strongly in } L_{L o c}^{s}\left(\mathbb{R}^{N}\right) \text { for any } s \in\left[1.2_{s}^{*}\right), \\ u_{n}(x) \rightarrow u(x) & \text { for a.e. } x \in \mathbb{R}^{N} .\end{cases}
$$

Hence it is standard to check that for any $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right) \subset H_{\varepsilon}$,

$$
\begin{align*}
\int_{\mathbb{R}^{N}} \varepsilon^{2 s}(-\Delta)^{\frac{s}{2}} u_{n}(-\Delta)^{\frac{s}{2}} \varphi d x & \rightarrow \int_{\mathbb{R}^{N}} \varepsilon^{2 s}(-\Delta)^{\frac{s}{2}} u(-\Delta)^{\frac{s}{2}} \varphi d x \\
\int_{\mathbb{R}^{N}} V(x) u_{n} \varphi d x & \rightarrow \int_{\mathbb{R}^{N}} V(x) u \varphi d x  \tag{2.11}\\
\int_{\mathbb{R}^{N}} g\left(x, u_{n}\right) \varphi d x & \rightarrow \int_{\mathbb{R}^{N}} g(x, u) \varphi d x
\end{align*}
$$

By (2.11), the density of $C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ in $H_{\varepsilon}$, and $J_{\varepsilon}^{\prime}\left(u_{n}\right) \rightarrow 0$, we obtain that the weak limit $u$ is a critical point of $J_{\varepsilon}$, then

$$
\begin{equation*}
\|u\|_{\varepsilon}^{2}=\int_{\mathbb{R}^{N}} g(x, u) u d x \tag{2.12}
\end{equation*}
$$

Next, we will prove two claims.
Claim $1 \lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}} g\left(x, u_{n}\right) u_{n} d x=\int_{\mathbb{R}^{N}} g(x, u) u d x$.
this claim, (2.9) and (2.12) imply that $\left\|u_{n}\right\|_{\varepsilon}^{2} \rightarrow\|u\|_{\varepsilon}^{2}$, which yields that $u_{n} \rightarrow u$ in $H_{\varepsilon}$.

To prove this claim, from Lemma 2.5, we know that: for each $\zeta>0$, there exists $R=R(\zeta)>0$ such that

$$
\begin{equation*}
\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{N} \backslash B_{R}(0)}\left(\varepsilon^{2 s} \int_{\mathbb{R}^{N}} \frac{\left|u_{n}(x)-u_{n}(y)\right|^{2}}{|x-y|^{N+2 s}} d y+V(x)\left|u_{n}\right|^{2}\right) d x<\zeta . \tag{2.13}
\end{equation*}
$$

By (2.13), $\left(g_{2}\right),\left(f_{1}\right),\left(f_{2}\right)$ and the Sobolev embedding, for $n$ large enough, we get that

$$
\begin{align*}
\int_{\mathbb{R}^{N} \backslash B_{R}(0)} g\left(x, u_{n}\right) u_{n} d x & \leq C_{1} \int_{\mathbb{R}^{N} \backslash B_{R}(0)}\left(u_{n}^{2}+u_{n}^{2_{s}^{*}}\right) d x \\
& \leq C_{2}\left(\zeta+\zeta_{s}^{2_{s}^{*} / 2}\right) \tag{2.14}
\end{align*}
$$

On the other hand, choosing $R$ large enough, we may assume that

$$
\int_{\mathbb{R}^{N} \backslash B_{R}(0)} g(x, u) u d x \leq \zeta .
$$

Therefore, from the last inequality and (2.14), we have

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N} \backslash B_{R}(0)} g\left(x, u_{n}\right) u_{n} d x=\int_{\mathbb{R}^{N} \backslash B_{R}(0)} g(x, u) u d x . \tag{2.15}
\end{equation*}
$$

By the definition of $g$, we obtain that

$$
g\left(x, u_{n}\right) u_{n} \leq P_{0} f\left(u_{n}\right) u_{n}+Q_{0} a^{2_{s}^{*}}+\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right) V_{0} u_{n}^{2} \forall x \in \mathbb{R}^{N} \backslash \Omega .
$$

Since the set $B_{R}(0) \cap\left(\mathbb{R}^{N} \backslash \Omega\right)$ is bounded, we can use the above estimates, $\left.\left(g_{1}\right),\left(g_{2}\right), 2.11\right)$ and Lebesgue's theorem to conclude that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{B_{R}(0) \cap\left(\mathbb{R}^{N} \backslash \Omega\right)} g\left(x, u_{n}\right) u_{n} d x=\int_{B_{R}(0) \cap\left(\mathbb{R}^{N} \backslash \Omega\right)} g(x, u) u d x . \tag{2.16}
\end{equation*}
$$

Claim $2 u_{n} \rightarrow u$ in $L^{2_{s}^{*}}(\Omega)$. If Claim 2 holds, by $\left(g_{2}\right),\left(f_{1}\right),\left(f_{2}\right),(2.11)$ and Lebesgue's theorem, we can obtain that

$$
\lim _{n \rightarrow \infty} \int_{B_{R}(0) \cap \Omega} g\left(x, u_{n}\right) u_{n} d x=\int_{B_{R}(0) \cap \Omega} g(x, u) u d x .
$$

Hence, Claim 1 follows from the above expression, (2.15) and (2.16).
It remains to prove Claim 2. By Phrokorovs theorem (see Bogachev [6], Theorem 8.6.2) we may suppose that there are positive measures $\mu, \nu$ such that

$$
\begin{equation*}
\left|(-\Delta)^{\frac{s}{2}} u_{n}(x)\right|^{2} d x \stackrel{*}{\rightharpoonup} \mu,\left|u_{n}(x)\right|^{2_{s}^{*}} d x \stackrel{*}{\rightharpoonup} \nu . \tag{2.17}
\end{equation*}
$$

Hence, by Lemma 2.6 we have an at most countable index set of distinct points $\left\{x_{j}\right\}_{j \in J}$, $\left\{\mu_{j}\right\}_{j \in J},\left\{\nu_{j}\right\}_{j \in J} \subset(0, \infty)$, and positive measures $\tilde{\mu}$ with support contained in $\Omega$ such that

$$
\begin{equation*}
\mu=\left|(-\Delta)^{\frac{s}{2}} u(x)\right|^{2} d x+\tilde{\mu}+\sum_{j \in J} \mu_{j} \delta_{x_{j}}, \quad \nu=|u(x)|^{2_{s}^{*}} d x+\sum_{j \in J} \nu_{j} \delta_{x_{j}} \text { and } S_{*} \nu_{j}^{2 / 2_{s}^{*}} \leq \mu_{j}, \tag{2.18}
\end{equation*}
$$

for all $j \in J$, where $\delta_{x_{j}}$ is the Dirac mass at $x_{j} \in \mathbb{R}^{N}$.
It suffices to show that $\left\{x_{j}\right\}_{j \in J} \cap \Omega=\emptyset$. If not, suppose that $x_{j} \in \Omega$ for some $j \in J$. For $\rho>0$, define the function $\psi_{\rho}(x):=\psi_{\rho}\left(\frac{x-x_{j}}{\rho}\right)$ where $\psi \in C_{0}^{\infty}\left(\overline{\mathbb{R}^{N}},[0,1]\right)$ is such that $\psi \equiv 1$ on $B_{\frac{1}{2}}(0), \psi \equiv 0$ on $\mathbb{R}^{N} \backslash B_{1}(0)$. We assume that $\rho$ is chosen in such way that the support of $\psi_{\rho}$ is contained in $\Omega$.

Since $\left\{u_{n} \psi_{\rho}\right\}$ is bounded, $\left\langle J_{\varepsilon}^{\prime}\left(u_{n}\right), u_{n} \psi_{\rho}\right\rangle=o_{n}(1)$, we have

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} \varepsilon^{2 s}(-\Delta)^{\frac{s}{2}} u_{n}(-\Delta)^{\frac{s}{2}} u_{n} \psi_{\rho} d x \leq \int_{\mathbb{R}^{N}} P_{0} f\left(u_{n}\right) u_{n} \psi_{\rho} d x+\int_{\mathbb{R}^{N}} Q_{0}\left|u_{n}\right|^{2_{s}^{*}} \psi_{\rho} d x+o_{n}(1) \tag{2.19}
\end{equation*}
$$

Using the fact that $\psi_{\rho}$ has compact support and $f$ has subcritical growth, we have

$$
\lim _{\rho \rightarrow 0} \lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}} P_{0} f\left(u_{n}\right) u_{n} \psi_{\rho} d x=\lim _{\rho \rightarrow 0} \int_{\mathbb{R}^{N}} P_{0} f(u) u \psi_{\rho} d x=0
$$

And, by (1.1) and (1.2), we can write

$$
\begin{align*}
& \int_{\mathbb{R}^{N}} \varepsilon^{2 s}(-\Delta)^{\frac{s}{2}} u_{n}(-\Delta)^{\frac{s}{2}} u_{n} \psi_{\rho} d x \\
= & \int_{\mathbb{R}^{N}} \varepsilon^{2 s} u_{n}(x)(-\Delta)^{\frac{s}{2}} u_{n}(x)(-\Delta)^{\frac{s}{2}} \psi_{\rho}(x) d x+\int_{\mathbb{R}^{N}} \varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{n}(x)\right|^{2} \psi_{\rho}(x) d x \\
& -2 \int_{\mathbb{R}^{N}} \varepsilon^{2 s}(-\Delta)^{\frac{s}{2}} u_{n}(x) \int_{\mathbb{R}^{N}} \frac{\left(u_{n}(x)-u_{n}(y)\right)\left(u_{n}(x) \psi_{\rho}(x)-u_{n}(y) \psi_{\rho}(y)\right)}{|x-y|^{N+s}} d x d y . \tag{2.20}
\end{align*}
$$

Now, we show that

$$
\begin{equation*}
\lim _{\rho \rightarrow 0} \lim _{n \rightarrow \infty}\left|\int_{\mathbb{R}^{N}} u_{n}(x)(-\Delta)^{\frac{s}{2}} u_{n}(x)(-\Delta)^{\frac{s}{2}} \psi_{\rho}(x) d x\right|=0 \tag{2.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{\rho \rightarrow 0} \lim _{n \rightarrow \infty}\left|\int_{\mathbb{R}^{N}}(-\Delta)^{\frac{s}{2}} u_{n}(x) \int_{\mathbb{R}^{N}} \frac{\left(u_{n}(x)-u_{n}(y)\right)\left(u_{n}(x) \psi_{\rho}(x)-u_{n}(y) \psi_{\rho}(y)\right)}{|x-y|^{N+s}} d x d y\right|=0 \tag{2.22}
\end{equation*}
$$

If (2.21) and (2.22) hold, we can use (2.17), 2.18), and take limits as $n \rightarrow \infty$ and $\rho \rightarrow 0$ in 2.20 to obtain that $Q_{0} \nu_{j} \geq \varepsilon^{2 s} \mu_{j}$. The proof of (2.21) and (2.22) is standard which can found in [5, Lemma 2.8 and Lemma 2.9], and we omit it. Then from the
last statement in (2.18), we get $\nu_{j} \geq \varepsilon^{N} \frac{1}{Q_{0}^{\frac{N}{2 s}}} S_{*}^{\frac{N}{2 s}}$, and hence we can use $\left(g_{3}\right),\left(f_{2}\right),\left(f_{3}\right)$, $(P),(Q)$ and $\left(\Omega_{1}\right)$ to obtain

$$
\begin{align*}
c= & J_{\varepsilon}\left(u_{n}\right)-\frac{1}{2}\left\langle J_{\varepsilon}^{\prime}\left(u_{n}\right), u_{n}\right\rangle+o_{n}(1)  \tag{2.23}\\
= & \int_{\mathbb{R}^{N} \backslash \Omega}\left(\frac{1}{2} g\left(x, u_{n}\right) u_{n}-G\left(x, u_{n}\right)\right) d x \\
& +\int_{\Omega} P(x)\left(\frac{1}{2} f\left(u_{n}\right) u_{n}-F\left(u_{n}\right) d x+\frac{s}{N} \int_{\Omega} Q(x) u_{n}^{2_{s}^{*}} d x+o_{n}(1)\right. \\
\geq & \alpha \int_{\Omega}\left(\frac{1}{2}-\frac{1}{\theta}\right) f\left(u_{n}\right) u_{n} d x+\frac{s}{N} \int_{\Omega}\left(Q(x)-Q_{0}\right) u_{n}^{2_{s}^{*}} d x+\frac{s}{N} \int_{\Omega} Q_{0} u_{n}^{2_{s}^{*}} d x+o_{n}(1) \\
\geq & \alpha\left(\frac{1}{2}-\frac{1}{\theta}\right) C_{0} \int_{\Omega} u_{n}^{q} d x-\frac{s}{N} \int_{\Omega}\left(Q_{0}-Q(x)\right) u_{n}^{2_{s}^{*}} d x+\frac{s}{N} \int_{\Omega} Q_{0} u_{n}^{2_{s}^{*}} d x+o_{n}(1) \\
\geq & \frac{s}{N} \int_{\Omega} Q_{0} u_{n}^{2_{s}^{*}} d x+o_{n}(1) \tag{2.24}
\end{align*}
$$

Since $Q \in L^{\infty}\left(\mathbb{R}^{N}\right), Q_{0}=\sup _{\mathbb{R}^{N}} Q$ and $\left\{u_{n}\right\}$ is bounded, then the last inequality holds for sufficiently large $C_{0}$. In [13], although it is not stated in the condition $\left(f_{2}\right)$ that the $C_{0}$ used is also sufficiently large. Taking the limit and from (2.18) one has

$$
\begin{equation*}
c \geq \frac{s}{N} Q_{0} \sum_{\left\{j \in J: x_{j} \in \Omega\right\}} \psi_{\rho}\left(x_{j}\right) \nu_{j}=\frac{s}{N} Q_{0} \sum_{\left\{j \in J: x_{j} \in \Omega\right\}} \nu_{j} \geq \varepsilon^{N} \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}} . \tag{2.25}
\end{equation*}
$$

which yields a contradiction. Hence, Claim 2 holds true and the lemma is proved.
Lemma 2.8. The functional $J_{\varepsilon}$ possesses a positive critical point $u_{\varepsilon} \in H_{\varepsilon}$ such that $J_{\varepsilon}\left(u_{\varepsilon}\right)=c_{\varepsilon}$ for $\varepsilon$ small.

Proof. Let $x_{0} \in \Omega$ be such that $V\left(x_{0}\right)=V_{0}$. By Lemma 2.3, we know that problem $\left(P_{V_{0}}\right)$ has a positive ground state solution. Let $w \in H_{0}$ be a least energy solution of problem ( $P_{V_{0}}$ ), then,

$$
c_{0}:=J_{0}(w)=\inf _{u \in H_{0} \backslash\{0\}} \sup _{t \geq 0} J_{0}(t u)=\inf _{u \in \mathcal{N}_{0}} J_{0}(u)
$$

Set $\tilde{w}(x):=w\left(\frac{x-x_{0}}{\varepsilon}\right)$. Then $c_{\varepsilon} \leq \sup _{t>0} J_{\varepsilon}(t \tilde{w})=J_{\varepsilon}\left(t_{0} \tilde{w}\right)$ for some $t_{0}>0$. We have

$$
\begin{aligned}
J_{\varepsilon}\left(t_{0} \tilde{w}\right)=\varepsilon^{N} & {\left[\frac{t_{0}^{2}}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} w\right|^{2}+V\left(x_{0}+\varepsilon x\right) w^{2}\right) d x-\int_{\mathbb{R}^{N}} G\left(x_{0}+\varepsilon x, t_{0} w\right) d x\right] } \\
=\varepsilon^{N} & {\left[J_{0}\left(t_{0} w\right)+\frac{t_{0}^{2}}{2} \int_{\mathbb{R}^{N}}\left(V\left(x_{0}+\varepsilon x\right)-V_{0}\right) w^{2} d x+\int_{\mathbb{R}^{N}} P_{0} F\left(t_{0} w\right) d x\right.} \\
& \left.+\frac{Q_{0}}{2_{s}^{*}} \int_{\mathbb{R}^{N}}\left|t_{0} w\right|^{2_{s}^{*}} d x-\int_{\mathbb{R}^{N}} G\left(x_{0}+\varepsilon x, t_{0} w\right) d x\right] .
\end{aligned}
$$

From the Lebesgue's theorem we have $\int_{\mathbb{R}^{N}}\left(V\left(x_{0}+\varepsilon x\right)-V_{0}\right) w^{2} d x \rightarrow 0$ as $\varepsilon \rightarrow 0^{+}$. From $\left(f_{1}\right),\left(f_{2}\right)$ and $\left(g_{2}\right)$ we have $G\left(x, t_{0} w\right) \leq C_{1}|w|^{2}+C_{2}|w|^{2_{s}^{*}}$. Again, by the Lebesgue's theorem the following convergence holds

$$
\int_{\mathbb{R}^{N}} G\left(x_{0}+\varepsilon x, t_{0} w\right) d x \rightarrow \int_{\mathbb{R}^{N}} P_{0} F\left(t_{0} w\right)+\frac{Q_{0}}{2_{s}^{*}} \int_{\mathbb{R}^{N}}\left|t_{0} w\right|^{2_{s}^{*}} d x
$$

as $\varepsilon \rightarrow 0^{+}$. Hence

$$
\begin{equation*}
c_{\varepsilon} \leq J_{\varepsilon}\left(t_{0} \tilde{w}\right)=\varepsilon^{N}\left(J_{0}\left(t_{0} w\right)+o(1)\right) \leq \varepsilon^{N}\left(c_{0}+o(1)\right) . \tag{2.26}
\end{equation*}
$$

By Lemma 2.1, we get $c_{\varepsilon}<\varepsilon^{N} \frac{s}{N Q_{0}^{\frac{N-2 s}{2 s}}} S_{*}^{\frac{N}{2 s}}$ for $\varepsilon$ small enough. Since $J_{\varepsilon}$ satisfies the mountain-pass geometry conditions. Thus there exists a sequence $\left\{u_{n}\right\} \subset H_{0}$ such that $J_{\varepsilon}\left(u_{n}\right) \rightarrow c_{\varepsilon}$ and $J_{\varepsilon}^{\prime}\left(u_{n}\right) \rightarrow 0$. From Lemma 2.7, we obtain $J_{\varepsilon}$ satisfies the PalaisSmale condition at level $c_{\varepsilon}$. Then, the functional $J_{\varepsilon}$ possesses a nontrivial critical point $u_{\varepsilon} \in H_{\varepsilon}$ such that $J_{\varepsilon}\left(u_{\varepsilon}\right)=c_{\varepsilon}$, similar to the proof of Lemma 2.3, we get $u_{\varepsilon} \in H_{\varepsilon}$ is a positive critical point such that $J_{\varepsilon}\left(u_{\varepsilon}\right)=c_{\varepsilon}$.

### 2.3 Proof of theorem 1.1

Next we shall prove our main result. The idea is to show that the solution obtained in Lemma 2.8 satisfy the estimate $u_{\varepsilon} \leq a, \forall x \in \Omega$ for $\varepsilon$ small enough. This fact implies that the solution is indeed solution of the original problem $\left(P_{\varepsilon}\right)$.

Lemma 2.9. There is $C>0$ such that

$$
\int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{\varepsilon}\right|^{2}+V(x)\left|u_{\varepsilon}\right|^{2}\right) d x \leq C \varepsilon^{N} .
$$

Proof. Indeed, we have $\left\langle J_{\varepsilon}^{\prime}\left(u_{\varepsilon}\right), u_{\varepsilon}\right\rangle=0$, that is,

$$
\int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{\varepsilon}\right|^{2}+V(x)\left|u_{\varepsilon}\right|^{2}\right) d x=\int_{\mathbb{R}^{N}} g\left(x, u_{\varepsilon}\right) u_{\varepsilon} d x .
$$

By 2.26 ) and $\left(g_{3}\right)$, we have

$$
\begin{aligned}
& \frac{1}{2} \int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{\varepsilon}\right|^{2}+V(x)\left|u_{\varepsilon}\right|^{2}\right) d x \\
= & J_{\varepsilon}\left(u_{\varepsilon}\right)+\int_{\mathbb{R}^{N}} G\left(x, u_{\varepsilon}\right) d x \\
\leq & \varepsilon^{N}\left(c_{0}+o(1)\right)+\frac{1}{\theta} \int_{\Omega} g\left(x, u_{\varepsilon}\right) u_{\varepsilon} d x+\frac{1}{2}\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right) \int_{\mathbb{R}^{N}} V(x)\left|u_{\varepsilon}\right|^{2} d x \\
\leq & C_{1} \varepsilon^{N}+\left(\frac{1}{\theta}+\frac{1}{2}\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right)\right) \int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{\varepsilon}\right|^{2}+V(x)\left|u_{\varepsilon}\right|^{2}\right) d x .
\end{aligned}
$$

Therefore, we have

$$
\left(\frac{1}{2}-\frac{1}{\theta}-\frac{1}{2}\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right)\right) \int_{\mathbb{R}^{N}}\left(\varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}} u_{\varepsilon}\right|^{2}+V(x)\left|u_{\varepsilon}\right|^{2}\right) d x \leq C \varepsilon^{N} .
$$

Moreover $\frac{1}{2}-\frac{1}{\theta}-\frac{1}{2}\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right)=\frac{1}{2}\left(\frac{\theta-2}{\theta}-\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right)\right)>0$ and the proof is completed.
Lemma 2.10. If $\varepsilon_{n} \rightarrow 0^{+}$and $\left\{x_{n}\right\} \subset \bar{\Omega}$ are such that $u_{\varepsilon_{n}}\left(x_{n}\right) \geq \gamma>0$, then $\lim _{n \rightarrow \infty} V\left(x_{n}\right)=V_{0}$.
Proof. Assume by contradiction, passing to a subsequence, that $x_{n} \rightarrow \bar{x} \in \bar{\Omega}$ and $V(\bar{x})>V_{0}$. Let $v_{n}(x):=u_{\varepsilon_{n}}\left(x_{n}+\varepsilon_{n} x\right)$. Obviously, $v_{n} \in H_{\varepsilon}$ satisfy the following equation

$$
\begin{equation*}
(-\Delta)^{s} v_{n}+V\left(x_{n}+\varepsilon_{n} x\right) v_{n}=g\left(x_{n}+\varepsilon_{n} x, v_{n}\right) \text { in } \mathbb{R}^{N} . \tag{2.27}
\end{equation*}
$$

The associated energy functional is given by

$$
J_{n}(u)=\frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+V\left(\varepsilon_{n} x+x_{n}\right) u^{2}\right) d x-\int_{\mathbb{R}^{N}} G\left(\varepsilon_{n} x+x_{n}, u\right) d x .
$$

From Lemma 2.9, we have that $\left\{v_{n}\right\}$ is bounded in $H_{\varepsilon}$ and therefore $v_{n} \rightharpoonup v$ in $H_{\varepsilon}$ for some $v \in H_{\varepsilon}$. From $\left(f_{1}\right)-\left(f_{3}\right)$, it is easy to get that $J_{0}\left(t v_{n}\right)>0$ for $t>0$ small enough and $J_{0}\left(t v_{n}\right) \rightarrow-\infty$ as $t \rightarrow \infty$. Thus, there exists $t_{n}>0$ such that $J_{0}\left(t_{n} v_{n}\right)=$ $\max _{t \geq 0} J_{0}\left(t v_{n}\right)$. Set $\tilde{v}_{n}:=t_{n} v_{n}$, therefore, $c_{0} \leq J_{0}\left(\tilde{v}_{n}\right)$. Since $\left\{v_{n}\right\}$ satisfy equation 2.27), we have $J_{n}\left(v_{n}\right)=\max _{t \geq 0} J_{n}\left(t v_{n}\right)$, thus

$$
\begin{aligned}
c_{0} \leq J_{0}\left(\tilde{v}_{n}\right)= & \frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}_{n}\right|^{2}+V_{0} \tilde{v}_{n}^{2}\right) d x-\int_{\mathbb{R}^{N}} P_{0} F\left(\tilde{v}_{n}\right) d x-\frac{Q_{0}}{2_{s}^{*}} \int_{\mathbb{R}^{N}}\left|\tilde{v}_{n}\right|^{2_{s}^{*}} d x \\
\leq & \frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}_{n}\right|^{2}+V\left(\varepsilon_{n} x+x_{n}\right) \tilde{v}_{n}^{2}\right) d x-\int_{\mathbb{R}^{N}} P\left(\varepsilon_{n} x+x_{n}\right) F\left(\tilde{v}_{n}\right) d x \\
& -\frac{1}{2_{s}^{*}} \int_{\mathbb{R}^{N}} Q\left(\varepsilon_{n} x+x_{n}\right)\left|\tilde{v}_{n}\right|^{2_{s}^{*}} d x \\
\leq & \frac{t_{n}^{2}}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} v_{n}\right|^{2}+V\left(\varepsilon_{n} x+x_{n}\right) v_{n}^{2}\right) d x-\int_{\mathbb{R}^{N}} G\left(\varepsilon_{n} x+x_{n}, t_{n} v_{n}\right) d x \\
= & J_{n}\left(t_{n} v_{n}\right) \leq J_{n}\left(v_{n}\right)=\varepsilon_{n}^{-N} J_{\varepsilon_{n}}\left(u_{\varepsilon_{n}}\right) \leq c_{0}+o(1),
\end{aligned}
$$

which implies $\lim _{n \rightarrow \infty} J_{0}\left(\tilde{v}_{n}\right)=c_{0}$, where the last inequality is from (2.26). Moreover, $\left\{\tilde{v}_{n}\right\}$ is bounded and $\tilde{v}_{n} \rightharpoonup \tilde{v}$. We claim that $\left\{\tilde{v}_{n}\right\}$ satisfies the following limits

$$
J_{0}\left(\tilde{v}_{n}\right) \rightarrow c_{0} \text { and } J_{0}^{\prime}\left(\tilde{v}_{n}\right) \rightarrow 0
$$

In fact, using Ekeland's variational Principle in [10], there exists a sequence $\left\{\nu_{n}\right\} \subset \mathcal{N}_{0}$ satisfying $\nu_{n}=\tilde{v}_{n}+o_{n}(1), J_{0}\left(\nu_{n}\right) \rightarrow c_{0}$ and $J_{0}^{\prime}\left(\nu_{n}\right)-\lambda_{n} \Phi^{\prime}\left(\nu_{n}\right)=o_{n}(1)$, where $\lambda_{n}$ is a real number and $\Phi\left(\nu_{n}\right)=\left\langle J_{0}^{\prime}\left(\nu_{n}\right), \nu_{n}\right\rangle$. Thus, by the definition of $\Phi\left(\nu_{n}\right)$ and $\left\{\nu_{n}\right\} \subset \mathcal{N}_{0}$, we have that

$$
\left\langle\Phi^{\prime}\left(\nu_{n}\right), \nu_{n}\right\rangle=\int_{\mathbb{R}^{N}} P_{0}\left(f\left(\nu_{n}\right) \nu_{n}-f^{\prime}\left(\nu_{n}\right)\left|\nu_{n}\right|^{2}\right) d x-\left(2_{s}^{*}-2\right) \int_{\mathbb{R}^{N}} Q_{0}\left|\nu_{n}\right|^{2_{s}^{*}} d x
$$

$$
\begin{equation*}
\leq \int_{\mathbb{R}^{N}} P_{0}\left(f\left(\nu_{n}\right) \nu_{n}-f^{\prime}\left(\nu_{n}\right)\left|\nu_{n}\right|^{2}\right) d x \text {. } \tag{2.28}
\end{equation*}
$$

Since $\left\{\nu_{n}\right\}$ is bounded and $\nu_{n} \nrightarrow 0$, Lemma 2.2 guarantees the existence of a sequence $\left\{y_{n}\right\} \subset \mathbb{R}^{N}$ such that $\tilde{\nu}_{n}=\nu_{n}\left(\cdot+y_{n}\right)$ is a bounded sequence in $H_{0}$ and $\tilde{\nu}_{n} \rightharpoonup \tilde{\nu}$ for some $\tilde{\nu} \neq 0$. Hence, there exists a subset $\Lambda \subset \mathbb{R}^{N}$ having positive measure, such that $\tilde{\nu}>0$ a.e. in $\Lambda$. Assume by contradiction that $\lim \sup _{n \rightarrow \infty}\left\langle\Phi^{\prime}\left(\nu_{n}\right), \nu_{n}\right\rangle=0$. Then, taking into account (2.28), $\left(f_{4}\right)$ and Fatou's Lemma, we get $0>\int_{\Lambda}\left(f(\tilde{\nu}) \tilde{\nu}-f^{\prime}(\tilde{\nu})|\tilde{\nu}|^{2}\right) \geq 0$ which gives a contradiction. Hence $\lim \sup _{n \rightarrow \infty}\left\langle\Phi^{\prime}\left(\nu_{n}\right), \nu_{n}\right\rangle<0$, implying that $\lambda_{n}=o_{n}(1)$, thus $J_{0}\left(\nu_{n}\right) \rightarrow c_{0}, J_{0}^{\prime}\left(\nu_{n}\right) \rightarrow 0$. Without loss of generalization, we may assume that $J_{0}\left(\tilde{v}_{n}\right) \rightarrow c_{0}, J_{0}^{\prime}\left(\tilde{v}_{n}\right) \rightarrow 0$. Thus, $J_{0}^{\prime}(\tilde{v})=0$, by Fatou's Lemma we get

$$
\begin{aligned}
c_{0} \leq J_{0}(\tilde{v})= & J_{0}(\tilde{v})-\frac{1}{\theta}\left\langle J_{0}^{\prime}(\tilde{v}), \tilde{v}\right\rangle \\
= & \left(\frac{1}{2}-\frac{1}{\theta}\right)\left(\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}\right|^{2}+V_{0} \tilde{v}^{2}\right) d x\right) \\
& +P_{0} \int_{\mathbb{R}^{N}}\left(\frac{1}{\theta} f(\tilde{v}) \tilde{v}-F(\tilde{v})\right) d x+\left(\frac{1}{\theta}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{N}} Q_{0}|\tilde{v}|^{2_{s}^{*}} d x \\
\leq & \liminf _{n \rightarrow \infty}\left(\frac{1}{2}-\frac{1}{\theta}\right)\left(\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}_{n}\right|^{2}+V_{0} \tilde{v}_{n}^{2}\right) d x\right) \\
& +P_{0} \int_{\mathbb{R}^{N}}\left(\frac{1}{\theta} f\left(\tilde{v}_{n}\right) \tilde{v}-F\left(\tilde{v}_{n}\right)\right) d x+\left(\frac{1}{\theta}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{N}} Q_{0}\left|\tilde{v}_{n}\right|^{2_{s}^{*}} d x \\
= & \liminf _{n \rightarrow \infty}\left(J_{0}\left(\tilde{v}_{n}\right)-\frac{1}{\theta}\left\langle J_{0}^{\prime}\left(\tilde{v}_{n}\right), \tilde{v}_{n}\right\rangle\right) \leq c_{0} .
\end{aligned}
$$

Then, $\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}_{n}\right|^{2}+V_{0} \tilde{v}_{n}^{2}\right) d x=\int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}\right|^{2}+V_{0} \tilde{v}^{2}\right) d x$. Hence, $\tilde{v}_{n} \rightarrow \tilde{v}$ in $H^{s}\left(\mathbb{R}^{N}\right)$.

Taking into account that $V(\bar{x})>V_{0}, P(\bar{x}) \leq P_{0}, Q(\bar{x}) \leq Q_{0}, \tilde{v}_{n} \rightarrow \tilde{v}$ and Fatou's Lemma, we obtain

$$
\begin{aligned}
c_{0}=J_{0}(\tilde{v})< & \liminf _{n \rightarrow \infty}\left\{\frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} \tilde{v}_{n}\right|^{2}+V\left(\varepsilon_{n} x+x_{n}\right) \tilde{v}_{n}^{2}\right) d x\right. \\
& \left.\quad-\int_{\mathbb{R}^{N}} P\left(\varepsilon_{n} x+x_{n}\right) F\left(\tilde{v}_{n}\right) d x-\frac{1}{2_{s}^{*}} \int_{\mathbb{R}^{N}} Q\left(\varepsilon_{n} x+x_{n}\right)\left|\tilde{v}_{n}\right|^{2_{s}^{*}} d x\right\} \\
\leq & \liminf _{n \rightarrow \infty}\left\{\frac{t_{n}^{2}}{2} \int_{\mathbb{R}^{N}}\left(\left|(-\Delta)^{\frac{s}{2}} v_{n}\right|^{2}+V\left(\varepsilon_{n} x+x_{n}\right) v_{n}^{2}\right) d x\right. \\
& \left.\quad-\int_{\mathbb{R}^{N}} G\left(\varepsilon_{n} x+x_{n}, t_{n} v_{n}\right) d x\right\} \\
= & \liminf _{n \rightarrow \infty} J_{n}\left(t_{n} v_{n}\right) \leq \liminf _{n \rightarrow \infty} J_{n}\left(v_{n}\right)=\liminf _{n \rightarrow \infty} \varepsilon_{n}^{-N} J_{\varepsilon_{n}}\left(u_{\varepsilon_{n}}\right) \leq c_{0}+o(1),
\end{aligned}
$$

which yields a contradiction. Thus the proof is completed.
Lemma 2.11. There holds $\lim _{\varepsilon \rightarrow 0^{+}} m_{\varepsilon}=0$, where $m_{\varepsilon}:=\max _{\partial \Omega} u_{\varepsilon}$.

Proof. Assume by contradiction that $m_{\varepsilon} \nrightarrow 0$. Let $x_{\varepsilon} \in \partial \Omega \subset \bar{\Omega}$ be such that $u_{\varepsilon}\left(x_{\varepsilon}\right)=$ $m_{\varepsilon}$. Therefore, up to a subsequence we have $u_{\varepsilon_{n}}\left(x_{\varepsilon_{n}}\right) \geq \gamma>0$ and $x_{\varepsilon_{n}} \rightarrow x_{0} \in \partial \Omega$, by Lemma 2.10 we have

$$
\min _{\partial \Omega} V \leq \lim _{n \rightarrow \infty} V\left(x_{\varepsilon_{n}}\right)=V\left(x_{0}\right)=V_{0}<\min _{\partial \Omega} V,
$$

which gets a contradiction.
Proof of theorem 1.1. Let $u_{\varepsilon}$ be a positive critical point for $J_{\varepsilon}$. From Lemma 2.11, there exists $\varepsilon_{0}$ such that for any $\varepsilon \in\left(0, \varepsilon_{0}\right), m_{\varepsilon}<a$. Therefore $u_{\varepsilon}(x)<a$ for $x \in \partial \Omega$. Thus, in view of the maximum principle, we obtain

$$
u_{\varepsilon}(x) \leq a \text { for } x \in \Omega .
$$

Taking $\left(u_{\varepsilon}-a\right)_{+}=\max \left\{u_{\varepsilon}-a, 0\right\}$ as a test function for $J_{\varepsilon}$, we have

$$
\begin{align*}
0=J_{\varepsilon}^{\prime}\left(u_{\varepsilon}\right)\left(\left(u_{\varepsilon}-a\right)_{+}\right)= & \int_{\mathbb{R}^{N} \backslash \Omega} \varepsilon^{2 s}\left|(-\Delta)^{\frac{s}{2}}\left(u_{\varepsilon}-a\right)_{+}\right|^{2} d x \\
& +\int_{\mathbb{R}^{N} \backslash \Omega} c(x)\left(u_{\varepsilon}-a\right)_{+}^{2}+c(x) a\left(u_{\varepsilon}-a\right)_{+} d x \tag{2.29}
\end{align*}
$$

where $c(x)=V(x)-\frac{g\left(x, u_{\varepsilon}\right)}{u_{\varepsilon}}$. Moreover, for $x \in \mathbb{R}^{N} \backslash \Omega$, we get $\frac{g\left(x, u_{\varepsilon}\right)}{u_{\varepsilon}} \leq\left(\frac{P_{0}}{K}+\frac{Q_{0}}{L}\right) V_{0}$. Hence, $c(x)>0$ for $x \in \mathbb{R}^{N} \backslash \Omega$. So every term in the last identity of (2.29) is 0 . Therefore, $\left(u_{\varepsilon}-a\right)_{+}=0$ and $u_{\varepsilon}(x) \leq a$ for $x \in \mathbb{R}^{N} \backslash \Omega$. Thus, $g\left(x, u_{\varepsilon}\right)=P(x) f\left(u_{\varepsilon}\right)+$ $Q(x)\left|u_{\varepsilon}\right|^{2_{s}^{*}-2} u_{\varepsilon}$ and $u_{\varepsilon}$ is a solution of $\left(P_{\varepsilon}\right)$. Similar arguments as [13], we can obtained that if $\eta_{\varepsilon} \in \mathbb{R}^{N}$ denotes the global maximum point of $u_{\varepsilon}$, then

$$
u_{\varepsilon}(x) \leq \frac{C \varepsilon^{N+2 s}}{\varepsilon^{N+2 s}+\left|x-\eta_{\varepsilon}\right|^{N+2 s}}
$$

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